

# The generalized Bernstein-Vazirani algorithm for determining an integer string

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We present the generalized Bernstein-Vazirani algorithm for determining a restricted integer string. Given the set of real values  $\{a_1, a_2, a_3, \dots, a_N\}$  and a function  $g : \mathbf{R} \rightarrow \mathbf{Z}$ , we shall determine the following values  $\{g(a_1), g(a_2), g(a_3), \dots, g(a_N)\}$  simultaneously. The speed of determining the values is shown to outperform the classical case by a factor of  $N$ . The method determines the maximum of and the minimum of the function  $g$  that the finite domain is  $\{a_1, a_2, a_3, \dots, a_N\}$ .

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## I. INTRODUCTION

In 1993, the Bernstein-Vazirani algorithm was published [1, 2]. This work can be considered an extension of the Deutsch-Jozsa algorithm [3–5]. In 1994, Simon’s algorithm [6] and Shor’s algorithm [7] were discussed. In 1996, Grover [8] provided the highest motivation for exploring the computational possibilities offered by quantum mechanics.

The original Bernstein-Vazirani algorithm [1, 2] determines a bit string. It is extended to determining the values of a function [9, 10]. The values of the functions are restricted in  $\{0, 1\}$ . By using the extension, we can consider quantum algorithm of calculating a multiplication [10].

By extending the Bernstein-Vazirani algorithm more, we give an algorithm of determining the values of a function that are extended to the natural numbers  $\mathbf{N}$  [11]. That is, the extended algorithm determines a natural number string instead of a bit string. So we have the generalized Bernstein-Vazirani algorithm for determining a restricted natural number string. By using the extension, quantum algorithm for determining a homogeneous linear function is studied.

Here, by extending the quantum algorithm more and more, we present an algorithm of determining the values of a function that are extended to the integers  $\mathbf{Z}$ . That is, the extended algorithm determines an integer string instead of a natural number string.

In this article, we present the generalized Bernstein-Vazirani algorithm for determining an integer string. Given the set of real values  $\{a_1, a_2, a_3, \dots, a_N\}$  and a function  $g : \mathbf{R} \rightarrow \mathbf{Z}$ , we shall determine the following values  $\{g(a_1), g(a_2), g(a_3), \dots, g(a_N)\}$  simultaneously. The speed of determining the values is shown to outperform the classical case by a factor of  $N$ . The method determines the maximum of and the minimum of the function

$g$  that the finite domain is  $\{a_1, a_2, a_3, \dots, a_N\}$ . Our argumentations provide a new insight into the importance of the original Bernstein-Vazirani algorithm.

## II. THE QUANTUM ALGORITHM FOR DETERMINING THE MAXIMUM OF AND THE MINIMUM OF A FUNCTION

Let us suppose that the following sequence of real values is given

$$a_1, a_2, a_3, \dots, a_N. \quad (1)$$

Let us now introduce a function

$$g : \mathbf{R} \rightarrow \mathbf{Z}. \quad (2)$$

Our goal is of determining the following values

$$g(a_1), g(a_2), g(a_3), \dots, g(a_N). \quad (3)$$

We can determine the maximum of and the minimum of the function  $g$  that the finite domain is  $\{a_1, a_2, a_3, \dots, a_N\}$ . Recall that in the classical case, we need  $N$  queries, that is,  $N$  separate evaluations of the function (2). In our quantum algorithm, we shall require a single query.

We introduce a positive integer  $d$ . Throughout the discussion, we consider the problem in the modulo  $d$ . Assume the following

$$-(d-1) \leq \overbrace{g(a_1), g(a_2), g(a_3), \dots, g(a_N)}^N \leq d-1 \quad (4)$$

where  $g(a_j) \in \{-(d-1), \dots, -1, 0, 1, \dots, d-1\}$ , and we define

$$g(a) = (g(a_1), g(a_2), g(a_3), \dots, g(a_N)) \quad (5)$$

where each entry of  $g(a)$  is an integer in the modulo  $d$ . Here  $g(a) \in \{-(d-1), \dots, -1, 0, 1, \dots, d-1\}^N$ . We define  $f(x)$  as follows

$$\begin{aligned} f(x) &= g(a) \cdot x \text{ mod } d \\ &= g(a_1)x_1 + g(a_2)x_2 + \dots + g(a_N)x_N \text{ mod } d \end{aligned} \quad (6)$$

where  $x = (x_1, \dots, x_N) \in \{-(d-1), \dots, -1, 0, 1, \dots, d-1\}^N$ . Let us follow the quantum states through the algorithm.

