## A Fourier derivative collocation method for the solution of the Navier–Stokes problem

Daniel Thomas Hayes, dthayes83@gmail.com

September 26, 2024

A proposed solution to the millennium problem on the existence and smoothness of the Navier–Stokes equations.

## 1. Introduction

The Navier–Stokes equations are thought to govern the motion of a fluid in  $\mathbb{R}^3$ , [1–5]. Let  $\mathbf{u} = \mathbf{u}(\mathbf{x},t) \in \mathbb{R}^3$  be the fluid velocity and let  $p = p(\mathbf{x},t) \in \mathbb{R}$  be the fluid pressure, each dependent on position  $\mathbf{x} \in \mathbb{R}^3$  and time  $t \ge 0$ . I take the externally applied force acting on the fluid to be identically zero. The fluid is assumed to be incompressible with constant viscosity v > 0 and to fill all of  $\mathbb{R}^3$ . The Navier–Stokes equations can then be written as

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla)\mathbf{u} = \nu \nabla^2 \mathbf{u} - \nabla p, \tag{1}$$

$$\nabla \cdot \mathbf{u} = 0 \tag{2}$$

with initial condition

$$\mathbf{u}(\mathbf{x},0) = \mathbf{u}^{\circ} \tag{3}$$

where  $\mathbf{u}^{\circ} = \mathbf{u}^{\circ}(\mathbf{x}) \in \mathbb{R}^{3}$ . In these equations

$$\nabla = \left(\frac{\partial}{\partial \mathbf{x}_1}, \frac{\partial}{\partial \mathbf{x}_2}, \frac{\partial}{\partial \mathbf{x}_3}\right) \tag{4}$$

is the gradient operator and

$$\nabla^2 = \sum_{i=1}^3 \frac{\partial^2}{\partial \mathbf{x}_i^2} \tag{5}$$

is the Laplacian operator. When  $\nu = 0$  equations (1), (2), (3) are called the Euler equations. When  $\nabla p = \mathbf{0}$  equations (1), (3) are called the Burgers equations. Solutions of (1), (2), (3) are to be found with

$$\mathbf{u}^{\circ}(\mathbf{x} + e_i) = \mathbf{u}^{\circ}(\mathbf{x}) \tag{6}$$

for  $1 \le i \le 3$  where  $e_i$  is the  $i^{th}$  unit vector in  $\mathbb{R}^3$ . The initial condition  $\mathbf{u}^{\circ}$  is a given  $C^{\infty}$  divergence-free vector field on  $\mathbb{R}^3$ . A solution of (1), (2), (3) is then accepted to be physically reasonable [3] if

$$\mathbf{u}(\mathbf{x} + e_i, t) = \mathbf{u}(\mathbf{x}, t), \quad p(\mathbf{x} + e_i, t) = p(\mathbf{x}, t) \tag{7}$$

on  $\mathbb{R}^3 \times [0, \infty)$  for  $1 \le i \le 3$  and

$$\mathbf{u}, p \in C^{\infty} \left( \mathbb{R}^3 \times [0, \infty) \right). \tag{8}$$

## 2. Solution of the Navier-Stokes problem

**Theorem**. Take  $\nu > 0$ . Let  $\mathbf{u}^{\circ}$  be any smooth, divergence-free vector field satisfying (6). Then there exist smooth functions  $\mathbf{u}$ , p on  $\mathbb{R}^3 \times [0, \infty)$  that satisfy (1), (2), (3), (7), (8).

**Proof.** A Fourier derivative collocation method is as follows. Let **u**, p be given by

$$\mathbf{u} = \sum_{\mathbf{L} = -\infty}^{\infty} \mathbf{u}_{\mathbf{L}} e^{ik\mathbf{L} \cdot \mathbf{x}},\tag{9}$$

$$p = \sum_{\mathbf{L} = -\infty}^{\infty} p_{\mathbf{L}} e^{ik\mathbf{L} \cdot \mathbf{x}}$$
 (10)

