

sigmoid functionization of magnetization in the spin-glass Ising model.

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Abstract

The magnetization of the spin-glass Ising model can be expressed using a sigmoid function. In the ground state, the magnetization is determined by solving a set of nonlinear simultaneous equations, each corresponding to a magnetization. As the magnetization of the ground state in the spin-glass Ising model constitutes an NP-complete problem, the P=NP problem can be reformulated as solving these nonlinear simultaneous equations. If practical computation yields result that are feasible, it can be essentially considered as P=NP. Furthermore, all interacting systems in nature can be represented by sigmoid functions, and the ground state can be obtained by solving nonlinear simultaneous equations.

1 INTRODUCTION

The Ising model is the simplest and most fundamental interaction model. Each spin (node) interacts with each other, taking states of either -1 or 1, and evolves continuously at temperature T, eventually converging to a state called the ground state at T=0. This ground state has the lowest energy in the system, with each state's expected value being the order parameter taking -1 or 1. It represents how the system settles due to interactions. In complex systems, the resultant states of interactions serve as solutions to combinatorial optimization problems (NP problems). In the real world, solutions to interaction systems and combinatorial optimization problems have wide applications across various fields. The spin-glass Ising model is a variant of the Ising model where interactions include both positive and negative components individually. It can model various interaction systems and combinatorial optimization problems. Now, the magnetization of this spin-glass Ising model has been represented using a sigmoid function. Furthermore, in the ground state, magnetization can be expressed by a set of nonlinear simultaneous equations for each magnetization. Combinatorial optimization problems (NP problems) and interacting systems can be reduced to solving these nonlinear simultaneous equations. Below, we present the derivation of these equations.

2 RESULTS

At all temperatures, the following results were obtained.

$$\langle n_i \rangle = \left\langle \frac{1}{1 + e^{-\xi_i t}} \right\rangle \quad (1)$$

$$\xi_i = \varepsilon_{i,i} + \sum_{j=1 \neq i}^N \varepsilon_{i,j} n_j, \quad t = \frac{1}{k_B T} \quad (2)$$

Also, at the ground state (T=0), the following results were obtained.

$$\langle n_i \rangle = \frac{1}{1 + e^{-\tilde{\xi}_i t}} \quad (3)$$

$$\tilde{\xi}_i = \varepsilon_{i,i} + \sum_{j=1 \neq i}^N \varepsilon_{i,j} \langle n_j \rangle, \quad t = \frac{1}{k_B T} \quad (4)$$

3 THEORY

Let there be state variables for each node $i = 1$ to N , each with the following states.

$$n_i = 0 \text{ or } 1 \quad (5)$$

The model, represented by the following Hamiltonian of the system, is called the Ising model [1]. (Although the Ising model is generally characterized by $n = -1$ or 1 , it can be simplified for calculations as $n = 0$ or 1 .)

$$-H = \sum_{i \leq j} \varepsilon_{i,j} n_i n_j \quad (6)$$

However, ε_{ij} can take any value within the specified range (spin-glass Ising model [1]).

$$\varepsilon_{i,j} = -1 \sim 1 \quad (7)$$

The order parameter $\langle n_i \rangle$ is determined by the following.

$$\langle n_i \rangle = \frac{Z_i}{Z} \quad (8)$$

$$Z = \sum_{\{n\}} e^{-Ht}, \quad Z_i = \sum_{\{n\}} n_i e^{-Ht}, \quad t = \frac{1}{k_B T} \quad (9)$$

From equations (8) and (9) and (5),

$$\begin{aligned} \langle n_i \rangle &= \langle (1 - n_i) e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t} \rangle \\ \langle n_i \rangle + \langle n_i e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t} \rangle &= \langle e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t} \rangle \\ \langle n_i \left(1 + e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t} \right) \rangle &= \langle e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t} \rangle \\ \langle n_i \rangle &= \left\langle \frac{e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t}}{1 + e^{\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{\varepsilon_{i,j} n_j t}} \right\rangle \\ \langle n_i \rangle &= \left\langle \frac{1}{1 + e^{-\varepsilon_{i,i} t} \prod_{j=1 \neq i}^N e^{-\varepsilon_{i,j} n_j t}} \right\rangle \\ \langle n_i \rangle &= \left\langle \frac{1}{1 + e^{-\tilde{\xi}_i t}} \right\rangle \quad (10) \end{aligned}$$

$$\xi_i = \varepsilon_{i,i} + \sum_{j=1 \neq i}^N \varepsilon_{i,j} n_j, \quad t = \frac{1}{k_B T} \quad (11)$$

