

# **Reissner-Nodstrom Solution and Energy-Momentum Density's Conservation Law**

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## **ABSTRACT**

We find Reissner-Nodstrom solution hold by the energy-momentum density's conservation law (Noether's theorem) of electromagnetic field in general relativity theory.

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**Key words:**General relativity theory;

**Reissner-Nodstrom solution;**

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## 1. Introduction

Our article's aim is that we find Reissner-Nordstrom solution hold by the energy-momentum density's conservation law (Noether's theorem) of electromagnetic field.

In general relativity theory, the energy-momentum tensor  $\mathcal{T}^{\mu\nu}$  of the electromagnetic field is

$$\mathcal{T}^{\mu\nu} = \frac{1}{4\pi c} (\mathcal{F}^{\mu\rho} \mathcal{F}^{\nu\rho} - \frac{1}{4} g^{\mu\nu} \mathcal{F}_{\rho\sigma} \mathcal{F}^{\rho\sigma}) \quad (1)$$

In this time, spherical coordinates is in general relativity theory

$$ds^2 = -A(t, r) dt^2 + (B(t, r))^2 dr^2 + r^2 \sin^2 \theta d\phi^2 \quad (2)$$

Hence, Faraday tensors  $\mathcal{F}^{\mu\nu}, \mathcal{F}_{\mu\nu}$  are in general relativity theory.

$$\mathcal{F}^{\mu\nu} = \begin{pmatrix} 0 & -E(r, t) & 0 & 0 \\ E(r, t) & 0 & 0 & 0 \\ 0 & 0 & 0 & -B(t, r) \\ 0 & 0 & B(t, r) & 0 \end{pmatrix} \quad (3-i)$$

$$\mathcal{F}_{\mu\nu} = \begin{pmatrix} 0 & E(t, r) & 0 & 0 \\ -E(t, r) & 0 & 0 & 0 \\ 0 & 0 & 0 & -B(t, r) \\ 0 & 0 & B(t, r) & 0 \end{pmatrix} \quad (3-ii)$$

Hence, Einstein-Maxwell equations is

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = -\frac{8\pi G}{c^4} \mathcal{T}_{\mu\nu} = -\frac{2G}{c^5} (\mathcal{F}_{\mu\rho} \mathcal{F}^{\nu\rho} - \frac{1}{4} g_{\mu\nu} \mathcal{F}_{\rho\sigma} \mathcal{F}^{\rho\sigma}) \quad (4)$$

$$\mathcal{F}^{\mu\nu}{}_{;\nu} = 0 \quad (5)$$

$$\frac{\partial \mathcal{F}_{\mu\nu}}{\partial x^\rho} + \frac{\partial \mathcal{F}_{\nu\rho}}{\partial x^\mu} + \frac{\partial \mathcal{F}_{\rho\mu}}{\partial x^\nu} = 0 \quad (6)$$

## 2. Reissner-Nordstrom solution and Noether's theorem

In this time, energy-momentum tensor  $\mathcal{T}^{\mu\nu}$  is

$$\begin{aligned} \mathcal{T}^{\mu\nu} &= \frac{1}{4\pi c} (\mathcal{F}^{\mu\rho} \mathcal{F}^{\nu\rho} - \frac{1}{4} g^{\mu\nu} \mathcal{F}_{\rho\sigma} \mathcal{F}^{\rho\sigma}) \\ &= \frac{1}{4\pi c} (g^{\mu\nu} \mathcal{F}_{\nu\rho} \mathcal{F}^{\nu\rho} - \frac{1}{4} g^{\mu\nu} \mathcal{F}_{\rho\sigma} \mathcal{F}^{\rho\sigma}) \end{aligned} \quad (7)$$

