# Nuclear Quantum Gravity – A First Review

U. V. S. Seshavatharam<sup>1</sup> & S. Lakshminarayana<sup>2</sup>

<sup>1</sup>Hon. Faculty, I-SERVE, S. No-42, Hitex Rd., Hitech city, Hyderabad-84, Telangana, India <sup>2</sup>Department of Nuclear Physics, Andhra University, Visakhapatnam-03, AP, India Corresponding Emails: seshavatharam.uvs@gmail.com; lnsrirama@gmail.com;

Abstract: To have a unified model of nuclear quantum gravity, it seems quite reasonable to consider a large nuclear gravitational constant,  $G_s \approx (3.3 \pm 0.03) \times 10^{28} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}$ . In this context, we show practical applications pertaining to micro physics as well as macro physics. We would like to suggest that: (1)  $G_s$  plays a crucial role in understanding quantum theory of light, photoelectric work functions, superconductivity, nuclear binding energy, nuclear root mean square charge radii, root mean square radius of proton, neutron life time, neutron-proton mass difference, nuclear stability, nuclear magnetic dipole moments, weak coupling angle, Fermi's weak coupling constant, proton melting point and total energy of electron in Hydrogen atom etc.; (2) Nuclear binding energy can be understood with a single energy coefficient of magnitude 10 MeV. (3) Newtonian gravitation al constant  $G_N$  and the proposed  $G_s$  play a joint role in understanding neutron star mass generation as well as proton mass generation; and (4) Considering  $G_s$  as a characteristic feature of magnetism, celestial bodies 'mass dependent' magnetic dipole moments can be estimated. (5) Magnitude of  $G_N$  can certainly be estimated from microscopic elementary physical constants.

Keywords: Final unification, nuclear gravitational constant, Newtonian gravitational constant, nuclear quantum gravity.

#### 1. Introduction

In this paper, we review our recently published views on nuclear quantum gravity [1] for better presentation on nuclear stability and binding energy [2-9], root mean square radius of proton, neutron life time, issues connected with coupling constants and other minor changes that may help in developing this new subject. The most desirable cases of any unified description [10-13] are:

- 1. To implement gravity in microscopic physics.
- 2. To develop a model of quantum gravity.
- 3. To simplify the complicated issues of known physics.
- 4. To predict new effects, arising from a combination of the fields inherent in the unified description.

In this context, for a better understanding, we would like to review our recent publications [14-25] in the following way:

"By replacing  $\left(\frac{\hbar c}{m_p^2}\right)$  with a large gravitational constant

of magnitude,  $G_s \approx 3.32956 \times 10^{28}$  assumed to be associated with nuclear structure, a basic model of 'nuclear quantum gravity' can be developed". Qualitatively, our assumption is not new and is having a long standing history [26-34]. For more information, readers are strongly encouraged to see Abdus Salam's 'Strong gravity' concept [32]. Recently, O. F. Akinto and Farida Tahir elaborated their work on 'modified strong gravity concepts' pertaining to QCD and general relativity in arXiv preprint [33]. In 2013, Roberto Onofrio [34] proposed a very interesting concept: Weak interactions are peculiar manifestations of quantum gravity at the Fermi scale, and that the Fermi coupling constant is related to the Newtonian constant of gravitation. In his opinion, at atto-meter scale, Newtonian gravitational constant seems to reach a magnitude of  $8.205 \times 10^{22}$  m<sup>3</sup>kg<sup>-1</sup>sec<sup>-2</sup>.

#### 2. To unite nuclear and sub nuclear interactions

The modern theory of strong interaction is Quantum chromodynamics (QCD) [35-38]. It explores baryons and mesons in broad view with 6 quarks and 8 gluons. According to QCD, the four important properties of strong interaction are: 1) color charge;2) confinement; 3) asymptotic freedom ;4) short-range nature( $\leq 10^{-15}$  m). Color charge is assumed to be responsible for the strong force to act on quarks via the force carrying agent, gluon. Experimentally it is well established that, strength of strong force depends on the energy of the interaction or the distance between particles. At lower energies or longer distances: a) color charge strength increases; b) strong force becomes 'stronger'; c) nucleons can be considered as fundamental nuclear particles and quarks seem to be strongly bound within the nucleons leading to 'Quark confinement'. At high energies or short distances: a) color charge strength decreases; b) strong force gets 'weaker'; 3) colliding protons generate 'scattered free quarks' leading to 'Quark Asymptotic freedom'. Based on these points, low energy nuclear scientists assume 'strong interaction' as a strange nuclear interaction associated with binding of nucleons and its implications were not considered. High energy nuclear scientists consider nucleons as composite states of quarks and try to understand the nature and strength of strong interaction  $(\alpha_s)$  at sub nuclear level. According to QCD,  $(\alpha_s)$  decreases with increasing interaction energy. By definition, at low energy scales,  $\alpha_s \approx 1$  and by experiments and observations, at 80 to 90 GeV energy scales,  $\alpha_s \approx 0.1186$ .

At this juncture, one important question to be answered and reviewed at basic level is: How to understand nuclear interactions in terms of sub nuclear interactions? Unfortunately, 1) At 1.2 fm scale, there is no practical evidence or applications for the basic definition of  $\alpha_s \approx 1$ . 2) With current concepts of QCD, one cannot

explain the observed nuclear binding energy scheme. 3) Famous nuclear models like, Liquid drop model and Fermi gas model [39-41] are lagging in answering this question. To find a way, we would like to suggest that, by implementing the 'strong coupling constant' of magnitude 0.1186 in low energy nuclear physics, nuclear charge radius, Fermi's weak coupling constant and strong coupling constant can be studied in a unified picture. Proceeding further, close to beta stability line, nuclear binding energy can be addressed with a single energy coefficient of (8.9 to 10.0) MeV [3-6].

### 5. Microscopic and macroscopic Applications of the large nuclear gravitational constant

With the proposed assumption, quantum theory of light [42], photo electric work functions [43], super conductivity [44] nuclear physics, sub nuclear (particle) physics, electroweak theory, physics pertaining to planetary dipole magnetic moments, nuclear astrophysics and Bohr's theory of Hydrogen atom can be studied in a unified approach.

1) Quantum theory of light can be understood with,

$$
hc \cong \sqrt{\left(\frac{m_p}{m_e}\right)\left(\frac{e^2}{4\pi\varepsilon_0}\right)\left(G_s m_p^2\right)}.
$$

2) Magnetic flux quantum in super conductivity can be understood with,

$$
\Phi_0 \cong \frac{h}{2e} \cong \frac{1}{2} \sqrt{\left(\frac{m_p}{m_e}\right) \left(\frac{\mu_0}{4\pi}\right)} \left(G_s m_p^2\right)
$$

3)  $G_F$  being the Fermi's weak coupling constant,

$$
G_s m_p m_e \cong \left(\frac{\pi^3 \varepsilon_0 G_F^2 c^8}{e^2 G_s^2}\right)^{\!\!\frac{1}{3}}.
$$

4) Photoelectric work functions can be understood with,  $\frac{1}{C}$  m m  $\frac{1}{1}$ 

 $\frac{3}{3} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  to  $-A^3 \left( \frac{G_s m_p}{2r_Z} \right)$  $S_Z \approx -Z^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  to  $-A^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  $W_Z \approx -Z^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  to  $- A^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  $\approx -Z^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  to  $-A^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right)$  where

 $r_Z$  is the radius of atom [45], Z is the atomic number and  $\vec{A}$  is the mass number.

5) Near to beta stability line, nuclear binding energy can be understood with,  $\frac{1}{2} \left( e^2 G_s m_p^3 \left/ 4 \pi \varepsilon_0 \hbar^2 \right) \approx 10.09 \text{ MeV}.$ 

$$
\frac{1}{2} (e^{i\theta} G_s m_p / 4\pi \epsilon_0 n) = 10.09 \text{ MeV}.
$$
  
6) Nucleon mass difference, nuclear stability and neutron  
life time can be understood with

$$
\[ e^2 G_s m_p^3 \Big/ 4\pi \varepsilon_0 \big( \hbar/2 \big)^2 \, \] \cong 80.7 \text{ MeV}.
$$

7) Nuclear charge radius can be addressed with,  $R_0 \cong \frac{2G_s m_p}{r} \cong 1.239$  fm Or

$$
\kappa_0 \equiv \frac{e^2}{c^2} \equiv 1.239 \text{ m or}
$$
\n
$$
\left(\ln \sqrt{\frac{G_s}{G_N}}\right) \sqrt{\frac{e^2 G_s}{4\pi \varepsilon_0 c^4}} \approx 1.374 \text{ fm.}
$$

8) For medium, heavy and super heavy atomic nuclei, root mean square charge radii can be fitted with

For medium, heavy and super heavy atomic nuclei, root  
mean square charge radii can be fitted with  

$$
R_{(Z,A)} \cong \left\{ Z^{1/3} + \left(\sqrt{Z(A-Z)}\right)^{1/3}\right\} \left(\frac{G_s m_p}{c^2}\right)
$$

$$
\cong \left\{ Z^{1/3} + \left[Z(A-Z)\right]^{1/6}\right\} \left(\frac{G_s m_p}{c^2}\right).
$$
Hadronic melting points [46] can be understood with,  

$$
T_{hadron} \cong \frac{hc^3}{8\pi k_B G_s m_{hadron}}.
$$
Strong coupling constant can be understood with.

