

# Quark string of elementary particle

T.Spiegelman

tnts2595@gmail.com, south korea, ulsan, dongsan-ro-78

July 19, 2018

## Abstract

This paper suggests quark combination of elementary particle based on AdS/CFT correspondence[1]. Through this, we can define quark conservation law and majonara particle. Tension of closed string which diverge to infinity confine quarks as in color confinement of strong interaction.

## 1 Structure of quark string

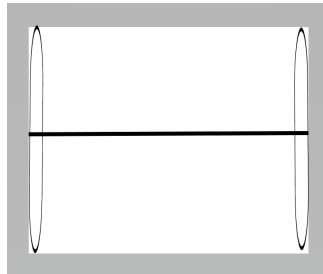


Figure 1: Non-vibrating state of open string is straight line with two D-brane[2]

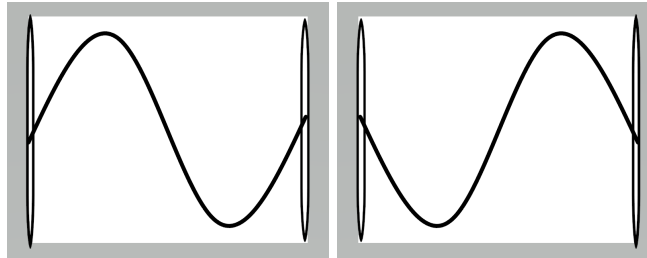
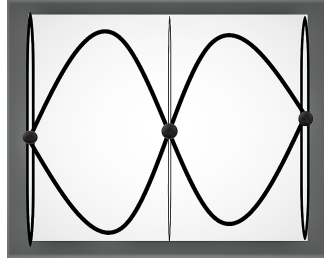
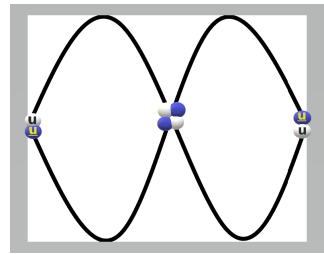


Figure 2: Vibrating state of open string

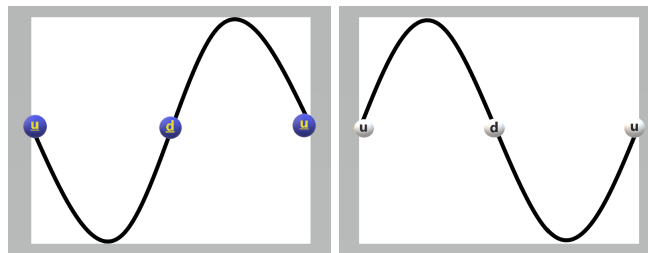
Two open strings interact and become a closed string.



A closed string disintegrate from D-brane and generate quark on lattice. State of  $q\bar{q}$  pair of open string generate four quarks of two  $q\bar{q}$  pairs on edges of a closed string. By interact of strings, quark of different style are generated in center lattice.

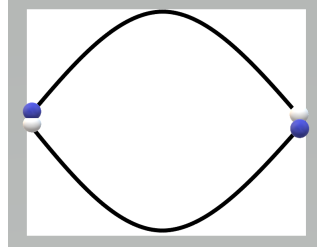


Disintegrated a closed string disintegrate itself to  $L$  and  $\bar{L}$ . For example, quark string of electron and positron are shown to below picture.



But in reality, it will be changed into close string.

An open string of two quarks; gauge boson is made as below picture.



Tension of string which confine quarks will be introduced at chapter 4.

## 2 Quark combination of elementary particle

A photon is a string of two edges. In pair production,  $0 \longleftrightarrow q\bar{q}$  where quarks are generated from lattices of a closed string.

Lepton consists of three quarks (total spin :  $\pm\frac{1}{2}$ ) is given by

$$L_e \begin{cases} e^- * \overline{uud} \\ e^+ * uud \\ 0(\nu) * \overline{udd} \\ 0(\bar{\nu}) * udd \end{cases}$$

These quark combinations are make possible of quark conservation law in particle disintegration. It replaces quark style conversion ability of W boson with quark conservation law by added quark combinations of W boson and leptons.

Define the quark conservation box as  $\begin{pmatrix} u & d \\ \bar{u} & \bar{d} \end{pmatrix}$ .