respectively. Here  $\mathbf{u}_{\mathbf{L}} = \mathbf{u}_{\mathbf{L}}(t) \in \mathbb{C}^3$ ,  $p_{\mathbf{L}} = p_{\mathbf{L}}(t) \in \mathbb{C}$ ,  $\mathbf{i} = \sqrt{-1}$ ,  $k = 2\pi$ , and  $\sum_{\mathbf{L} = -\infty}^{\infty}$  denotes the sum over all  $\mathbf{L} \in \mathbb{Z}^3$ . The initial condition  $\mathbf{u}^{\circ}$  is a Fourier series [2] of which is convergent for all  $\mathbf{x} \in \mathbb{R}^3$ . Equations (1), (2) can be written as

$$\frac{\partial \mathbf{u}_{i}}{\partial t} + \sum_{i=1}^{3} \mathbf{u}_{j} \frac{\partial \mathbf{u}_{i}}{\partial \mathbf{x}_{j}} = \nu \sum_{i=1}^{3} \frac{\partial^{2} \mathbf{u}_{i}}{\partial \mathbf{x}_{j}^{2}} - \frac{\partial p}{\partial \mathbf{x}_{i}} \text{ for } i = 1, 2, 3,$$
(11)

and

$$\sum_{j=1}^{3} \frac{\partial \mathbf{u}_{j}}{\partial \mathbf{x}_{j}} = 0 \tag{12}$$

respectively. In this method we have for a quantity q that

$$\left[\frac{\partial q}{\partial \mathbf{x}_j}\right] = \left[G_j\right] \lceil q \rceil \tag{13}$$

valid at  $\mathbf{x} = \mathbf{x}^*_n$  for n = 1, 2, ..., N. Here  $\left[G_j\right]$  is a known constant  $N \times N$  matrix with  $\left[G_j\right]_{m,n} = G_{j,m,n}$  and  $\lceil r \rceil$  means to vectorise r where the components are equal to  $r|_{\mathbf{x}=\mathbf{x}^*_n}$ , n = 1, 2, ..., N. We denote  $q|_{\mathbf{x}=\mathbf{x}^*_n} = \lceil q \rceil_n = q_n$ . Then

$$\left[\frac{\partial \mathbf{u}_{i}}{\partial \mathbf{x}_{j}}\right]_{n} = \sum_{\alpha=1}^{N} G_{j,n,\alpha} u_{i,\alpha}, \quad \left[\frac{\partial p}{\partial \mathbf{x}_{i}}\right]_{n} = \sum_{\alpha=1}^{N} G_{i,n,\alpha} p_{,\alpha}, \quad (14)$$

$$\left[\frac{\partial^2 \mathbf{u}_i}{\partial \mathbf{x}_j^2}\right]_n = \sum_{\alpha=1}^N G_{j,n,\alpha} \left[\frac{\partial \mathbf{u}_i}{\partial \mathbf{x}_j}\right]_{\alpha} = \sum_{\alpha=1}^N \sum_{\beta=1}^N G_{j,n,\alpha} G_{j,\alpha,\beta} u_{i,\beta},\tag{15}$$

and

$$\left[\frac{\partial \mathbf{u}_i}{\partial t}\right]_n = \frac{\partial}{\partial t} \left[\mathbf{u}_i\right]_n = \frac{\partial}{\partial t} u_{i,n}.$$
 (16)

Equations (11), (12) at  $\mathbf{x} = \mathbf{x}^*_n$  imply

$$\frac{\partial}{\partial t} \left[ \mathbf{u}_i \right]_n + \sum_{j=1}^3 \left[ \mathbf{u}_j \right]_n \left[ \frac{\partial \mathbf{u}_i}{\partial \mathbf{x}_j} \right]_n = \nu \sum_{j=1}^3 \left[ \frac{\partial^2 \mathbf{u}_i}{\partial \mathbf{x}_j^2} \right]_n - \left[ \frac{\partial p}{\partial \mathbf{x}_i} \right]_n$$
(17)

and

$$\sum_{j=1}^{3} \left[ \frac{\partial \mathbf{u}_{j}}{\partial \mathbf{x}_{j}} \right]_{n} = 0 \tag{18}$$