9) Hadronic melting points [46] can be understood with,

.

$$
T_{hadron} \cong \frac{\hbar c^3}{8\pi k_B G_s m_{hadron}}
$$

10) Strong coupling constant can be understood with,

 <sup>0</sup> <sup>2</sup> 2 2 1 0.1152 s s p c G m or 2 0 1 exp 0.1185 4 5. s p s s e e G m m 

- 11) Weak coupling angle can be understood with, 2  $a \approx 4\pi\epsilon_0$  $\sin^2 \theta_W \approx \frac{4\pi \varepsilon_0 G_s m_p m_e}{r^2} \approx 0.2198.$  $G_s m_p m_e$ e  $\theta_W \cong \frac{4\pi\varepsilon_0 G_s m_p m_e}{2} \cong$
- 12) Up and down quark mass ratio can be understood with, 0 2  $\left(\frac{u}{d}\right) \approx \sqrt{\frac{4\pi\epsilon_0 G_s m_p m_e}{e^2}} \approx \sin\theta_W \approx 0.469$  $m_{\mu}$   $4\pi\varepsilon_0 G_s m_{\rho} m_e$  $m_d$   $\int - \sqrt{1 - e^2}$  $\left(\frac{m_u}{m}\right) \approx \sqrt{\frac{4\pi \varepsilon_0 G_s m_p m_e}{2}} \approx \sin \theta_W \approx 0$  $\left(m_d\right)^{-1}$
- 13) Proton's magnetic dipole moment can be understood with,

$$
\mu_{proton} \approx \left(\frac{m_u}{m_d}\right) \left(\frac{eG_s m_p}{c}\right) \approx 1.396 \times 10^{-26} \text{ J/T esla}
$$

14) Neutron's magnetic dipole moment can be understood with,

$$
\mu_{neutron} \approx \left(\frac{m_u}{m_d}\right)^{\frac{3}{2}} \left(\frac{eG_s m_n}{c}\right) \approx 9.573 \times 10^{-27} \text{ J/T esla}
$$

15) Ratio of proton to neutron magnetic diploe moment can be addressed with

$$
\frac{\mu_{proton}}{\mu_{neutron}} \approx \sqrt{\frac{m_u}{m_d}} \approx \sqrt{\sin \theta_W} \approx 0.685
$$

16) Ratio of proton-electron magnetic dipole moments can

be understood with, 
$$
\left(\frac{\mu_{proton}}{\mu_{electron}}\right) \approx \frac{G_s m_p m_e}{\hbar c}
$$
.

- antium theory of light can be understood with,<br>  $\sin^2 \theta_W = \frac{\sin^2 \theta_W = \frac{\sin^2 \theta_W}{c^2}}{\cos^2 \theta_W} = 0.2198.$ <br>  $\therefore \frac{m}{2m} = \sqrt{\left(\frac{m_p}{m_c}\right)\left(\frac{e^2}{4\pi c_0}\right)(G_s m_p^2)}$ ,  $\left(\frac{m_u}{m_d}\right) = \sqrt{\frac{4\pi c_0}{c^2}}\frac{G_m m_p m_c}{c^2} = \sin \theta_W = 0.469$ <br>
dest  $\mathbb{E}[\sqrt{\frac{m_e}{m_e}}] \left(\frac{\pi_{E}}{m_e}\right) \left(\frac{4\pi_{E}}{4\pi_{E}}\right)^{(1/2)}$  Find the number conductivity can be  $\frac{1}{2}$  in  $\frac{1}{e^2}$   $\frac{1}{e^2}$  and  $\theta_W$  and  $\theta_W$  and  $\theta_W$  and  $\theta_W$  is  $\frac{1}{2}$  in  $\theta_W$  being the Fermi is super 17) Nuclear Planck mass can be defined as,  $m_{\text{mol}} \approx \sqrt{\hbar c/G_s} \approx 546.6 \text{ MeV}/c^2$ . Based on this new mass unit, a quantized model mechanism can be developed for understanding the hadronic mass spectrum. In our recent publication [14], by considering 546.6 MeV/ $c<sup>2</sup>$  as a characteristic neutral hadronic fermion, we have developed a toy model for understanding the hadronic mass spectrum.
	- 18)  $m_{npl} \approx \sqrt{\hbar c/G_s} \approx 546.6 \text{ MeV}/c^2 \text{ can be considered}$ as a characteristic dark matter candidate [33].
	- 19) Total energy of electron in Hydrogen atom can be understood with,

viXra:1710.0212v2 Dedicated to Dr. Abdus Salam  $\begin{array}{|l|}\n 2\n \end{array}$ 

$$
-\frac{1}{2}\left(\frac{G_s m_p m_e}{\hbar c}\right)^2 \frac{\sqrt{m_p m_e c^2}}{2n^2} \approx -\frac{14.0}{n^2} \text{ eV, where}
$$

 $\left(\frac{1}{2n^2}\right)$  c can be considered as the probability of finding

electron in its orbits labeled as  $n = 1, 2, 3, \dots$ 20) Neutron star mass [47] or radius can be understood

$$
\quad\text{with,}\ \sqrt{\frac{G_{\!s}}{G_N}}\ .
$$

- 21) Mass dependent planetary magnetic dipole moments [48] can be understood with,  $\alpha$  $G_s m_p m_e \parallel e G_s M_{planet}$  $\left(\frac{G_s m_p m_e}{\hbar c}\right)\left(\frac{e G_s M_{planet}}{2c}\right).$ .
- 22) Root mean square radius of proton can be fitted with:

$$
R_p \approx \ln\left(\frac{e^2}{4\pi\varepsilon_0 G_N m_p^2}\right) \sqrt{\frac{e^2}{4\pi\varepsilon_0 G_s m_p^2}} \left(\frac{\hbar}{m_p c}\right) \le 0.87 \text{ fm}
$$
  
or  $2\pi R_p \approx 2 \ln\left(\frac{G_s}{G_N}\right) \sqrt{\frac{e^2 G_s}{4\pi\varepsilon_0 c^4}} \approx 0.8746 \text{ fm.}$ 

23) Neutron life time can be understood with,

$$
t_n \approx \left(\frac{G_s}{G_N}\right)^{\frac{2}{3}} \sqrt{\frac{4\pi\varepsilon_0 G_s m_p^2}{e^2}} \left(\frac{\hbar}{m_n c^2}\right) \approx 885.45 \text{ sec}
$$

24) Proton-electron mass ratio can be understood with,

$$
\begin{cases}\n\left(\frac{m_p}{m_e}\right) \cong \left(\left(\frac{G_s}{G_N}\right)\left(\frac{G_s m_e^2}{\hbar c}\right)\right)^{\frac{1}{10}} \\
\cong \left(\frac{G_s m_e^2}{G_N m_{npl}^2}\right)^{\frac{1}{10}} \cong 1836.3\n\end{cases}
$$
\n
$$
m_{m1} \cong \sqrt{\hbar c/G_s} \cong 546.6 \text{ MeV/g}
$$

where  $m_{npl} \approx \sqrt{\hbar c/G_s} \approx 546.6 \text{ MeV}/c^2$ .

### 6. To understand proton's melting point

With reference to Hawking black hole temperature formula [46], melting point of proton [36] can be understood with:

$$
T_{proton} \approx \frac{\hbar c^3}{8\pi k_B G_s m_p} \approx 0.15 \times 10^{12} \text{ K}
$$
 (1)

Based on this relation and with reference to up quark, quark melting points can be expressed with the following kind of relation.

$$
T_{quark} \cong \left(\frac{m_q}{m_{up}}\right) \frac{\hbar c^3}{8\pi k_B G_s m_{up}}\tag{2}
$$

where can be  $\frac{m_q}{q}$ up  $m_a$  $\left(\frac{m_q}{m_{up}}\right)$  re represents the ratio of mass of any

quark to mass of up quark. Based on this relation, for up quark of rest energy 2 MeV, its corresponding

$$
T_{up} \approx 69
$$
 Tera K and  $8\pi k_B T_{up} \approx 236$  MeV. This  
energy can be compared with currently believed QCD  
energy scale of 270 MeV.

