

# Axiomatic Particle Theory

## Fermions and Gravitation

Author Peter Hickman  
January 14, 2021

### **Abstract**

An axiomatic proposal for an underlying description of particles and their interactions. The existence of fundamental laws of physics is precluded and only random events exist at the fundamental level. Quarks and Leptons in 3 families are found along with spin 2 massless gravitation. Chiral sector - 3 chiral neutrinos can mix, no charged lepton mixing, quark mixing is allowed. The standard model groups  $U(1)$ ,  $SU(2)$  and  $SU(3)$  act on vacuum states to generate particles.

# 1 Introduction

An axiomatic proposal for an underlying description of particles and their interactions. The canonical approach in physics is to deduce equations of motion for the state of a system. The state is not deduced i.e particle content is assumed along with the electrical charges etc. However this approach does not provide a mechanism of how the state changed. As Feynmann noted [1], a particle does not calculate its trajectory in phase space. In addition the reason for the particle content is not addressed. In the Standard Model (SM) [2] of particle physics, the particle content and its interactions are assumed. The fermion content will be addressed in this paper.

In the construction of physics models / theories, assumptions about which mathematical structures are to be used are made. This freedom to choose the mathematical structure is not accessible by Nature. Since Axiomatic Set Theory is a foundation of Mathematics, this will be the foundation for an axiomatic model of nature at the fundamental level.

## 2 Axioms

**A1:** No equations of motion for a fundamental system

**A2:** preclude the creation of same element in succession

Consider a set  $S$  of elements  $\epsilon$  which represents the state of a system

**A1**  $\Rightarrow$  no specified relational measure on  $S$

**A1**  $\Rightarrow$  no functional relations on  $\epsilon$

**A1**  $\Rightarrow$   $\epsilon$  are independent

**A1**  $\Rightarrow$   $\epsilon$  are random

independent implies  $\Rightarrow$   $\epsilon$  are orthogonal

**A2**  $\Rightarrow$   $\epsilon$  form a CAR algebra

**A3:** minimum 2d CAR,  $\epsilon \rightarrow \epsilon_i : i = 1, 2$

**A3** requires inclusion of all 2d vector spaces  $V_k$

Basis  $e_{ki} : e_{ki}^2 = \pm 1$

The basis  $e_{ki} : e_{ki}^2 = -1$  act like imaginary units

$signature(V_k) \in \{(1, 1), (1, -1), (-1, -1), (-1, 1)\}$

Let  $V = \cup_k V_k$  is the set of all 2d spaces

$\epsilon_k$  are the 2d CAR elements on  $V_k$

Othorgonality of  $\epsilon_k \Rightarrow$  the 2d vector spaces  $V_k$  are orthogonal

All sets form a state  $S$  of the system

On  $V_k$  the state is  $(\epsilon_{k1}, 0)$  or  $(0, \epsilon_{k2})$

### 3 Real and Complex States

The real state  $\{(e_{11}, 0), (0, e_{22}), (e_{31}, 0), (e_{41}, 0)\} \in \mathbb{R}^{1,3}$

The real state  $\{(e_{11}, 0), (e_{21}, 0), (e_{31}, 0), (0, e_{42})\} \in \mathbb{R}^{3,1}$

These real states are not to be identified as space-time since no co-ordinate chart has been specified

The pair  $(e_{11}, e_{22}) \in \mathbb{C}$  and similarly other complex states can be found

From [2.19], the states  $V_2, V_4 \notin \mathbb{C}$

### 4 Vacuum States

As particles not defined yet the following  $\mathbb{C}^2$  states are considered to be Vacuum states

$\{(V_1, V_2), (V_3, V_4)\}$

$\{(V_1, V_3), (V_2, V_4)\}$

$\{(V_1, V_4), (V_2, V_3)\}$

as shown below the actions of  $U(1)$ ,  $SU(2)$  and  $SU(3)$  acts on the vacuum states to generate particles

### 5 Chirality

Adapting the definition of chirality [3]

Chirality  $c = -\sigma_0\sigma_1\sigma_2\sigma_3, \sigma_i^2 = -1$

On  $\mathbb{R}^{1,3}$ ,  $\sigma_0^2 = 1, c = -1$

On  $\mathbb{R}^{3,1}$ ,  $\sigma_0^2 = -1, c = -1$

as shown below Chiral  $-1$  fermions

### 6 Scalars

Representations  $e^{in\theta} \in U(1)$  can all be constructed from the fundamental representations  $n = 0, \pm 1$

Let  $\phi = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \in \mathbb{C}^2$

Action of  $U(1)$  on  $\phi$

$$\begin{pmatrix} \phi^n \\ \phi^0 \end{pmatrix} = U(1) \times \phi \rightarrow \begin{pmatrix} \phi_1 e^{in\theta} \\ \phi_2 \end{pmatrix}$$

Complex scalar doublets  $n = \pm 1 - \begin{pmatrix} \phi^\mp \\ \phi^0 \end{pmatrix}$  Higgs doublets when multiplied by dimensionful constant.