The input state is

$$|\psi_0\rangle = |0\rangle^{\otimes N} |d-1\rangle \quad (7)$$

where  $|0\rangle^{\otimes N}$  means  $\overbrace{|0, 0, \dots, 0\rangle}^N$ . We discuss the general Fourier transform of  $|0\rangle$

$$|0\rangle \rightarrow \sum_{y=-(d-1)}^{d-1} \frac{\omega^{y \cdot 0} |y\rangle}{\sqrt{2d-1}} = \sum_{y=-(d-1)}^{d-1} \frac{|y\rangle}{\sqrt{2d-1}} \quad (8)$$

where we have used  $\omega^0 = 1$ .

Subsequently let us define the wave function  $|\phi\rangle$  as follows

$$|\phi\rangle = \frac{1}{\sqrt{d}} (\omega^d |0\rangle + \omega^{d-1} |1\rangle + \dots + \omega |d-1\rangle) \quad (9)$$

where  $\omega = e^{2\pi i/d}$ . In the following, we discuss the Fourier transform of  $|d-1\rangle$

$$\begin{aligned} |d-1\rangle &\rightarrow \sum_{y=0}^{d-1} \frac{\omega^{y \cdot (d-1)} |y\rangle}{\sqrt{d}} = \sum_{y=0}^{d-1} \frac{\omega^{yd-y} |y\rangle}{\sqrt{d}} \\ &= \sum_{y=0}^{d-1} \frac{\omega^{d-y} |y\rangle}{\sqrt{d}} = |\phi\rangle \end{aligned} \quad (10)$$

where we have used  $\omega^{yd} = \omega^d = 1$ .

The general Fourier transform of  $|x_1 \dots x_N\rangle$  is as follows

$$\begin{aligned} &|x_1 \dots x_N\rangle \\ &\rightarrow \sum_{z_1=-(d-1)}^{d-1} \dots \sum_{z_N=-(d-1)}^{d-1} \frac{\omega^{z_1 x_1} |z_1\rangle}{\sqrt{2d-1}} \dots \frac{\omega^{z_N x_N} |z_N\rangle}{\sqrt{2d-1}} \\ &= \sum_{z \in K} \frac{\omega^{z \cdot x} |z\rangle}{\sqrt{(2d-1)^N}} \end{aligned} \quad (11)$$

where  $K = \{-(d-1), \dots, -1, 0, 1, \dots, d-1\}^N$  and  $z$  is  $(z_1, z_2, \dots, z_N)$ . Hence, for completeness,  $\sum_{z \in K}$  is a shorthand to the compound sum

$$\sum_{z_1 \in \{-(d-1), \dots, -1, 0, 1, \dots, d-1\}} \dots \sum_{z_N \in \{-(d-1), \dots, -1, 0, 1, \dots, d-1\}} \quad (12)$$

After the componentwise general Fourier transforms of the first  $N$  qudits state and after the Fourier transform of  $|d-1\rangle$  in (7)

$$\overbrace{G|0\rangle \otimes G|0\rangle \otimes \dots \otimes G|0\rangle}^N \otimes F|d-1\rangle \quad (13)$$

we have

$$|\psi_1\rangle = \sum_{x \in K} \frac{|x\rangle}{\sqrt{(2d-1)^N}} |\phi\rangle. \quad (14)$$

Here, the notation  $G|0\rangle$  means the general Fourier transform of  $|0\rangle$  and the notation  $F|d-1\rangle$  means the Fourier transform of  $|d-1\rangle$ .

We introduce  $SUM_{f(x)}$  gate

$$|x\rangle |j\rangle \rightarrow |x\rangle |(f(x) + j) \text{ mod } d\rangle \quad (15)$$

where

$$f(x) = g(a) \cdot x \text{ mod } d. \quad (16)$$

We have

$$SUM_{f(x)} |x\rangle |\phi\rangle = \omega^{f(x)} |x\rangle |\phi\rangle. \quad (17)$$

In what follows, we will discuss the rationale behind of the above relation (17). Now consider applying the  $SUM_{f(x)}$  gate to the state  $|x\rangle |\phi\rangle$ . Each term in  $|\phi\rangle$  is of the form  $\omega^{d-j} |j\rangle$ . We see

$$\begin{aligned} &SUM_{f(x)} \omega^{d-j} |x\rangle |j\rangle \\ &\rightarrow \omega^{d-j} |x\rangle |(j + f(x)) \text{ mod } d\rangle. \end{aligned} \quad (18)$$

We introduce  $k$  such as  $f(x) + j = k \Rightarrow d - j = d + f(x) - k$ . Hence (18) becomes

$$\begin{aligned} &SUM_{f(x)} \omega^{d-j} |x\rangle |j\rangle \\ &\rightarrow \omega^{f(x)} \omega^{d-k} |x\rangle |k \text{ mod } d\rangle. \end{aligned} \quad (19)$$