## 7. To fit and understand Fermi's weak coupling constant

Fermi's weak coupling constant [35] can be fitted with a relation of the kind,

$$
G_F \approx \left(\frac{m_e}{m_p}\right)^2 \hbar c R_0^2 \approx G_s m_e^2 \left(2 \sqrt{\frac{G_s \hbar}{c^3}}\right)^2
$$
  

$$
\approx \frac{4G_s^2 m_e^2 \hbar}{c^3} \approx 1.44021 \times 10^{-62} \text{ J.m}^3
$$
 (3)

where  $\sqrt{\frac{G_s \hbar}{c^3}} \approx 0.36$  fm.  $c^{\cdot}$  $\frac{\hbar}{c} \approx 0.36$  fm can be considered as the

characteristic Nuclear Planck length and  $2C_{\nu}$  $2\sqrt{\frac{G_s \hbar}{r}} \cong \frac{2G_s m_{npl}}{r}.$  $\sqrt{a^2 + b^2}$ 

$$
2\sqrt{\frac{G_s n}{c^3}} \cong \frac{2G_s m_{npl}}{c^2}.
$$

#### 8. To fit and understand strong coupling constant

Based on the proposed  $G_s$ , strong coupling constant can be fitted with the following kind of relation.

$$
\alpha_s \approx \left[\frac{\hbar c}{G_s m_p^2}\right]^2 \approx 0.115543\tag{4}
$$

Based on this relation,

$$
R_0 \cong \left(\frac{1}{\sqrt{\alpha_s}}\right) \frac{2\hbar}{m_p c} \cong 1.239 \text{ fm}
$$
 (5)

$$
\alpha_s F_W \cong \frac{4\hbar^3 m_e^2}{m_p^4 c} \tag{6}
$$

$$
\alpha_s \approx \frac{4\hbar^3 m_e^2}{m_p^4 c F_W} \approx \left(\frac{m_e}{m_p}\right)^2 \left(\frac{4\hbar^3 m_e^2}{m_p^2 c F_W}\right) \tag{7}
$$

Based on this relation and considering relativistic energy of proton, it is possible to show that, o fit and understand strong coupling constant<br>
on the proposed  $G_s$ , strong coupling constant<br>
can do with the following kind of relation.<br>  $\alpha_s \approx \left[ \frac{\hbar c}{G_s m_p^2} \right]^2 \approx 0.115543$  (4)<br>
on this relation,<br>  $R_0 \approx \left( \frac{1}{\sqrt{$  $\alpha_s \propto \left\{ \left[1 - \left(v/c\right)^2\right] / m_p^2 \right\}$ . Qualitatively, this kind of observation seems to be in-line with modern QCD concepts. m (5)<br>
(6)<br>
(6)<br>
(7)<br>
(7)<br>
(7)<br>
relativistic energy of<br>
to show that,<br>
ely, this kind of<br>
with modern QCD<br>
s<br>
s<br>
can construct two<br>
mits in the following<br>  $\left(m_p c^2\right)$ 

# 9. Two characteristic energy units

Based on the proposed  $G_s$ , one can construct two (characteristic and practical) energy units in the following way.

$$
E_X \approx \left(\frac{e^2}{4\pi\varepsilon_0\hbar c}\right) \left(\frac{G_s m_p^2}{\hbar c}\right) \left(m_p c^2\right)
$$

$$
\approx \left(\frac{e^2 G_s m_p^3}{4\pi\varepsilon_0\hbar^2}\right) \approx 20.174 \text{ MeV}
$$

viXra:1710.0212v2 Dedicated to Dr. Abdus Salam

(8)

By considering 'Uncertainty relation' and replacing  $(h)$ with  $(\hbar/2)$ , it is possible to construct another energy unit.

$$
E_Y \approx \left(\frac{e^2 G_s m_p^3}{4\pi \varepsilon_0 \left(\hbar/2\right)^2}\right) \approx \left(\frac{4e^2 G_s m_p^3}{4\pi \varepsilon_0 \hbar^2}\right)
$$
  
 
$$
\approx 4E_X \approx 80.696 \text{ MeV}
$$
 (9)

It may be noted that, with reference to Fermi gas model of the nucleus,  $E_X$  seems to represent the nucleon's mean kinetic energy per nucleon and  $2E_X \cong (E_Y/2) \cong \sqrt{E_X E_Y} \cong 40.35$  MeV seems to represent the depth of nuclear potential energy.

# 10. Fitting neutron-proton mass difference

Neutron-proton mass difference can be understood with the following relation:

$$
\left(\frac{m_n c^2 - m_p c^2}{m_e c^2}\right) \approx \ln \sqrt{\frac{E_Y}{m_e c^2}} \approx \ln \sqrt{\frac{4e^2 G_s m_p^3}{4\pi\varepsilon_0 \hbar^2 m_e c^2}}
$$

$$
\approx \ln \sqrt{\frac{80.696 \text{ MeV}}{0.511 \text{ MeV}}} \approx \ln (4\pi) \approx 2.531 \tag{10}
$$

#### 11. Fitting neutron life time

Neutron life time  $t_n$  can be understood with the following relation:

Fitting neutron life time  
\ntron life time 
$$
t_n
$$
 can be understood with the following  
\nion:  
\n $t_n \approx \exp \left\{ \frac{E_Y}{(m_n - m_p)c^2} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{m_nc^2} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{2\pi} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{2\pi} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{2\pi} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text{ MeV}}{1.2933 \text{ MeV}} \right\} \times \left( \frac{\hbar}{2\pi} \right)$   
\n $\approx \exp \left\{ \frac{80.696 \text$ 

This fitted value can be compared with material bottle and cold beam experimental results:  $(878.5 \pm 0.8)$  sec and  $(887.7 \pm 2.2)$  sec [49]. See section 22 for correlating bottle and cold beam experimental results.

#### 12. Understanding nuclear stability

Stable mass number corresponding to  $Z$  can be estimated with the following relation [2]:

$$
(A_s - 2Z) \approx \left\{ \frac{m_e c^2}{E_Y} \right\} Z^2 \approx \left( \frac{4\pi \varepsilon_0 \hbar^2 m_e c^2}{4e^2 G_s m_p^3} \right) Z^2
$$
  

$$
\approx \left( \frac{0.511 \text{ MeV}}{80.696 \text{ MeV}} \right) Z^2 \approx 0.00633 (Z)^2 \approx kZ^2
$$
 (12)

extrainty relation' and replacing (*h*)<br>
where  $k \approx \left\{ \frac{m_e c^2}{E_y} \right\} \approx \left( \frac{4\pi \varepsilon_0 h^2 m_e c^2}{4e^2 G_s m_p^3} \right) \approx 0.00633$ . Using<br>
let to construct another energy unit.<br>
this new number, nuclear binding can be estimated. where  $\frac{2}{2} \ge \left( \frac{4\pi \varepsilon_0 \hbar^2 m_e c^2}{4 a^2 G m^3} \right)$  $\left.\frac{4\pi\epsilon_0\hbar^2m_ec^2}{r^2}\right| \geq 0.00633.$ 4  $_e^e e^ \vert \sim \vert$   $4\pi \varepsilon_0 h^- m_e^e$  $\gamma$  J (4e<sup>-</sup> $G_s m_p$ )  $k \cong \left\{\frac{m_e c^2}{r}\right\} \cong \left\{\frac{4\pi\epsilon_0 \hbar^2 m_e c^2}{r^2}\right\}$  $E_Y \parallel \equiv \parallel 4e^2 G_s m_p^3$  $\geq \left\{\frac{m_e c^2}{E}\right\} \cong \left(\frac{4\pi \varepsilon_0 \hbar^2 m_e c^2}{4\pi^2 \varepsilon_0^3}\right) \cong 0$  $\left[ E_Y \right]$   $\left[ 4e^2 G_s m_p^3 \right]$  $\frac{\hbar^2 m_e c^2}{r^2}$  = 0.00633. Using

this new number, nuclear binding can be estimated. With further study, it is also possible to show that,

$$
k \approx \left\{ \frac{m_e c^2}{E_Y} \right\} \approx \frac{0.71 \text{ MeV}}{8.89 \text{ MeV}}.
$$
 (13)

where 0.71 MeV represents the coulombic energy coefficient and 8.89 MeV can be considered as the maximum binding energy per nucleon.