## 7 Chiral -1 Fermions

Rotation of elements of  $V_2$  or  $V_4$  induces  $\mathbb{C}^2 \leftrightarrow \mathbb{R}^{1,3}$   
The entangled state is

$$\psi_2 = \begin{pmatrix} (V_1, V_2) \\ (V_1, V_3) \\ (V_1, V_4) \end{pmatrix}$$

There are 3  $\mathbb{C}^2$ , let  $\psi = \begin{pmatrix} \psi_{1i} \\ \psi_2 \end{pmatrix} \in \mathbb{C}^2$   
For  $\psi_2 \notin \mathbb{C}^3$ , action of  $U(1)$  on  $\psi$

$$\begin{pmatrix} \psi_i^n \\ \psi_2 \end{pmatrix} = U(1) \times \psi \rightarrow \begin{pmatrix} \psi_{1i} e^{in\theta} \\ \psi_2 \end{pmatrix}$$

Complex vectors  $n = \pm 1 - \begin{pmatrix} \psi_i^\mp \\ \psi_0 \end{pmatrix}$  - 3 chiral lepton doublets  
For  $\psi_2 \in \mathbb{C}^3$ ,  $\psi_2 \rightarrow \psi_{2j}$ , action of  $U(1)$  on  $\psi$

$$\begin{pmatrix} \psi_i^n \\ \psi_j^m \end{pmatrix} = U(1) \times \psi \rightarrow \begin{pmatrix} \psi_{1i} e^{in\theta} \\ \psi_{2j} e^{im\theta} \end{pmatrix}$$

entangled state  $\psi_{2j}$  is a  $U(1)$  state  $\therefore \sum_j n_j = \pm 1$ , therefore the charge on each state  $\psi_{2j}$  is  $\pm \frac{1}{3}$ , the minimum non-zero difference  $n - m = -1$ , (The charge difference must be a fundamental representation of  $U(1)$ , thus no further charges), the charges on the quarks are  $\mp \frac{1}{3}, \pm \frac{2}{3}$

Note the entangled state is a triplet of quarks  
3 chiral quark doublets

## 8 Chiral +1 Fermions

Let  $\phi = \begin{pmatrix} e_1, e_2 \\ e_3, e_4 \end{pmatrix} \in \mathbb{C}^2$  where  $e_1^2 = e_3^2 = 1, e_2^2 = e_4^2 = -1$

In addition  $\begin{pmatrix} e_1, e_4 \\ e_3, e_2 \end{pmatrix} \in \mathbb{C}^2$  thus the following action on the state  $\phi$

$$su(2, \mathbb{C}) \left( \begin{pmatrix} e_1, e_2 \\ e_3, e_4 \end{pmatrix} \right) + su(2, \mathbb{C}) \left( \begin{pmatrix} e_1, e_4 \\ e_3, e_2 \end{pmatrix} \right)$$

The chiral state is U(1) invariant to the 2nd  $su(2, \mathbb{C})$  thus for example generate electron chiral +1 as below

$$\begin{pmatrix} e_L^- \\ \bar{\nu}_{eR} \end{pmatrix} + \begin{pmatrix} e_L^+ \\ e_R^- \end{pmatrix}_{-1}, \begin{pmatrix} e_L^- \\ \bar{\nu}_{eR} \end{pmatrix} + \begin{pmatrix} e_L^- \\ e_R^+ \end{pmatrix}_{+1}$$

where the  $\pm 1$  indicates the chirality

## 9 Gravitation spin 2

The  $su(2, \mathbb{C})$  acts on the components and or basis.

Thus 2  $su(2, \mathbb{C})$  act on the basis and 2  $su(2, \mathbb{C})$  act on the components each action generates  $s = \pm \frac{1}{2}$

$$(su(2, \mathbb{C}) \times su(2, \mathbb{C}) \times su(2, \mathbb{C}) \times su(2, \mathbb{C}))\mathbb{C}^2 \rightarrow \begin{pmatrix} +2 \\ -2 \end{pmatrix}$$

## 10 A Probability Distribution

$(c, \theta)$  for complex number  $ce^{i\theta}$

$(c, \theta)$  are random - construct a partition function

$$Z = \iint dcd\theta e^{-(c^2 + \theta^2)}$$

The partition function favours complex states with small amplitude and phase  $\theta$  For unitary transformations of states, (changes to the phase  $\theta$ ) small  $\theta \implies$  also Lie transformations of states

## 11 Conclusion

From minimal set of axioms, find 3 families of leptons and quarks along with chiral sector - 3 chiral neutrinos can mix, no charged lepton mixing, quark mixing is allowed. The standard model groups U(1), SU(2) and SU(3) act on vacuum states to generate particles. The states can be real or complex states. Future work involves understanding the boson sector of the Standard Model.

## References

- [1] Richard Feynman's Lectures on Physics (The Messenger Lectures)  
<https://archive.org> - last accessed Jan 2021

- [2] Pham Xuan Yem Quang Ho-Kim. Elementary Particles and Their Interactions, Springer-Verlag, 1998
- [3] Pham Xuan Yem Quang Ho-Kim. Elementary Particles and Their Interactions, 3 Fermions p84, Springer-Verlag, 1998