respectively. Equations (17), (18) imply

$$\frac{\partial}{\partial t}u_{i,n} + \sum_{j=1}^{3} \sum_{\alpha=1}^{N} u_{j,n}G_{j,n,\alpha}u_{i,\alpha} = \nu \sum_{j=1}^{3} \sum_{\alpha=1}^{N} \sum_{\beta=1}^{N} G_{j,n,\alpha}G_{j,\alpha,\beta}u_{i,\beta} - \sum_{\alpha=1}^{N} G_{i,n,\alpha}p_{,\alpha}$$
(19)

and

$$\sum_{j=1}^{3} \sum_{\alpha=1}^{N} G_{j,n,\alpha} u_{j,\alpha} = 0$$
 (20)

respectively. Let U be a matrix where  $U_{i,n} = u_{i,n}$  and let P be a matrix where  $P_{\alpha,n} = p_{,\alpha}$ . Then equations (19), (20) imply

$$\frac{\partial U}{\partial t} + U(A(n)U) = \nu U B(n) - A(n)^T P \tag{21}$$

and

$$trace(UA(n)) = 0 (22)$$

respectively. Herein A(n) and B(n) are matrices where

$$A(n)_{\alpha,j} = G_{j,n,\alpha} \tag{23}$$

and

$$B(n)_{\beta,n} = \sum_{i=1}^{3} \sum_{\alpha=1}^{N} G_{j,n,\alpha} G_{j,\alpha,\beta}.$$
 (24)

The i, n component of (21) recovers (19) since

$$[U(A(n)U)]_{i,n} = \sum_{l=1}^{N} U_{i,l}[A(n)U]_{l,n} = \sum_{l=1}^{N} U_{i,l} \left[ \sum_{m=1}^{3} A(n)_{l,m} U_{m,n} \right]$$

$$= \sum_{j=1}^{3} \sum_{\alpha=1}^{N} U_{i,\alpha} A(n)_{\alpha,j} U_{j,n} = \sum_{j=1}^{3} \sum_{\alpha=1}^{N} U_{i,\alpha} G_{j,n,\alpha} U_{j,n}, \quad (25)$$

$$[UB(n)]_{i,n} = \sum_{l=1}^{N} U_{i,l}B(n)_{l,n} = \sum_{\beta=1}^{N} U_{i,\beta}B(n)_{\beta,n} = \sum_{i=1}^{3} \sum_{\alpha=1}^{N} \sum_{\beta=1}^{N} U_{i,\beta}G_{j,n,\alpha}G_{j,\alpha,\beta}, \quad (26)$$

and

$$\begin{aligned}
\left[A(n)^{T}P\right]_{i,n} &= \sum_{l=1}^{N} A(n)_{i,l}^{T} P_{l,n} = \sum_{l=1}^{N} A(n)_{l,i} P_{l,n} = \sum_{l=1}^{N} G_{i,n,l} P_{l,n} \\
&= \sum_{\alpha=1}^{N} G_{i,n,\alpha} P_{\alpha,n} = \sum_{\alpha=1}^{N} G_{i,n,\alpha} P_{\alpha,\alpha}.
\end{aligned} (27)$$

Equation (22) recovers (20) since

trace(
$$UA(n)$$
) =  $\sum_{j=1}^{3} [UA(n)]_{j,j} = \sum_{j=1}^{3} \sum_{l=1}^{N} U_{j,l} A(n)_{l,j}$   
=  $\sum_{j=1}^{3} \sum_{l=1}^{N} U_{j,l} G_{j,n,l} = \sum_{j=1}^{3} \sum_{\alpha=1}^{N} U_{j,\alpha} G_{j,n,\alpha}$ . (28)

Let Q(n) be a matrix such that  $(A(n)^T P)Q(n) = 0$ . We have

$$\left[ \left( A(n)^T P \right) Q(n) \right]_{i,j} = \left[ A(n)^T \left( P Q(n) \right) \right]_{i,j} = \sum_{l=1}^N A(n)_{i,l}^T (P Q(n))_{l,j} \\
= \sum_{l=1}^N \sum_{m=1}^N A(n)_{i,l}^T P_{l,m} Q(n)_{m,j} = 0.$$
(29)