### 13. Understanding nuclear binding energy

Based on the new integrated model proposed by N. Ghahramany et al [3,4] and with reference to relation (12), it is possible to show that,  $Z \approx (40 \text{ to } 83)$ , close to the beta stability line,

<sup>1</sup>Uncertainty relation' and replacing (h)<sup>1</sup> where 
$$
k \approx \left\{ \frac{m_c c^2}{E_y} \right\} \approx \left( \frac{4\pi \epsilon_0 \hbar^2 m_c c^2}{4\epsilon^2 G_s m_p^2} \right) \approx 0.00633
$$
. Using  
possible to construct another energy unit.  

$$
\frac{e^2 G_s m_p^3}{4\pi \epsilon_0 (h/2)^2} = \left\{ \frac{4e^2 G_s m_p^3}{4\pi \epsilon_0 \hbar^2} \right\}
$$
 (9)  
 $E_X \approx 80.696$  MeV  
that, with reference to Fermi gas model of  
energy, we are nucleon's mean  
to represent the nucleous in the direction of the  
energy term. The nucleon is the mass of the  
energy term. The nucleon is the mass of the  
energy term. The nucleon is the mass of the  
energy term. The nucleon is the mass of the  
energy term. The decoupling energy is the mass of the  
phof nuclear potential energy.  

$$
= \sqrt{E_X E_Y} \approx 40.35
$$
 MeV seems to  
the mass difference. Based on the new integrated model proposed by N.  
mass difference can be understood with the  
mass difference as the data of the mass of the  
kinning energy. The mass of the  
kinning energy is the  
transluting at a [3,4] and with reference to relation (12),  
on:  

$$
\frac{c^2}{2} = \ln \sqrt{\frac{E_Y}{(m_c c^2)}} = \ln \sqrt{\frac{4e^2 G_s m_p^3}{4\pi \epsilon_0 \hbar^2 m_c c^2}}
$$
 (B)<sub>4,5</sub> =  $\left[ A_s - \left( \frac{N_s^2 - Z^2}{3Z} \right) \right] \times 9.5$  MeV  
from life time  
the t<sub>n</sub> can be understood with the following  
surbending (B) is the mass of the  
sublications [5,6] and first four terms of the semi empirical  
mass formula (SEMF), close to the beta stability line,  

$$
\left[ m_n - m_p \right] c^2 \right\} \times \left( \frac{\hbar}{m_c c^2} \right) = 877.9 \text{ sec}
$$
 (B)<sub>4,5</sub> =  $\left[ A_s - 4_s^{1/3} - \frac{k4_s \sqrt{N_s Z}}{3.42} - 1 \right] \times \left\{ \frac{E_X}{2} \right\}$  (15)  
29.33 MeV

2  $Z^2$ where,  $\left[\frac{N_s^2 - Z^2}{Z}\right] \cong kZA_s$ .  $\left[N_s^2 - Z^2\right]$  $\left[\frac{N_s - Z}{Z}\right] \cong kZA_s$ . Based on this strange and

simple relation and with reference to our recent publications [5,6] and first four terms of the semi empirical mass formula (SEMF), close to the beta stability line, for  $(Z = 2 \text{ to } 100)$ , it is possible to show that,

 1 3 1 3 1 3.42 2 1 10.09 MeV 3.35 s s s X <sup>A</sup> s s s s s s kA N Z E B A A kA N Z A A (15) 

where, 
$$
\left(\frac{k}{3.35}\right) \approx \alpha_s \left(\frac{a_c}{2a_a}\right)
$$
. See table-1.

For  $Z = 50$  and  $A = 100$  to 136, estimated binding energy range is (857 to 1140) MeV and can be compared with reference binding energy [40] range of (806 to 1105) MeV. It is for further study. With reference to SEMF, close to the beta stability line, it is also possible to show that,

$$
\frac{\left(A_s - 2Z\right)^2}{A_s} \cong \left(k^2 A_s N_s \sqrt{Z}\right) \tag{16}
$$

.

Let,  
\n
$$
\begin{cases}\na_v \equiv a_s \equiv a_a \approx 14.8 \text{ MeV} \approx (3/2) \times 10.0 \text{ MeV} \\
\text{and } a_c \equiv 0.71 \text{ MeV}\n\end{cases}
$$
\nIf so,

viXra:1710.0212v2 Dedicated to Dr. Abdus Salam

 $\mathbf{L}$ 

$$
(B)_{A_s} \approx \begin{bmatrix} A_s - A_s^{2/3} \\ -0.0473 \left[ \frac{Z(Z-1)}{A_s^{1/3}} \right] - \left( k^2 A_s N_s \sqrt{Z} \right) \end{bmatrix} \times 14.8 \text{ MeV}
$$
  
\n
$$
(17)
$$
\n  
\ncomparison with SEMF, by replacing  $A_s$  with  $A$  in  
\nlation (17) and by considering a multiplication factor of  
\n
$$
(17)
$$
\n
$$
(18)
$$
\n
$$
\left( \frac{\mu_{proton}}{\mu_{electron}} \right) \approx \frac{G_s m_p m_e}{\hbar c}.
$$
\nBased on this observation,

In comparison with SEMF, by replacing  $A_s$  with  $A$  in relation (17) and by considering a multiplication factor of (B)<sub>A,</sub>  $\approx \begin{bmatrix} A_s - A_s^{2/3} \\ -0.0473 \Big[ \frac{Z(Z-1)}{A_s^{1/3}} \Big] - (k^2 A_s N_s \sqrt{Z}) \end{bmatrix} \times 14.8 \text{ MeV}$ <br>
It is very interesting to note that,<br>
In comparison with SEMF, by replacing  $A_s$  with  $A$  in<br>  $(17)$ <br>
In comparison with SEMF, b  $(A_s/A)^{1-(Z/A)}$  associated with each term, binding energy of  $A$  can be estimated approximately. For relations (15) and (17), see figure 1 (red and violet curves respectively) for the estimated binding energy per nucleon close to beta stability line of  $Z=2$  to 100 compared with first four terms of the semi empirical mass formula (Green curve) where :  $a_v \approx 15.77 \text{ MeV}, a_s \approx 18.34 \text{ MeV},$  $a_a \approx 23.2$  MeV and  $a_c \approx 0.71$  MeV.  $\int_{A_4} \frac{A}{\epsilon} = A_3^{-2/3}$ <br>  $\int_{A_5} \frac{A}{\epsilon} = \left[ \frac{A_2 - A_3^{-2/3}}{A_2^{1/3}} \right] = (k^2 A_2 N_s \sqrt{Z})$ <br>  $\left[ k^2 A_3 N_s \sqrt{Z} \right]$ <br>  $= (k^2 A_4 N_s \sqrt{Z})$ <br>  $= (k^2 A_5 N$ and with SEMF, by replacing  $A_x^{(13)}$  [17)<br>
(17)<br>
(20)<br>
(17)<br>
(20)<br>
(17)<br>
(20)<br>
(17)<br>
(20)<br>
(17)<br>
(20)<br>
(17)<br>
(20)<br>
(17)<br>
(20

#### 14. Fitting medium, heavy and super heavy nuclear charge radii

For medium, heavy and super heavy atomic nuclei, nuclear charge radii [50-54] can be fitted with the following simple relation.

respectively) for the estimated binding energy per nucleon  
\nelse to beta stability line of 
$$
Z=2
$$
 to 100 compared with  
\nfirst four terms of the semi empirical mass formula (Green  
\ncurve) where :  $a_v \approx 15.77$  MeV,  $a_s \approx 18.34$  MeV,  
\n $a_d \approx 23.2$  MeV and  $a_c \approx 0.71$  MeV.  
\n14. Fitting medium, heavy and super heavy nuclear  
\ncharge radii  
\nFor medium, heavy and super heavy nuclear  
\ncharge radii  
\nFor medium, heavy and super heavy atomic nuclei, nuclear  
\ncharge radii  
\n $R_{(Z,A)} \approx \left\{Z^{1/3} + \left(\sqrt{Z(A-Z)}\right)^{1/3}\right\} \left(\frac{G_s m_p}{c^2}\right)$   
\n $\left(Z^{1/3} + \left[\frac{Z(A-Z)}{Z(A-Z)}\right]^{1/6}\right) \left(\frac{G_s m_p}{c^2}\right)$   
\n $\left(Z^{1/3$ 

See table-2. It may be noted that, this relation is free from arbitrary numbers and can be compared with the following relation available in recent literature [52].

$$
R_{(Z,N)} \cong \left\{ 1 - 0.349 \left( \frac{N - Z}{N} \right) \right\} N^{\frac{1}{3}} \quad 1.262 \text{ fm} \tag{19}
$$

#### 15. To understand neutron star mass and radius

A) If  $(M_{NS}, m_n)$  represent the masses of neutron star [33] and neutron, then,

$$
\frac{G_N M_{NS} m_n}{\hbar c} \approx \sqrt{\frac{G_s}{G_N}} \to M_{NS} \approx 3.175 M_{\odot}
$$
 (20)