Based on the above, equations (1) and (2) are obtained. Equation (10) (equivalent to equation (1)) is a sigmoid function.

Furthermore, at $t \rightarrow \infty$, From $\langle n \rangle = 0$ or 1

$$\langle n_a n_b \rangle = \langle n_a \rangle \langle n_b \rangle$$

From the above, using the relations in equation (10)'s expansion, equations (3) and (4) are obtained.

$$\begin{aligned} \left\langle \frac{1}{1 + e^{-\xi_i t}} \right\rangle &= \frac{1}{2} + \frac{1}{2} \langle \tanh \left(\frac{1}{2} \xi_i t \right) \rangle \\ &= \frac{1}{2} + \frac{1}{2} \left\{ \left\langle \left(\frac{1}{2} \xi_i t \right) \right\rangle - \frac{1}{3} \left\langle \left(\frac{1}{2} \xi_i t \right)^3 \right\rangle + \frac{2}{15} \left\langle \left(\frac{1}{2} \xi_i t \right)^5 \right\rangle - \dots \right\} \\ &= \frac{1}{2} + \frac{1}{2} \left\{ \left\langle \left(\frac{1}{2} \tilde{\xi}_i t \right) \right\rangle - \frac{1}{3} \left\langle \left(\frac{1}{2} \tilde{\xi}_i t \right)^3 \right\rangle + \frac{2}{15} \left\langle \left(\frac{1}{2} \tilde{\xi}_i t \right)^5 \right\rangle - \dots \right\} \\ &= \frac{1}{2} + \frac{1}{2} \tanh \left(\frac{1}{2} \tilde{\xi}_i t \right) \\ &= \frac{1}{1 + e^{-\tilde{\xi}_i t}} \\ \langle n_i \rangle &= \frac{1}{1 + e^{-\tilde{\xi}_i t}} \quad (3) \end{aligned}$$

$$\tilde{\xi}_i = \varepsilon_{i,i} + \sum_{j=1 \neq i}^N \varepsilon_{i,j} \langle n_j \rangle, \quad t = \frac{1}{k_B T} \quad (4)$$

This constitutes a set of nonlinear simultaneous equations for $\langle n_i \rangle$. In other words, the problem of finding the order parameters $\langle n_i \rangle$ of the ground state reduces to solving these nonlinear simultaneous equations.

4 DISCUSSION

In mathematics, there exists an unresolved problem known as the P vs NP problem. This problem, in the field of computer science, raises the question of whether if a solution to a problem can be efficiently verified, can it also be efficiently found? Here, "efficiently" means that the solution can be computed in polynomial time with respect to the size of the problem. "Polynomial time" refers to an algorithm that takes time proportional to n raised to a constant power (where n is the size of the problem). In essence, the P vs NP problem asks whether all problems for which solutions can be quickly verified (NP) are also problems for which solutions can be quickly found (P). Now, about this problem, the ground state of the spin-glass Ising model has been reduced to a solution of a set of nonlinear simultaneous equations. Since finding the ground state of the spin-glass Ising model is equivalent to NP-

complete problems, which can be converted to all NP problems, if the solution to these nonlinear simultaneous equations can be obtained through numerical computation, it can be said that $P=NP$ within a practical range.

REFERENCES

- [1] Hidetoshi Nishimori, Iwanami Shoten, New Physics Selection, "Spin Glass Theory and Information Statistical Mechanics", February 10, 2016.
- [2] Steven H. Strogatz, "Nonlinear Dynamics And Chaos", January 30, 2015.