AN ESTIMATE OF THE MASS OF NEUTRINOS FROM NUCLEAR MASS DEFECTS.

R. Heyrovska

Academy of Sciences of the Czech Republic, J. Heyrovsky Institute of Physical Chemistry, Dolejskova 3, 182 23 Prague 8, Czech Republic. E-mail : rheyrovs@jh-inst. cas. cz

#### ABSTRACT

During the synthesis of a deuteron nucleus from a neutron and a proton, an electron, positron pair and an antineutrino, neutrino pair are released. The electrostatic energy (order of MeV) of the former pair is released as gamma radiation and is equivalent to the mass defect. Based on the largest mass defect per nucleon (MDPN) as the criterion of nuclear stability, it is estimated that in general, for a nuclide  $X(Z,N)$  with a fractional mass defect f, the MDPN(nu) due to neutrinos/antineutrinos alone is equal to  $fk(n)$  [ $Z/(Z+N)$ ], where the proportionality constant  $k(n)$ = 0.7 micro atomic mass unit. Since f is of the order of 0.01, MDPN(nu) is in nano-atomic mass units, which is the right order of the mass of neutrinos.

#### INTRODUCTION

The mass of the tiniest of particles, the "neutrinos", is expected to be about 10<sup>-9</sup>th of that of a hydrogen atom! A precise knowledge of this mass is sought by cosmologists and physicists [1-7] to answer many queries about our Universe.

1

Painstaking experiments planned to answer the question [2-3}, are jeopardized [8] by the over-riding difficulties and uncertainties. In this context [8,9], this report brings for the first time a different "possible" approach: estimates of the mass defects per nucleon attributable to neutrinos/antineutrinos, calculated using Lhe largest mass defect per nucleon criterion for nuclear stability. Using the atomic mass data for the 105 most stable nuclides, three ranges of masses (all of which are less than a few eV) have been found, "possibly" pertaining to the three kinds [1,5], the electron (lightest), muon and tau neutrinos.

## Outline of the idea:

 $\mathcal{O}(\mathcal{E}^{\mathcal{E}})$  .

The conversion of a neutron (n) to a proton  $(p^+)$  and the converse process are known to be accompanied by the release of (electrons,  $e^r$ , and) antineutrinos ( $v$ ) and (positrons,  $e^r$ , and) neutrinos  $(v)$  respectively  $[6, 7, 10-12]$ :

n ---> p<sup>+</sup> + e<sup>-</sup> +  $\nu$  (spontaneous decay) ...(1  $p^+$  ---> n + e<sup>+</sup> +  $\nu$ ...(2)

Therefore, during the synthesis of a nucleus,  $\texttt{X}^\texttt{z+}$ , by the fusion of N neutrons and Z protons,

 $\text{Nn + Zp}^*$  --->  $X^{z+}$  + ( $\nu$  +  $\nu$ ) ...(3)

the masses of neutrinos and antineutrinos can be expected to form

 $\mathcal{Z}$ 

,iri.'ii,. :;i-.i iit;:..: ".,..

a finite fraction of the mass defects (MD) or the binding energy (BE) (energy equivalent of MD) of the nuclides.

# Outline of the method and results obtained:

The MD for a nuclide X(Z,N) of atomic mass A, characterized by Z electrons, Z protons and N neutrons of rest masses  $m_e$ ,  $m_p$ and m<sub>n</sub> each respectively, is:

$$
MD = A_{z,N} - A = fA_{z,N} \qquad (4)
$$

where  $A_{z,N}$  (>A) stands for the sum  $Z(m_e+m_p)+Nm_n$ , f= (1-r), where r is Lhe ratio A/A',N. f is similar to the packing fract,ion [(Ar,'- M) /M, where  $M = Z+N$ , is the mass number]. The values of the rest masses are  $[13]$ :  $m_e + m_p = 1.007825032u$  and  $m_n = 1.008664924u$ . One atomic mass unit (u) is exactly  $1/12$ th the atomic mass of a  $^{12}C$ isotope (standard), and it's energy equivalent is 931 MeV.