B) If  $R_{NS}$  represents the neutron star radius, then,

$$
\frac{R_{NS}}{\left(\sqrt{G_s \hbar/c^3}\right)} \approx \sqrt{\frac{G_s}{G_N}} \to R_{NS} \approx 8.06 \text{ km}
$$
 (21)

#### 16. To understand earth's magnetic dipole moment

Planet's earth's magnetic dipole moment can be understood with:

$$
\mu_{earth} \approx \left(\frac{\mu_{proton}}{\mu_{electron}}\right) * \left(\frac{eG_s M_{earth}}{2c}\right) \approx 8.15 \times 10^{22} \text{ J.Tesla}^{-1}
$$
\n(22)

It is very interesting to note that,

$$
\left(\frac{\mu_{proton}}{\mu_{electron}}\right) \approx \frac{G_s m_p m_e}{\hbar c}.
$$
 Based on this observation,  

$$
\mu_{earth} \approx \left(\frac{G_s m_p m_e}{\hbar c}\right) * \left(\frac{eG_s M_{earth}}{2c}\right)
$$

$$
\approx 8.566 \times 10^{22} \text{ J.Tesla}^{-1}
$$
(23)

Based on this relation, other solar planets, exo-planets and neutron star's "mass dependent" magnetic dipole moments can be estimated [48]. See table-3. It may be noted that, for 30 hot Jupiters, on an average, estimated value is roughly 0.2 times the reference value.

#### 17. Fitting the Newtonian gravitational constant and proton-electron mass ratio

It is noticed that,

$$
\frac{G_N}{G_s} \approx \left\{ \left( \frac{m_e}{m_p} \right)^{10} \left( \frac{G_s m_p^2}{\hbar c} \right) \right\} \text{ and}
$$
\n
$$
G_N \approx \left\{ \left( \frac{m_e}{m_p} \right)^{10} \left( \frac{G_s m_p^2}{\hbar c} \right) \right\} G_s \tag{24}
$$
\n
$$
\approx 6.67986 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}
$$

By considering  $(G_s, G_N)$  as basic features of final unification and by considering  $m_{npl} \approx \sqrt{\hbar c/G_s} \approx 546.6 \text{ MeV}/c^2$  as a characteristic hadronic fermion, it is possible to show that,

$$
\left(\frac{m_p}{m_e}\right) \cong \left(\frac{G_s m_e^2}{G_N m_{npl}^2}\right)^{\frac{1}{10}} \cong 1836.3 \text{ and}
$$
\n
$$
m_p \cong \left(\frac{G_s m_e^2}{G_N m_{npl}^2}\right)^{\frac{1}{10}} m_e \cong 1836.3 \times m_e
$$
\n(25)

Here, interesting point to be noted is that, in RHS,  $m_e$ seems to be associated with  $G_s$  and  $m_{npl}$  seems to be associated with  $G_N$ .

18. Alternative expression for strong coupling constant It is also noticed that,

viXra:1710.0212v2 Dedicated to Dr. Abdus Salam  $\begin{array}{|l|} \hline \end{array}$  5

$$
\exp\sqrt{\frac{e^2}{4\pi\varepsilon_0 G_s m_p m_e}} \approx \frac{1}{\alpha_s} \approx 8.43562
$$
\n
$$
\rightarrow \alpha_s \approx 0.11855
$$
\n(26)

This can be compared with the recommended value of 0.1186 and can be given some consideration. 19. Understanding the total energy of electron in

#### hydrogen atom

Let  $r_0$  be the characteristic imaginary distance between proton and electron. In a quantum gravitational approach,

$$
r_o \approx \left(\frac{\hbar c}{G_s m_p m_e}\right) \left(\frac{\hbar}{\sqrt{m_p m_e c}}\right)
$$
  

$$
\approx \frac{\hbar^2}{G_s m_p^{3/2} m_e^{3/2}} \approx 5.6161 \times 10^{-12} \text{ m}
$$
 (27)

It may be noted that, this length is 9.42 times less than the Bohr radius  $a<sub>o</sub>$ . It can be approximated by the following relation.

$$
\frac{a_o}{r_o} \approx 2 \left( \frac{e^2}{4\pi \varepsilon_0 G_s m_p m_e} \right) \approx 2 \times 4.5475 \approx 9.075 \quad (28)
$$

By considering  $\left(\frac{1}{2n^2}\right)$  $\left(\frac{1}{2n^2}\right)$  as the probability of finding electron in its orbits labeled as  $n = 1, 2, 3, \dots$ , potential energy of electron can be understood with the following relation.

$$
r_o \approx \left(\frac{\hbar c}{G_s m_p m_e}\right) \approx 5.6161 \times 10^{-12} \text{ m}
$$
\n
$$
\approx \frac{\hbar^2}{G_s m_p^3 / 2 m_e^3} \approx 5.6161 \times 10^{-12} \text{ m}
$$
\n
$$
\approx \frac{\hbar^2}{G_s m_p^3 / 2 m_e^3} \approx 5.6161 \times 10^{-12} \text{ m}
$$
\n
$$
\approx \frac{\hbar^2}{G_s m_p^3 / 2 m_e^3} \approx 5.6161 \times 10^{-12} \text{ m}
$$
\n
$$
\approx \frac{\hbar^2}{G_s m_p^3 / 2 m_e^3} \approx 5.6161 \times 10^{-12} \text{ m}
$$
\n
$$
\approx \frac{\hbar^2}{G_s m_p^3 / 2 m_e^3} \approx 5.6161 \times 10^{-12} \text{ m}
$$
\n
$$
\approx \frac{\hbar^2}{G_s m_p m_e} \approx 2 \left(\frac{e^2}{4 \pi \epsilon_0 G_s m_p m_e}\right) \approx 2 \times 4.5475 \approx 9.075 \text{ (28)} \text{ Wg} \approx \frac{1}{4} \left(\frac{G_s m_p m_e}{2r_Z}\right) \approx (2 \times 1)^{\frac{1}{3}} \left(\frac{G_s m_e}{2r_Z}\right) \approx (2 \times 1)^{\frac{1}{3}} \left(\frac{G_s m_e}{2r_Z}\right) \approx (2 \times 1)^{\frac{1}{3}} \left(\frac{G_s m_e}{2r_Z}\right) \approx 2 \times 4.5475 \approx 9.075 \text{ (28)} \text{ Wg} \approx \frac{1}{4} \left(\frac{G_s m_p m_e}{2r_Z}\right) \approx (2 \times 1)^{\frac{1}{3}} \left(\frac{G_s m_e}{2r_Z}\right) \approx (2 \times 1)^{\frac{1}{3}} \left(\frac{G_s m_e}{2r_Z}\right) \approx (2 \times 1)^{\frac{1}{3}} \left(\frac{G_s m_e}{2r_Z}\right) \approx 2 \times 4.5475 \approx 9.075 \text{ (28)} \text{ Wg} \approx \text{m}
$$
\n
$$
\approx \left(\frac{G_s m_e m_e}{2n^2}\right) \
$$

This can be compared with  $-\left(\frac{27.2}{r^2}\right)$  eV n  $-\left(\frac{27.2}{n^2}\right)$  eV . With reference

to Virial theorem, corresponding kinetic energy can be understood with the following relation.

$$
(E_{kin})_n \approx \frac{1}{2} \left| \left( \frac{1}{2n^2} \right) \frac{G_s m_p m_e}{r_0} \right|
$$

$$
\approx \left[ \left( \frac{G_s m_p m_e}{\hbar c} \right)^2 \frac{\sqrt{m_p m_e c^2}}{4n^2} \right] \approx \frac{14.014}{n^2} \text{ eV}
$$
(30)

Thus, binding energy or total energy of electron can be understood with,

$$
(E_{tot})_n \approx (E_{pot})_n - (E_{kin})_n \approx -\left(\frac{1}{2n^2}\right) \frac{G_s m_p m_e}{2r_0}
$$
that, recommended value of  $G_N$  seems to be fitted with  
lower limit of the rms radius of proton i.e.  

$$
\approx -\left[\left(\frac{G_s m_p m_e}{\hbar c}\right)^2 \frac{\sqrt{m_p m_e c^2}}{4n^2}\right] \approx -\frac{14.014}{n^2} \text{ eV}
$$