Then (21) implies

$$\left(\frac{\partial U}{\partial t}\right)Q(n) + (U(A(n)U))Q(n) = (vUB(n))Q(n). \tag{30}$$

Equation (30) is the same as we would get for the Burgers equations. Now we consider a matrix Riccati equation problem.

$$\frac{\partial X}{\partial t} = aX + bY,\tag{31}$$

$$\frac{\partial Y}{\partial t} = cX + dY,\tag{32}$$

with

$$X = (U\lambda)Y. \tag{33}$$

Then we get

$$\left(\frac{\partial U}{\partial t}\lambda\right)Y + (U\lambda)\frac{\partial Y}{\partial t} = a(U\lambda)Y + bY \tag{34}$$

which implies

$$\left(\frac{\partial U}{\partial t}\lambda\right)Y + (U\lambda)[c(U\lambda)Y + dY] = a(U\lambda)Y + bY \tag{35}$$

implying

$$\frac{\partial U}{\partial t}\lambda + (U\lambda)c(U\lambda) + (U\lambda)d = a(U\lambda) + b. \tag{36}$$

We then let a = b = 0,  $\lambda = Q(n)$ ,  $c = Q(n)^{-1}A(n)$ ,  $d = -\nu Q(n)^{-1}B(n)Q(n)$  to recover (30). Then (31) implies

$$X = X|_{t=0}. (37)$$

Equation (32) implies

$$\frac{\partial Y}{\partial t} = cX|_{t=0} + dY \tag{38}$$

and so

$$\frac{\partial}{\partial t} \left( e^{-dt} Y \right) = e^{-dt} c X|_{t=0}$$
 (39)

which integrating with respect to t yields

$$e^{-dt}Y = \int_0^t e^{-d\tau} cX|_{t=0} d\tau + Y|_{t=0}$$
 (40)

which implies

$$e^{-dt}Y = \left[ \left( e^{-d\tau} (-d)^{-1} \right) cX|_{t=0} \right]_0^t + Y|_{t=0}$$
 (41)

to obtain

$$Y = (e^{-dt})^{-1} \left[ \left( (e^{-dt} - I)(-d)^{-1} \right) cX|_{t=0} + Y|_{t=0} \right].$$
 (42)

Equation (33) then implies

$$U\lambda = X|_{t=0}Y^{-1}$$

$$= (U|_{t=0}\lambda)Y|_{t=0} \left\{ \left[ \left( \left( e^{-dt} - I \right) (-d)^{-1} \right) c(U|_{t=0}\lambda Y|_{t=0}) + Y_{t=0} \right]^{-1} e^{-dt} \right\}$$

$$= (U|_{t=0}\lambda) \left\{ \left[ \left( \left( e^{-dt} - I \right) (-d)^{-1} \right) c(U|_{t=0}\lambda) + I \right]^{-1} e^{-dt} \right\}. \tag{43}$$

No blowup is possible since the Burgers equations are regular.  $\Box$  For the Euler equations we have

$$U\lambda = (U|_{t=0}\lambda) [c(U|_{t=0}\lambda)t + I]^{-1}$$
(44)

and blowup is possible since for odd N the equation

$$\det\left(c(U|_{t=0}\lambda)t + I\right) = 0\tag{45}$$

can have a solution t where  $0 < t < \infty$ .

## References

- [1] Batchelor G. 1967. *An introduction to fluid dynamics*. Cambridge U. Press, Cambridge.
- [2] Doering C. 2009. The 3D Navier–Stokes problem. *Annu. Rev. Fluid Mech.* **41**: 109–128.
- [3] Fefferman C. 2000. Existence and smoothness of the Navier–Stokes equation. *Clay Mathematics Institute*. Official problem description.
- [4] Ladyzhenskaya O. 1969. *The mathematical theory of viscous incompressible flows*. Gordon and Breach, New York.
- [5] Tao T. 2013. Localisation and compactness properties of the Navier–Stokes global regularity problem. *Analysis and PDE*. **6**: 25–107.