In case of  $\beta^-$  decay[3] :

$$n \longrightarrow p + \bar{\nu} + e^-$$

where the quark conservation box is  $\begin{pmatrix} u & d \\ \bar{u} & \bar{d} \end{pmatrix}$ , quark conservation is given by

$$n \begin{pmatrix} 1 & 2 \\ 0 & 0 \end{pmatrix} \longrightarrow p + \bar{\nu} + e^- \begin{pmatrix} 3 & 3 \\ 2 & 1 \end{pmatrix},$$

and  $\begin{pmatrix} 3 & 3 \\ 2 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 2 \\ 0 & 0 \end{pmatrix}$  because of  $0 \leftrightarrow q\bar{q}$

In case of pion decay[4],

$$\pi^- \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \longrightarrow \mu^- + \bar{\nu}_\mu \begin{pmatrix} 1 & 2 \\ 2 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\pi^+ \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \longrightarrow \mu^+ + \nu_\mu \begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\pi^0 \begin{pmatrix} 1 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} [5] \longrightarrow \gamma + \gamma \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} = \gamma' \begin{pmatrix} 1 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix}$$

Also we get quark system of W boson,

$$W_e \begin{cases} e^- * \bar{u}d \\ e^+ * u\bar{d} \end{cases}$$

$$W^- \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \longrightarrow e^- + \bar{\nu}_e \begin{pmatrix} 1 & 2 \\ 2 & 1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$W^+ \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \longrightarrow e^+ + \nu_e \begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

### 3 Majonara particle

Antiparticle is the structure of antiquarks which are symmetry to particle. It gives quark condition of majonara particle

$$AB = \overline{AB}$$

where  $B$  is  $\bar{A}$ . Thus, For the two quarks system, majonara particle is given by

$$A\bar{A} = \overline{AA}.$$

It means non-charge gauge bosons are antiparticle themselves.

Possible quark combinations of three quarks fermion are given by

$$[aab, \overline{aab}, aaa, \overline{aaa}, bbb, \overline{bbb}, abb, \overline{abb}]$$

where  $a$  is  $+\frac{2}{3}$  quark and  $b$  is  $-\frac{1}{3}$  quark. In short, fermion cannot be majonara particle because its quarks number is odd, while majonara particle consist of a pair of particle-antiparticle.

## 4 Tension of quark confinement

Describe the motion of a particle scattered on lattice. If we consider only gravitational field, it will be described as projectile motion.

$$x = (v_0 \cos \alpha)t$$

$$y = -\frac{1}{2}gt^2 + (v_0 \sin \alpha)t$$

Writing  $x$  as a function of  $t$ ,

$$t = \frac{x}{v_0 \cos \alpha}.$$

If we insert this value of  $t$  in the equation for  $y$ , we obtain

$$y = -\frac{g}{2v_0^2 \cos^2 \alpha} \left( x - \frac{v_0^2 \sin \alpha \cos \alpha}{g} \right)^2 + \frac{v_0^2 \sin^2 \alpha}{2g},$$

and the vertex is

$$\left( \frac{v_0^2 \sin \alpha \cos \alpha}{g}, \frac{v_0^2 \sin^2 \alpha}{2g} \right).$$

Assume that scattered particle follow the principle of least action[6]. Angle  $\alpha$  of closed string is adjusted to specific value to scatter a particle most far from lattice in most low energy. Most far distance from lattice, in case of distance from lattice to vertex  $r$  is given by

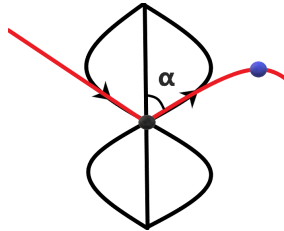
$$r = \frac{v_0^2 \sqrt{\sin 2\alpha + \sin^2 \alpha}}{2g},$$

and  $\alpha$  is  $\frac{\pi}{4}$  for the max value of  $r$ .

Tension of closed string is given by

$$T = \frac{\sin^4 \alpha}{d^2 \cos 2\alpha} [7]$$

where  $d$  is length of closed string.



According to the principle of least action, closed string maintain its angle to  $\frac{\pi}{4}$  until it interact with other strings. Thus,

$$T \rightarrow \infty,$$

closed string confine quarks.

## 5 References

- [1] Maldacena, Juan M. (1998). The large N limit of superconformal field theories and supergravity. *Advances in Theoretical and Mathematical Physics* 2 (1): 231252.
- [2] Johnson, Clifford V. (2003). *D-Branes*. Cambridge Monographs on Mathematical Physics. Cambridge University Press.
- [3] Konya, J.; Nagy, N. M. (2012). *Nuclear and Radio-chemistry*. Elsevier. pp. 7475.
- [4] C. Amsler et al. (2008): Particle listings  $\pi^\pm$
- [5] R. Bjorklund; W. E. Crandall; B. J. Moyer; H. F. York (1950). "High Energy Photons from Proton-Nucleon Collisions". *Physical Review*. 77 (2): 213218.
- [6] Gray, Chris G (2009). Principle of least action. *Scholarpedia* 4 (12): 8291
- [7] Zeyad Al-Zhour, Amin A. Aqel, Abdallah Ibrahim(2008). "Mathematical properties of DNA structure in 3-dimensional space", section 2, p. 7(220)