In Table 1 are given the atomic mass (A) data taken from [13] for the most abundant stable (or the one with the longest half-life) nuclide of every element from Z=1 to Z=105. Nuclides marked with an \* occur in only one isotopic form. The neutron number N in column 3 is the rest mass based number [14],  $N \ge (A Zm_H$ ) / $m_n > N-1$ , where  $m_H$  stands for the sum  $m_b+m_e$ .

The mass defect per nucleon, MDPN, is an index of nuclear stability and is given bY

MDPN = MD/ (Z+N) =  $f_{A_{Z,N}}$ / (Z+N) =  $f_{A_{Z,N}}$  ... (5)

where  $m_{z,N}$  stands for the mean mass per nucleon,  $A_{z,N}/(Z+N)$ . These are given in column 5, Table 1. Nuclides with large MDPN are considered more stable than others. Note that Fe has the largest MDPN.

Using the accepted. largest MDPN as the criterion of nuclear stability, the equation for the atomic mass of any nuclide,

 $A = Zm_{H} + N'm_{n}$  ... (6)

where  $N \ge N' > N-1$ , was shown [14] to give the highest MDPN.

Therefore, the mass defect  $MD = (N-N')m_n$  and the MDPN defined by Eq.5 is equal to  $[14]$ ,

$$
MDPN = (N-N') m_n / (Z+N)
$$
 (7)

This equation implies that the formation of a nuclide of atomic mass A from the total mass  $A_{z,N}$  takes place at the expense of an amount of mass  $(N-N')$ m<sub>n</sub> during the fusion of the neutrons with protons. This is also in accord with the earlier view [12], that during the fusion of neutrons and protons to form a nucleus, the protons are restored at the expense of (the unstable) neutrons.

As neutrons break-up into protons, electrons and antineutrinos (cf: Eq.1), the mass  $m_n$  in Eq.7 splits into  $m_H$  and  $\Delta m_n$ . Thus,

MDPN =  $(N-N') m_H / (Z+N)$  +  $(N-N') \Delta m_H / (Z+N)$  ...(8) MDPN(1) MDPN(2)

4

where  $\Delta m_n = m_n - m_H$  (= 0.00084u). The values of MDPN(2) are of the order of  $10^{-6}u$  as can be seen from column 6, Table 1. As  $\Delta m_n/m_{\rm H}$  =  $8.4x10^{-4}$ , MDPN $(2)$  << MDPN $(1)$ .

 $\Box$ 

The MDPN resulting from the break-up of a mass  $\Delta m_n$  from the mean mass per nucleon,  $m_{z,N}$ , (cf: Eq.5), is given by,

MDPN = 
$$
f[m_{z,N} - \Delta m_n]
$$
 +  $f \Delta m_n$  ... (9)  
MDPN (1') MDPN (2')

where each of the two terms on the right are nearly equal to the corresponding ones in Eq.8.

For nuclear stabilty, Eq. (9) should conform with Eq. (8), which is based on the highest MDPN criterion. Therefore, the terms MDPN(1') and MDPN(2') gain and lose masses respectively and become equal to MDPN(1) and MDPN(2) . Thus:

 $MDPN = MDPN(1) + MDPN(2)$ = MDPN  $(1') + \{ [fZ/(Z+N)] (\Delta m_n)^2/m_n \}$ + MDPN(2') - { $[fZ/(Z+N)](\Delta m_n)^2/m_n$ } ...(10)

The second term on the right within each of the  $\{ \}$  brackets is of the order of 10<sup>-9</sup>u, since it is the product of the fraction  $[fZ/(Z+N)]$ , which is about 10<sup>-2</sup>, and the constant, k(n) =  $\Delta m_n^2/m_n$  = 6.995x10-7u. Note that, this term is in the expected range of the mass of neutrinos/antineutrinos! On considering this as the mass defect due to neutrinos/antineutrinos, MDPN $(v)$ ,

 $\mathfrak{S}$ 

## $MDPN(V) = [fk(n) Z/(Z+N)]$

I

the actual values for various nuclides are as shown in the last column in Table 1. The +MDPN( $\nu$ ) and -MDPN( $\nu$ ) terms probably indicate that the antineutrinos released by neutrons in term (2') are absorbed by the protons in term  $(1')$ , and therby the MDPN conforms with the sum MDPN(1) +MDPN(2) of Eq. (8) .