This can be compared with the experimental total energy of

This can be completed with the experimental total energy of  
\n
$$
-\left(\frac{13.6}{n^2}\right) \text{ eV} \cdot \text{Based on this coincidence,}
$$
\n
$$
G_s \approx \left(\frac{\sqrt{2}e^2}{4\pi\epsilon_0 m_p^{5/4} m_e^{3/4}}\right) \approx 3.27125 \times 10^{28} \text{ m}^3 \text{ kg}^{-1} \text{sec}^{-2}
$$
\n(32)  
\n20. Fitting and understanding the photoelectric work functions  
\nfunctions  
\nBased on  $(G_s m_p m_e)$ , photoelectric work functions can be  
\nestimated with the following relation.  
\n
$$
W_Z \approx -Z^{\frac{1}{3}} \left(\frac{G_s m_p m_e}{2r_Z}\right) \qquad (33)
$$
\nwhere  $r_Z$  is the radius of atom and Z is the atomic  
\nnumber. With reference to light, medium and heavy atomic  
\nexperimental data range, above relation can be slightly  
\nmodified with  $A^{\frac{1}{3}}$  where A is the atomic mass number.  
\n
$$
W_Z \approx -A^{\frac{1}{3}} \left(\frac{G_s m_p m_e}{2r_Z}\right) \approx (2Z)^{\frac{1}{3}} \left(\frac{G_s m_p m_e}{2r_Z}\right) \qquad (34)
$$
\nSee table-4.  
\n21. To estimate  $G_N$  with RMS radius of proton

# 20. Fitting and understanding the photoelectric work functions

Based on  $(G_s m_p m_e)$ , photoelectric work functions can be estimated with the following relation.

$$
W_Z \approx -Z^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right) \tag{33}
$$

where  $r_Z$  is the radius of atom and Z is the atomic number. With reference to light, medium and heavy atomic experimental data range, above relation can be slightly modified with  $A^3$  where  $A$  is the atomic mass number. 1

$$
W_Z \approx -A^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right) \approx (2Z)^{\frac{1}{3}} \left( \frac{G_s m_p m_e}{2r_Z} \right) \tag{34}
$$

See table-4.

# 21. To estimate  $G_N$  with RMS radius of proton

Correlating elementary physical constants of different areas of physics is interesting and uncertain. With reference to the above semi empirical relations, in an optimistic approach, we tried to correlate the root mean square radius of proton and the Newtonian gravitational constant in the following way.

$$
R_p \cong \ln\left(\frac{e^2}{4\pi\epsilon_0 G_N m_p^2}\right) \sqrt{\frac{e^2}{4\pi\epsilon_0 G_s m_p^2}} \left(\frac{\hbar}{m_p c}\right) \cong 0.87 \text{ fm}
$$
\n(35)

elation.   
\n
$$
\frac{a_o}{r_o} \approx 2 \left( \frac{e^2}{4\pi \epsilon_0 G_s m_p m_e} \right) \approx 2 \times 4.5475 \approx 9.075 \text{ (28)}
$$
\n
$$
W_Z \approx -A \frac{1}{3} \left( \frac{G_s m_p m_e}{2r_Z} \right) \approx (22) \frac{1}{3} \left( \frac{G_s m_p m_e}{2r_Z} \right) \approx (22) \frac{1}{3} \left( \frac{G_s m_p m_e}{2r_Z} \right) \quad (34)
$$
\nSec table-4.

\nby considering  $\left( \frac{1}{2r^2} \right)$  as the probability of finding  
\nelectron in its orbits labeled as  $n = 1, 2, 3, \ldots$ , potential  
\nenergy of electron can be understood with the following  
\ntheory gives its interesting elementary physical constants of different areas  
\nrelation.

\n(*E<sub>pot</sub>*)<sub>n</sub> 
$$
\approx -\left( \frac{1}{2n^2} \right) \frac{G_s m_p m_e}{r_0}
$$
\n
$$
\approx -\left( \frac{G_s m_p m_e}{\hbar c} \right)^2 \frac{\sqrt{m_p m_e c^2}}{2n^2} \approx -\frac{28.03}{n^2} \text{ eV}
$$
\nThis can be compared with  $-\left( \frac{27.2}{n^2} \right)$  eV. With reference

\nWiral theorem, corresponding kinetic energy can be

\n(*E<sub>kin</sub>*)<sub>n</sub> 
$$
\approx \frac{1}{2} \left( \frac{1}{2n^2} \right) \frac{G_s m_p m_e}{r_0}
$$
\n
$$
\approx \left[ \left( \frac{G_s m_p m_e}{\hbar c} \right)^2 \frac{\sqrt{m_p m_e c^2}}{4n^2} \right] \approx \frac{28.03}{r_0}
$$
\n(*E<sub>kin</sub>*)<sub>n</sub> 
$$
\approx \frac{1}{2} \left( \frac{1}{2n^2} \right) \frac{G_s m_p m_e}{r_0}
$$
\n(*E<sub>kin</sub>*)<sub>n</sub> 
$$
\approx \frac{1}{2} \left( \frac{1}{2n^2} \right) \frac{G_s m_p m_e}{r_0
$$

PDG recommended [55] value of RMS radius of proton is,  $R_p \approx (0.8751 \pm 0.0061)$  fm. It is very interesting to note that, recommended value of  $G_N$  seems to be fitted with lower limit of the rms radius of proton i.e.  $R_p \approx (0.8751 - 0.0061) \approx 0.869$  fm. From relation (36), estimated  $G_N \approx 7.2092 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}$ .

Considering the recommended value of  $G_N \approx 6.67408 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}$ , from relation (35), estimated  $R_p \approx 0.8698073$  fm.

Based on the values estimated from relations (35) and (36), it is possible to say that, there exists a very tight correlation in between Newtonian gravitational constant and RMS radius of proton.

Alternatively, independent of proton rest mass and reduced Planck's constant, we noticed that,

$$
2\pi R_p \approx 2\ln\left(\frac{G_s}{G_N}\right) \sqrt{\frac{e^2 G_s}{4\pi \varepsilon_0 c^4}} \approx 0.8746 \text{ fm}
$$
  

$$
\rightarrow R_p \approx \frac{1}{\pi} \ln\left(\frac{G_s}{G_N}\right) \sqrt{\frac{e^2 G_s}{4\pi \varepsilon_0 c^4}}
$$
(37)

Based on this relation,

$$
\ln\left(\frac{G_s}{G_N}\right) \cong \pi \sqrt{\frac{4\pi \varepsilon_0 c^4 R_p^2}{e^2 G_s}} \text{ and}
$$
\n
$$
G_N \cong \left\{ \exp\left(\pi \sqrt{\frac{4\pi \varepsilon_0 c^4 R_p^2}{e^2 G_s}}\right)\right\}^{-1} G_s
$$
\n(38)

If, recommended  $R_p \approx 0.8751$  fm, estimated  $G_N \approx 6.3793785 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}.$ 

# 22. To understand neutron life time controversy and to fit  $G_N$

With reference to  $(Gs, G_N)$  and with reference to bottle experiments and beam experiments, it is also possible to express  $t_n$  in the following way.

To fit with bottle experiments,

$$
(t_n)_{bottle} \approx \left(\frac{G_s}{G_N}\right)^{\frac{2}{3}} \sqrt{\frac{4\pi\varepsilon_0 G_s m_p^2}{e^2}} \left(\frac{\hbar(1-\alpha)}{m_n c^2}\right) \tag{39}
$$

$$
\approx 878.985 \text{ sec}
$$

To fit with beam experiments,

$$
(t_n)_{beam} \approx \left(\frac{G_s}{G_N}\right)^{\frac{2}{3}} \sqrt{\frac{4\pi\varepsilon_0 G_s m_p^2}{e^2}} \left(\frac{\hbar}{m_n c^2}\right)
$$
\n
$$
\approx 885.45 \text{ sec}
$$
\n(40)

 Interesting point to be noted is that, results of bottle experiments and beam experiments can be correlated with a factor of the kind,  $(1 - \alpha)$ .

$$
\frac{(t_n)_{bottle}}{(t_n)_{beam}} \approx (1 - \alpha)
$$
\n(41)\n  
\n(39) and (40), magnitude of  $G_N$  can be following relations.

Based on relations (39) and (40), magnitude of  $G_N$  can be estimated with the following relations.

$$
\frac{(t_n)_{bottle}}{(t_n)_{beam}} \approx (1-\alpha) \qquad (41)
$$
\nused on relations (39) and (40), magnitude of  $G_N$  can be  
\ntimated with the following relations.