Hence,  $\mathcal{T}^{\mu\nu}$  is

$$\begin{aligned}
T^{\mu\nu} &= \frac{E^2 + B^2}{8\pi c} \begin{pmatrix} -g^{00} & 0 & 0 & 0 \\ 0 & -g^{11} & 0 & 0 \\ 0 & 0 & g^{22} & 0 \\ 0 & 0 & 0 & g^{33} \end{pmatrix} \\
&= \frac{E^2 + B^2}{8\pi c} \begin{pmatrix} 1/A & 0 & 0 & 0 \\ 0 & -1/B & 0 & 0 \\ 0 & 0 & 1/r^2 & 0 \\ 0 & 0 & 0 & 1/r^2 \sin^2 \theta \end{pmatrix}
\end{aligned} \tag{8}$$

Therefore, energy-momentum tensor  $T_{\mu\nu}$  is

$$\begin{aligned}
T_{\mu\nu} &= \frac{1}{4\pi c} (F_{\mu\rho} F_{\nu}^{\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma}) \\
&= \frac{1}{4\pi c} (g_{\mu\nu} F_{\nu\rho} F^{\nu\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma})
\end{aligned} \tag{9}$$

Hence, if we calculate  $T_{\mu\nu}$ ,

$$T_{\mu\nu} = \frac{E^2 + B^2}{8\pi c} \begin{pmatrix} A & 0 & 0 & 0 \\ 0 & -B & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{pmatrix} \tag{10}$$

Eq(5) is

$$\frac{\partial F^{\mu\nu}}{\partial x^\nu} + \Gamma^\mu_{\rho\nu} F^{\rho\nu} + \Gamma^\nu_{\rho\nu} F^{\mu\rho} = \frac{\partial F^{\mu\nu}}{\partial x^\nu} + \Gamma^\nu_{\rho\nu} F^{\mu\rho} \tag{11}$$

In this time,

$$\Gamma^\nu_{\rho\nu} = \frac{1}{\sqrt{-g}} \frac{\partial}{\partial x^\rho} \sqrt{-g} \tag{12}$$

Eq(5) is

$$F^{\mu\nu}_{;\nu} = \frac{1}{\sqrt{-g}} \frac{\partial}{\partial x^\nu} (\sqrt{-g} F^{\mu\nu}) = 0 \tag{13}$$

Hence, Eq(13) is

$$\frac{\partial}{\partial r} (r^2 E) = \zeta, \quad E = k \frac{Q}{r^2} \tag{14}$$

Eq(6) is

$$\frac{\partial F_{23}}{\partial r} + \frac{\partial F_{31}}{\partial x^2} + \frac{\partial F_{12}}{\partial x^3} - \frac{\partial B}{\partial r} = 0, B = 0 \quad (15)$$

Hence, we know the following formula.

$$E = k \frac{Q}{r^2}, B = 0 \quad (16)$$

The solution of Eq(4),Eq(5),Eq(6) is Reissner-Nodstrom solution. Hence,

$$A = 1 - \frac{C_1}{r} + \frac{C_2}{r^2}, B = 1 / (1 - \frac{C_1}{r} + \frac{C_2}{r^2}) \quad (17)$$

In this time, if energy-momentum tensor  $T^{\mu\nu}$  satisfy the energy-momentum conservation law (Noether's theorem)

$$\begin{aligned} T^{\mu\nu}_{;\nu} &= T^{00}_{;\nu} + T^{ii}_{;\nu} = T^{00}_{;\nu} + \frac{\partial T^{ii}}{\partial x^\nu} + 2\Gamma^i_{ij}T^{ji} \\ &= \frac{E^2 + B^2}{8\pi c} \frac{\partial(-g^{11})}{\partial r} + 2 \cdot \frac{1}{2} g^{11} \left( \frac{\partial g_{11}}{\partial r} \right) \cdot \frac{E^2 + B^2}{8\pi c} \cdot -g_{11} \\ &= \frac{E^2 + B^2}{8\pi c} \left( -\frac{C_1}{r^2} + \frac{2C_2}{r^3} \right) + \frac{E^2 + B^2}{8\pi c} \left( \frac{C_1}{r^2} - \frac{2C_2}{r^3} \right) = 0 \end{aligned} \quad (18)$$

### 3. Conclusion

Reissner-Nodstrom solution hold by the energy-momentum density's conservation law of electromagnetic field in general relativity theory.

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