Now, when  $k < d$  we have  $|k \text{ mod } d\rangle = |k\rangle$  and thus, the terms in  $|\phi\rangle$  such that  $k < d$  are transformed as follows

$$SUM_{f(x)} \omega^{d-j} |x\rangle |j\rangle \rightarrow \omega^{f(x)} \omega^{d-k} |x\rangle |k\rangle. \quad (20)$$

Also, as  $f(x)$  and  $j$  are bounded above by  $d-1$ ,  $k$  is strictly less than  $2d$ . Hence, when  $d \leq k < 2d$  we have  $|k \text{ mod } d\rangle = |k-d\rangle$ . Now, we introduce  $m$  such that  $k-d = m$  then we have

$$\begin{aligned} &\omega^{f(x)} \omega^{d-k} |x\rangle |k \text{ mod } d\rangle = \omega^{f(x)} \omega^{-m} |x\rangle |m\rangle \\ &= \omega^{f(x)} \omega^{d-m} |x\rangle |m\rangle. \end{aligned} \quad (21)$$

Hence the terms in  $|\phi\rangle$  such that  $k \geq d$  are transformed as follows

$$SUM_{f(x)} \omega^{d-j} |x\rangle |j\rangle \rightarrow \omega^{f(x)} \omega^{d-m} |x\rangle |m\rangle. \quad (22)$$

Hence from (20) and (22) we have

$$SUM_{f(x)} |x\rangle |\phi\rangle = \omega^{f(x)} |x\rangle |\phi\rangle. \quad (23)$$

Therefore, the relation (17) holds.

We have  $|\psi_2\rangle$  by operating  $SUM_{f(x)}$  to  $|\psi_1\rangle$

$$SUM_{f(x)} |\psi_1\rangle = |\psi_2\rangle = \sum_{x \in K} \frac{\omega^{f(x)} |x\rangle}{\sqrt{(2d-1)^N}} |\phi\rangle. \quad (24)$$

After the general Fourier transform of  $|x\rangle$ , using the previous equations (11) and (24) we can now evaluate  $|\psi_3\rangle$  as follows

$$\begin{aligned} |\psi_3\rangle &= \sum_{z \in K} \sum_{x \in K} \frac{(\omega)^{x \cdot z + f(x)} |z\rangle}{(2d-1)^N} |\phi\rangle \\ &= \sum_{z \in K} \sum_{x \in K} \frac{(\omega)^{x \cdot z + g(a) \cdot x} |z\rangle}{(2d-1)^N} |\phi\rangle. \end{aligned} \quad (25)$$

Because we have

$$\sum_{x \in K} (\omega)^x = 0 \quad (26)$$

we may notice

$$\begin{aligned} \sum_{x \in K} (\omega)^{x \cdot (z + g(a))} &= (2d-1)^N \delta_{z+g(a), 0} \\ &= (2d-1)^N \delta_{z, -g(a)}. \end{aligned} \quad (27)$$

Therefore, the above summation is zero if  $z \neq -g(a)$  and the above summation is  $(2d-1)^N$  if  $z = -g(a)$ . Thus we have

$$\begin{aligned} |\psi_3\rangle &= \sum_{z \in K} \sum_{x \in K} \frac{(\omega)^{x \cdot z + g(a) \cdot x} |z\rangle}{(2d-1)^N} |\phi\rangle \\ &= \sum_{z \in K} \frac{(2d-1)^N \delta_{z, -g(a)} |z\rangle}{(2d-1)^N} |\phi\rangle \\ &= -|(g(a_1), g(a_2), g(a_3), \dots, g(a_N))\rangle |\phi\rangle \end{aligned} \quad (28)$$

from which

$$|(g(a_1), g(a_2), g(a_3), \dots, g(a_N))\rangle \quad (29)$$

can be obtained. That is to say, if we measure the first  $N$  qudits state of the state  $|\psi_3\rangle$ , that is,  $|(g(a_1), g(a_2), g(a_3), \dots, g(a_N))\rangle$ , then we can retrieve the following values

$$g(a_1), g(a_2), g(a_3), \dots, g(a_N) \quad (30)$$

using a single query. The method determines the maximum of and the minimum of the function  $g$  that the finite domain is  $\{a_1, a_2, a_3, \dots, a_N\}$ .

### III. CONCLUSIONS

In conclusion, we have presented the generalized Bernstein-Vazirani algorithm for determining an integer string. Given the set of real values  $\{a_1, a_2, a_3, \dots, a_N\}$  and a function  $g: \mathbf{R} \rightarrow \mathbf{Z}$ , we shall have determined the following values  $\{g(a_1), g(a_2), g(a_3), \dots, g(a_N)\}$  simultaneously. The speed of determining the values has been shown to outperform the classical case by a factor of  $N$ . The method has determined the maximum of and the minimum of the function  $g$  that the finite domain is  $\{a_1, a_2, a_3, \dots, a_N\}$ .

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