

Klein-Gordon Equation and Wave Function in Cosmological Special Theory of Relativity

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ABSTRACT

In the Cosmological Special Theory of Relativity, we study energy-momentum relations, Klein-Gordon equation and wave function.

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1. Introduction

Our article's aim is that we make Klein-Gordon equation and wave function in cosmological special theory of relativity.

At first, space-time relations are in cosmological special theory of relativity (CSTR).[7]

$$\begin{aligned}
 ct &= \gamma \left(ct + \frac{v_0}{c} \Omega(t) x \right), \quad x \Omega(t_0) = \gamma (\Omega(t_0) x' + v_0 \Omega(t_0) t') \\
 \Omega(t_0) y &= \Omega(t_0) y', \\
 \Omega(t_0) z &= \Omega(t_0) z', \quad \gamma = 1 / \sqrt{1 - \frac{v_0^2}{c^2} \Omega^2(t_0)}, \quad t_0 \text{ is cosmological time}
 \end{aligned} \quad (1)$$

Therefore, proper time is [7]

$$\begin{aligned}
 d\tau^2 &= dt^2 - \frac{1}{c^2} \Omega^2(t_0) [dx^2 + dy^2 + dz^2] \\
 &= dt'^2 - \frac{1}{c^2} \Omega^2(t') [dx'^2 + dy'^2 + dz'^2], \quad t_0 \text{ is cosmological time}
 \end{aligned} \quad (2)$$

Hence, energy-momentum relations are by the fact that energy-momentum are 4-vector in CSTR,

$$\begin{aligned}
 E &= \gamma (E' + v_0 \Omega^2(t_0) p_x'), \quad p_x \Omega(t_0) = \gamma (\Omega(t_0) p_x' + \frac{v_0}{c^2} \Omega(t_0) E') \\
 \Omega(t_0) p_y &= \Omega(t_0) p_y', \\
 \Omega(t_0) p_z &= \Omega(t_0) p_z', \quad \gamma = 1 / \sqrt{1 - \frac{v_0^2}{c^2} \Omega^2(t_0)}, \quad E = m_0 c^2 \frac{dt}{d\tau}, \quad \vec{p} = m_0 \frac{d\vec{x}}{d\tau}
 \end{aligned} \quad (3)$$

Therefore, energy-momentum-mass relation is in CSTR,

$$m_0^2 c^4 = E^2 - \Omega^2(t_0) p^2 c^2 \quad (4)$$

2. Klein-Gordon Equation and Wave Function in CSTR

According to [7], matter wave function is in CSTR,

$$\begin{aligned}
 \phi &= \phi_0 \exp i\Phi = \phi_0 \exp i \left[\frac{\omega t}{\sqrt{\Omega(t_0)}} - \vec{k} \cdot \vec{x} \sqrt{\Omega(t_0)} \right] \\
 &= \phi_0' \exp i\Phi' = \phi_0' \exp i \left[\frac{\omega' t'}{\sqrt{\Omega(t_0)}} - \vec{k}' \cdot \vec{x}' \sqrt{\Omega(t_0)} \right] \\
 \phi_0 &\text{ is amplitude, } \omega \text{ is angular frequency, } k = |\vec{k}| \text{ is wave number.}
 \end{aligned} \quad (5)$$

If we use Eq(1) in Eq(5), we obtain angular frequency-wave number relation.

$$\omega' = \gamma (\omega - v_0 \Omega(t_0) k_1), \quad k_1' = \gamma \left(k_1 - \frac{v_0}{c^2} \Omega(t_0) \omega \right)$$

$$k_2' = k_2, k_3' = k_3, \gamma = 1 / \sqrt{1 - \frac{V_0^2}{c^2} \Omega^2(t_0)} \quad (6)$$

In this time, if we define energy-momentum by angular frequency-wave number,

$$E = \hbar \omega, \vec{p} = \frac{\hbar \vec{k}}{\Omega(t_0)} \quad (7)$$

Hence, we obtain the angular frequency-wave number relation about the energy-momentum-mass relation in CSTR,

$$m_0^2 c^4 = E^2 - \Omega^2(t_0) p^2 c^2 = \hbar^2 \omega^2 - \hbar^2 k^2 c^2 \quad (8)$$

We obtain next result by the transformation of the angular frequency-wave number relation, Eq(6) in CSTR.

$$m_0^2 c^4 = \hbar^2 \omega^2 - \hbar^2 k^2 c^2 = \hbar^2 \omega'^2 - \hbar^2 k'^2 c^2 \quad (9)$$

If we define the differential operator about energy-momentum in CSTR,

$$E = i\hbar \sqrt{\Omega(t_0)} \frac{\partial}{\partial t}, \vec{p} = -i\hbar \frac{1}{\Omega(t_0) \sqrt{\Omega(t_0)}} \vec{\nabla} \quad (10)$$

If we apply Eq(10) to Eq(4),

$$m_0^2 c^4 = E^2 - \Omega^2(t_0) p^2 c^2 = \hbar^2 [-\Omega(t_0) (\frac{\partial}{\partial t})^2 + \frac{1}{\Omega(t_0)} c^2 \nabla^2]$$

We finally obtain Klein-Gordon equation in CSTR.

$$\frac{m_0^2 c^2}{\hbar^2} \phi = [-\Omega(t_0) \frac{1}{c^2} (\frac{\partial}{\partial t})^2 + \frac{1}{\Omega(t_0)} \nabla^2] \phi \quad (11)$$

Wave function, Eq(5) satisfy Klein-Gordon equation, Eq(11) in CSTR.

3. Conclusion

We are able to describe free particle by Klein-Gordon equation and wave function in CSTR.

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