

Yukawa Potential and Extended Klein-Gordon Equation in Rindler Space-Time

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ABSTRACT

Yukawa potential satisfy Proca equation or Klein-Gordon equation. If we represent Yukawa potential in Rindler space-time, this Yukawa potential satisfy the extended Klein-Gordon equation in Rindler space-time. We understand Yukawa force in Rindler space-time.

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1. Introduction

Atom's nucleus force understand by Yukawa potential. We study Yukawa potential in Rindler Space-time time.

At first, Yukawa potential V describes nucleus's combine force in semi-classical method.

$$V = -\frac{g^2}{r} \exp\left(-\frac{m_\pi r c}{\hbar}\right)$$

$$g \text{ is real number, } m_\pi \text{ is the meson's mass} \quad (1)$$

Klein-Gordon equation is satisfied by Yukawa potential V .

$$-\partial_i \partial^i V + \frac{m^2 c^2}{\hbar^2} V = -\nabla^2 V + \frac{m_\pi^2 c^2}{\hbar^2} V = 0$$

$$V = -\frac{g^2}{r} \exp\left(-\frac{m_\pi r c}{\hbar}\right), i = 1, 2, 3 \quad (2)$$

2. Yukawa potential from Extended Klein-Gordon Equation in Rindler-Space-Time

Rindler coordinates are

$$ct = \left(\frac{c^2}{a_0} + \xi^1\right) \sinh\left(\frac{a_0}{c} \xi^0\right), \quad x = \left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0}{c} \xi^0\right) - \frac{c^2}{a_0}$$

$$y = \xi^2, z = \xi^3 \quad (3)$$

If we write Yukawa potential V in inertial frame,

$$V = \frac{g^2}{r} \exp\left(\frac{m_\pi r c}{\hbar}\right) \quad (4)$$

If we rewrite Yukawa potential V_ξ in Rindler space-time,

$$V_\xi = -\frac{g^2}{\sqrt{x^2 + y^2 + z^2}} \exp\left(-\frac{m_\pi c}{\hbar} \sqrt{x^2 + y^2 + z^2}\right)$$

$$= -\frac{g^2}{\sqrt{\left\{\left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0}{c} \xi^0\right) - \frac{c^2}{a_0}\right\}^2 + (\xi^2)^2 + (\xi^3)^2}} \exp\left[-\frac{m_\pi c}{\hbar} \sqrt{\left\{\left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0}{c} \xi^0\right) - \frac{c^2}{a_0}\right\}^2 + (\xi^2)^2 + (\xi^3)^2}\right] \quad (5)$$

This Yukawa potential satisfy the extended Klein-Gordon equation. At first, energy and momentum are in Rindler space-time[1],

$$E_\xi = i\hbar \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)} \frac{\partial}{\partial \xi^0}, \quad \vec{p}_\xi = -i\hbar \vec{\nabla}_\xi \quad (6)$$

Energy-Momentum equation is in Rindler space-time[1],

$$E_\xi^2 = \vec{p}_\xi \cdot \vec{p}_\xi c + m^2 c^4 \quad (7)$$

Hence, normal Klein-Gordon equation is in Rindler-spacetime,

$$\begin{aligned} & \frac{m^2 c^2}{\hbar^2} V_\xi + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_\xi - \nabla_\xi^2 V_\xi \\ &= \frac{m_\pi^2 c^2}{\hbar^2} V_\xi + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_\xi - \nabla_\xi^2 V_\xi = 0 \end{aligned} \quad (8)$$

In this time, we focus the gauge Λ equation in Rindler space-time[1],

$$\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \frac{\partial^2}{(\partial \xi^0)^2} \Lambda - \nabla_\xi^2 \Lambda - \frac{\partial \Lambda}{\partial \xi^1} \frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} = 0 \quad (9)$$

Hence, Eq(8) change extended Klein-Gordon equation in Rindler space-time.

Extended Klein-Gordon Equation is in Rindler space-time,

$$\begin{aligned} & \frac{m^2 c^2}{\hbar^2} V_\xi + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_\xi - \nabla_\xi^2 V_\xi - \frac{\partial V_\xi}{\partial \xi^1} \frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \\ &= \frac{m_\pi^2 c^2}{\hbar^2} V_\xi + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \frac{\partial^2}{(\partial \xi^0)^2} V_\xi - \nabla_\xi^2 V_\xi - \frac{\partial V_\xi}{\partial \xi^1} \frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} = 0 \end{aligned} \quad (10)$$

Eq(5), Yukawa potential V_ξ satisfy Eq(10), extended Klein-Gordon equation in Rindler space-time.

Yukawa force \vec{f} is

$$\vec{f} = -\vec{\nabla} V = -\frac{g^2}{r^3} [\exp(-\frac{m_\pi r c}{\hbar})] (1 + \frac{m_\pi r c}{\hbar}) \vec{r} \quad (11)$$

In this time, Yukawa force \vec{f}_ξ is Rindler space-time,

$$\vec{f}_\xi = -\vec{\nabla}_\xi V_\xi = -\frac{g^2}{r^3} [\exp(-\frac{m_\pi r c}{\hbar})] (1 + \frac{m_\pi r c}{\hbar}) (x \cosh(\frac{a_0 \xi^0}{c}), \xi^2, \xi^3) \quad (12)$$

Hence, according to Yukawa force \vec{f}_ξ in Rindler space-time, the nuclear force strongly act in accelerated frame rather than inertial frame in x-axis.

3. Conclusion

We found Yukawa potential mechanism in Rindler Space-time. We understand nuclear force in Rindler space-time.

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