\n
$$
G_N \approx \left\{ \sqrt{\frac{4\pi\epsilon_0 G_s m_p^2}{e^2} \left( \frac{\hbar}{m_n c^2 (t_n)_{bottle}} \right) \right\}^{\frac{3}{2}} G_s \qquad (42)
$$
\n
$$
\approx 6.753392 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}.
$$
\n
$$
G_N \approx \left\{ \sqrt{\frac{4\pi\epsilon_0 G_s m_p^2}{e^2} \left( \frac{\hbar}{m_n c^2 (t_n)_{beam}} \right) \right\}^{\frac{3}{2}} G_s \qquad (43)
$$
\n
$$
\approx 6.648678 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}.
$$
\nused on relations (11) and (39),

\n
$$
G_N \approx \left\{ \sqrt{\frac{e^2}{4\pi\epsilon_0 G_s m_p^2}} \times \exp \left( \frac{E_Y}{(m_n - m_p) c^2} \right) \right\}^{-\frac{3}{2}} G_s
$$
\n
$$
\approx 6.76075 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}.
$$
\n(44)

Based on relations (11) and (39),

$$
G_N \cong \left\{ \sqrt{\frac{e^2}{4\pi\varepsilon_0 G_s m_p^2}} \times \exp\left(\frac{E_Y}{(m_n - m_p)c^2}\right) \right\}^{-\frac{3}{2}} G_s
$$
  
\n
$$
\cong 6.76075 \times 10^{-11} \text{ m}^3 \text{kg}^{-1} \text{sec}^{-2}.
$$
\n(44)

# 23. Discussion

sed on this relation,<br>  $\ln\left(\frac{G_x}{G_N}\right) = \pi \sqrt{\frac{4\pi \epsilon_0 c^4 R_p^2}{e^2 G_s}}$  and<br>  $G_N \approx \left\{ \exp\left(\pi \sqrt{\frac{4\pi \epsilon_0 c^4 R_p^2}{e^2 G_s}}\right)\right\}^{-\frac{1}{2}} G_s$ <br>  $\qquad G_N \approx \left\{ \exp\left(\pi \sqrt{\frac{4\pi \epsilon_0 c^4 R_p^2}{e^2 G_s}}\right)\right\}^{-1}$ <br>  $\qquad G_N \approx \left\{ \sqrt{\frac{4\pi \epsilon_0 c^4 R_p^$ Fig. (a)  $G_X = 6.3793785 \times 10^{-11}$  m<sup>3</sup> kg<sup>-1</sup>sec<sup>2</sup>.<br>
2. **To understand neutron life time controversy and to**<br>
2. **To understand neutron life time controversy and to**<br>
2. **To understand neutron life time controversy and t** It is true that, unless stringent requirements are met, in general, speculative alternatives to currently accepted theories cannot be accepted. Scientific papers having content that lie outside the mainstream of current research must justify by including a clear, detailed discussion of the motivation for the new speculation, with reasons for introducing any new concepts. If the new formulation results are in contradiction with the accepted theory, then there must both be a discussion of which experiments could be done to verify that the conventional theory needs improvement, and also an analysis showing the consistency of the new theory with the existing experiments. In this context, we would like to appeal that, in this paper, we presented a variety of relations pertaining to nuclear and electroweak coupling constants. It is clear from the above relations that we could satisfactorily fit the nuclear data through semi-empirical relations. This sincere attempt is to be ascertained by the scientific community. The problem is with "our understanding" and "our perception" by using which the current 'scientific standards' and 'procedures' can be reviewed for a better understanding of nature. We would like to appeal that,

- 1) With respect to currently believed String theory and Quantum gravity models - proposed assumption and proposed semi empirical relations, can be given some consideration in developing a 'workable model' of TOE.
- 2) Magnitude of  $G_N$  can be estimated from microscopic elementary constants by considering expressions like

viXra:1710.0212v2 Dedicated to Dr. Abdus Salam  $\begin{array}{|c|c|c|c|c|c|}\n\hline\n\end{array}$ 

$$
\left(\frac{G_s}{G_N}\right)
$$
 or  $\ln\left(\frac{G_s}{G_N}\right)$  or  $\ln\left(\frac{e^2}{4\pi\varepsilon_0 G_N m_p^2}\right)$ . We are

working in this new direction.

3) Magnitude of  $G<sub>s</sub>$  can be estimated with any of the following three relations:

$$
\left(\frac{G_s}{G_N}\right)
$$
 or  $\ln\left(\frac{G_s}{G_N}\right)$  or  $\ln\left(\frac{e^2}{G_N}\right)$  or  $\ln\left(\frac{e^2}{4\pi\epsilon_0 G_N m_p^2}\right)$ . We are  
working in this new direction.  
Magnitude of  $G_s$  can be estimated with any of the  
following three relations:  
following three relations:  

$$
\left(\frac{m_p}{m_e}\right)\left(\frac{e^2}{4\pi\epsilon_0}\right)\left(G_s m_p^2\right)
$$
  

$$
h = \sqrt{\frac{m_p}{m_e c^2}}\left(\frac{m_p}{m_{\pi c_0}}\right)\left(\frac{e^2}{4\pi\epsilon_0}\right)\left(\frac{e^2}{4\pi\epsilon_0}\right)
$$
  

$$
h = \sqrt{\frac{4\epsilon^2 G_S m_p^3}{4\pi\epsilon_0 \hbar^2 m_e c^2}}\right)
$$
  

$$
= \ln\left(\frac{m_n c^2}{\hbar}\right) = \ln\left(\frac{(m_n c^2) t_n}{\hbar}\right)
$$
  

$$
= \ln\left(\frac{m_n c^2}{\hbar}\right)
$$
  

$$
\left(\frac{m_n c^2 - m_p c^2}{4\pi\epsilon_0 \hbar^2 m_e c^2}\right) \equiv \ln\left(\frac{(m_n c^2) t_n}{\hbar}\right)
$$
  

$$
= \ln\left(\frac{m_n c^2}{\hbar}\right)
$$
  

$$
= \ln\left(\frac
$$

# 24. Conclusion

We would like to appeal the science community that:

- 1) So far, the whole subject of nuclear physics and particle physics is being studied independent of 'gravity'.
- 2) Background of whole experimental apparatus of nuclear and particle physics is 'gravity' only.
- As of today, 'string theory and its sister models' seem to be completely theoretical in nature and beyond the scope of observed four dimensions.
- 4) To have a 'theory of everything, it is inevitable to unite gravity and other three atomic interactions and is beyond the scope of current experimental physics.

Based on these points, even though it is in its budding stage, in a broad view, our work can be recommended for further research.

Acknowledgements: Author Seshavatharam U.V.S is indebted to professors brahmasri Dr. M. Nagaphani Sarma, Chairman, shri K.V. Krishna Murthy, former chairman, Institute of Scientific Research in Vedas (I-SERVE), Hyderabad, India and Shri K.V.R.S. Murthy, former scientist IICT (CSIR), Govt. of India, Director, Research and Development, I-SERVE, for their valuable guidance and great support in developing this subject.

#### References

- [1] Seshavatharam, U. V. S. & Lakshminarayana, S., On the Basics of Nuclear Quantum Gravity. Prespacetime Journal. (2017). 8(6): 777-791.
- [2] Chowdhury, P.R. et al. Modified Bethe-Weizsacker mass formula with isotonic shift and new driplines. Mod. Phys. Lett. A20 p.1605-1618. (2005).
- [3] N. Ghahramany et al. New scheme of nuclide and nuclear binding energy from quark-like model. Iranian Journal of Science & Technology. 2011; A3: 201-208.
- [4] N.Ghahramany et al. New approach to nuclear binding energy in integrated nuclear model. Journal of Theoretical and Applied Physics 6:3 (2012).
- [5] U. V. S Seshavatharam, S. Lakshminarayana. Simplified Form of the Semi-empirical Mass Formula.

Prespacetime Journal. (2017). 8(7): 881-810.