Figure 1 shows a plot of MDPN( $\nu$ ) vs the atomic number, Z. The three regions probably indicate the three types of light neutrinos, the electron (lightest), muon and the tau neutrinos (or different proportions of the three kinds of neutrinos).

From Eq.11, the total mass defect of the nuclide  $X_{z,N}$  due to neutrinos,  $MD(V)$  [=  $(Z+N)MDPN(V)$ ] is found to be,

$$
MD(\nu) = f k(n) Z \qquad \qquad \ldots (12)
$$

Figure 2 shows that  $MD(V)$  varies smoothly with Z (since fZ varies smoothly with Z). Note again the three regions.

On dividing Eq. (11) by Eq. (5), one finds that  $MDPN(\nu)$  (or  $MD(V)$ ) forms a well-defined fraction of MDPN (or MD),

MDPN (
$$
\nu
$$
) /MDPN = MD( $\nu$ ) /MD = (Z/A<sub>z,N</sub>) k(n) ... (13)

independent of f. Therefore, this fraction can be calculated for any value of Z and N. Fig. 3 shows the linear dependence of this fraction on  $Z/A_{Z,N}$ .

...(11)

### References:

t\_\_

1. J. N. Bahcall, Neutrino Astrophysics (Cambridge Univ. Press, 1989).

2. G. Taubes, Science 267 (1995) <sup>789</sup>

3. J. N. Wilford, 'The New York Times', New York, Jan. 31, 1995.

4. R. S. Raghavan, Science 267 (1995) 45.

5. J. N. Bahcall, et al. NaLure 375 (1995) 29'

6. L. N. Kraus, Sci. Amer. 255 (Dec. 1986) 50 '

7. D. N. Schramm and G. Steigman, Sci. Amer. 258 (June 1988) 44. 8. J. Glanz, Science 269 (1995) L67L.

g. "Some Unso1ved Problems in Astrophysics", Conference in honour of J. N. Bahcall's 60th Birthday, Princeton, 1995. 10. M. K. Moe and s. P. Rosen, Sci. Amer . 26L (Nov. L989) 30.

11. H. Harari, Sci. Amer. 248 (April, 1983) 48.

12. R. Heyrovska, Proc. II, J.Heyrovsky Centennial Congress and 41st Mtg. of the ISE, Prague, 1990.

13. G. Audi and A. H. Wapstra, Nuclear Physics A, Dec.

1993.

14. R. Heyrovska, J. Chem. Educn. 69 (1992) 742.

#### Acknowledgement

The author is grateful to Albert, Emil and David Heyrovskys for their interest and discussions and to Neela Heyrovska (to whom this work is dedicated) for her encouragement to finish and publish this work.

Table 1: Atomic Mass (A) from (13); and MDPN( $\dot{V}$ ), Eq.  $\theta$ . الماسات

 $\frac{1}{2}$ 

 $\ddot{\phantom{1}}$ 



 $\overline{7}$ 



\* Nuclides which occur in only one isotopic form

8

## Legends for Figures:

## Figure 1:

The mass defects per nucleon due to neutrinos / antineutrinos, MDPN( $\nu$ ) (cf: Eq.11) vs the atomic number Z.

# Figure 2:

The total mass defect of the nuclide due to neutrinos / antineutrinos, MD( $\nu$ ) [= (Z+N)MDPN( $\nu$ ), cf: Eq.12] vs the atomic number, Z.

## Figure 3:

The linear dependence of the fraction of mass defect,  $MD(V)/MD$ , due to neutrinos on the ratio  $Z/A_{z,N}$  (cf: Eq.13).



 $Fig. 1 (R.4)$ 



 $\epsilon$ 



# $Fig. 3/7}$