- [6] U. V. S Seshavatharam, S. Lakshminarayana. On the role of strong coupling constant and nucleons in understanding nuclear stability and binding energy. Journal of Nuclear Sciences. (2017); 4(1): 7-18.
- $\frac{e^2}{\sqrt{2\epsilon_0}G_Nm_\rho^2}$ . Prespacetime Journal. (2017). 8(7): 881-810.<br>  $\frac{e^2}{\sqrt{2\epsilon_0}G_Nm_\rho^2}$ . We are [6] U. V. S. Seshavatharam, S. Lakshminarayana. On the<br>
role of strong coupling constant and nucleons in<br>
unders [7] U. V. S. Seshavatharam, Lakshminarayana S. Understanding Nuclear Stability, Binding Energy and Magic Numbers with Fermi Gas Model. Journal of Applied Physical Science International, 4 (2) pp.51-59 (2015).
	- [8] U. V. S. Seshavatharam and S. Lakshminarayana. Understanding nuclear stability and binding energy with nuclear and electromagnetic gravitational constants. 61st DAE-BRNS Symposium on Nuclear Physics. A136. 332
	- [9] U. V. S. Seshavatharam and S. Lakshminarayana. On the role of strong interaction in understanding nuclear beta stability line and nuclear binding energy. Proceedings of the DAE-BRNS Symp. On Nucl. Phys. 60 (2015) A32 ,118.
	- [10] K. Becker, M. Becker and J. H. Schwarz. String Theory and M-theory: A Modern Introduction. Cambridge University Press, (2006).
	- [11] Edward Witten. What Every Physicist Should Know About String Theory. GR Centennial Celebration, Strings 2015, Bangalore, India. (2015).
	- [12] Juan M. Maldacena. Gravity, Particle Physics and Their Unification. Int.J.Mod.Phys. A15S1 840-852 (2000)
	- [13] Tilman Sauer. Einstein's Unified Field Theory Program. The Cambridge Companion to Einstein, M. Janssen, C. Lehner (eds), Cambridge University Press. 2013. (Chapter 9) http://philsciarchive.pitt.edu/3293/1/uft.pdf
	- [14] U. V. S. Seshavatharam and S. Lakshminarayana. Scale independent workable model of final unification. Universal journal of physics and application. 10(6): 198-206, 2016
	- [15] U. V. S. Seshavatharam and S. Lakshminarayana. Towards a workable model of final unification. International Journal of Mathematics and Physics 7, No1,117-130 (2016).
	- [16] U. V. S. Seshavatharam and S. Lakshminarayana. To Validate the Role of Electromagnetic and Strong Gravitational Constants via the Strong Elementary Charge. Universal Journal of Physics and Application 9(5): 216-225, 2015
	- [17] U. V. S. Seshavatharam and Lakshminarayana S. To confirm the existence of nuclear gravitational constant, Open Science Journal of Modern Physics. 2(5): 89-102 (2015)
	- [18] U. V. S. Seshavatharam and Lakshminarayana S. Final unification with Schwarzschild's Interaction. Journal of Applied Physical Science International 3(1): 12-22 (2015).
	- [19] Seshavatharam, U. V. S. and Lakshminarayana, S. The Possible Role of Newtonian, Strong & Electromagnetic Gravitational Constants in Particle Physics. Prespacetime journal, Vol 7, Issue 5, pp. 857-888 (2016)
	- [20] U. V. S. Seshavatharam and S. Lakshminarayana. Final unification with three gravitational constants associated with nuclear, electromagnetic and gravitational interactions. International Journal of Advanced Astronomy. 4 (2) 105-109 (2016)
	- [21] U. V. S. Seshavatharam and S. Lakshminarayana. Strong Nuclear Gravitational Constant

viXra:1710.0212v2 Dedicated to Dr. Abdus Salam 8

and the Origin of Nuclear Planck Scale. Progress in Physics. Vol -3, 31-38, July 2010.

- [22] U. V. S. Seshavatharam et al. Understanding the constructional features of materialistic atoms in the light of strong nuclear gravitational coupling. Materials Today: 3/10PB, Proceedings 3 (2016) pp. 3976-3981
- [23] U. V. S. Seshavatharam and S. Lakshminarayana. Understanding the basics of final unification with three gravitational constants associated with nuclear, electromagnetic and gravitational interactions. Journal of Nuclear Physics, Material Sciences, Radiation and Applications Vol-4, No-1, 1-19, (2017)
- [24] U. V. S. Seshavatharam and S. Lakshminarayana. Understanding nuclear structure with Schwarzschild interaction and Avogadro number. Proceedings of the DAE-BRNS Symp. On Nucl. Phys. 60 (2015). F8, 852
- [25] U.V. S. Seshavatharam and S. Lakshminarayana. Analytical estimation of the gravitational constant with atomic and nuclear physical constants. Proceedings of the DAE-BRNS Symp. On Nucl. Phys. 60 (2015) F7, 850
- [26] C. J. Isham, A. Salam and J. Strathdee, 2+ Nonet as Gauge particles for SL(6, C) symmetry, Phys. Rev. D8, 8 (1973).
- [27] A.Salam and J. Strathdee, Class of solutions for the strong-gravity equations, Phys. Rev. D16, 8 (1977).
- [28] Sivaram and K. P. Sinha, Strong Spin-two Interaction and General Relativity, Phys. Rep. (Rev. Sect.Phys. Lett.) 51, 3 (1979).
- [29] Dj. Sijacki and Y. Ne'eman, QCD as an effective strong gravity, Phys. Lett. B247, 4 (1990).
- [30] Y. Ne'eman and Dj. Sijacki, Proof of pseudo-gravity as QCD approximation for the hadron IR region and J\_M2 Regge trajectories, Phys. Lett. B276 (1992).
- [31] Usha Raut and K. P. Sinha. Strong gravity and the Yukawa field. International Journal of Theoretical Physics, Vol20, Issue 1, pp 69-77, (1981)
- [32] Salam A, Sivaram C. Strong Gravity Approach to QCD and Confinement. Mod. Phys. Lett, A8(4), 321- 326. (1993)
- [33] O. F. Akinto, Farida Tahir. Strong Gravity Approach to QCD and General Relativity. arXiv:1606.06963v3
- [34] Roberto Onofrio. Proton radius puzzle and quantum gravity at the Fermi scale. EPL 104, 20002 (2013)
- [35] S. Bethke and G.P. Salam. Quantum chromodynamics. K.A. Olive et al. (Particle Data Group), Chin. Phys. C, 38, 090001 (2014) and 2015 update.
- [36] Helmut Satz. The Quark-Gluon Plasma.Nucl.Phys.A862-863:4-12,2011
- [37] A.V. Manohar, C.T. Sachrajda. Quark masses. C. Patrignani et al. (Particle Data Group), Chin. Phys. 2016; C( 40), 100001.
- [38] David J. Gross. Twenty Five Years of Asymptotic

Freedom. Nucl.Phys.Proc.Suppl. 74 (1999) 426-446

- [39] Weizsäcker, Carl Friedrich von, On the theory of nuclear masses; Journal of Physics 96 (1935) pages 431- 458.
- [40] W. D. Myers et al. Table of Nuclear Masses according to the 1994 Thomas-Fermi Model.(from nsdssd.lbl.gov)
- [41] J.A. Maruhn et al., Simple Models of Many-Fermion Systems, Springer-Verlag Berlin Heidelberg 2010. Chapter 2, page:45-70.
- [42] Planck, M. (1914). The Theory of Heat Radiation. Masius, M. (transl.) (2nd ed.). P.Blakiston's Son & Co.
- [43] S. HAŁAS.100 years of work function. Materials Science-Poland, Vol. 24, No. 4, 2006
- [44] Deaver, Bascom; Fairbank, William (July 1961). Experimental Evidence for Quantized Flux in Superconducting Cylinders. Physical Review Letters. 7(2): 43-46.
- [45] The Periodic Table of the Elements (including Atomic Radius)www.sciencegeek.net/tables/AtomicRadius.pd f
- [46] S. Hawking, Particle Creation by Black Hole, Commun. Math. Phys. 43, 199-220(1975)
- [47] Sebastien Guillot et al. Measurement of the Radius of Neutron Stars with High S/N Quiescent Low-mass Xray Binaries in Globular Clusters. Astrophys. J. 772  $(2013)$
- [48] Hector Javier Durand-Manterola. Dipolar Magnetic Moment of the Bodies of the Solar System and the Hot Jupiters. Planetary and Space Science 57:1405-1411 (2009)
- [49] C. L. Morris et al. A new method for measuring the neutron lifetime using an in situ neutron detector. Report number: LA-UR-16- 27352.https://arxiv:1610.04560.
- [50] I.Angeli, K.P. Marinovab. Table of experimental nuclear ground state charge radii: An update. Atomic Data and Nuclear Data Tables 99 (2013) 69–95.
- [51] I. Angeli. A consistent set of nuclear rms charge radii: properties of the radius surface R(N,Z). Atomic Data and Nuclear Data Tables 87 (2004) 185–206
- [52] T. Bayram et .al. New parameters for nuclear charge radius formulas. ACTA PHYSICA POLONICA B. Vol. 44,No 8, 1791-1799 (2013)
- [53] Ning Wang and Tao Li. Shell and isospin effects in nuclear charge radii. Phys.Rev.C88:011301(R),2013
- [54] Zhongzhou Ren. Nuclear charge radii of exotic nuclei and superheavy nuclei from experimental decay data. http://cyclotron.tamu.edu/she2015/assets/pdfs/presenta tions/Ren\_SHE\_2015\_TAMU.pdf
- [55] C. Patrignani et al. (Particle Data Group), Chin. Phys. C, 40, 100001 (2016) and 2017 update

	Proton number	<b>Mass</b> number	<b>Neutron</b> number	<b>Estimated</b> binding energy <b>Relation (15)</b> (MeV)	(1 to 4) terms of SEMF binding energy (MeV)
		4		14.1	16.0
		$\sigma$		31.8	31.7
				49.8	48.5
viXra:1710.0212v2				<b>Dedicated to Dr. Abdus Salam</b>	

Table 1: To estimate nuclear binding energy close to beta stability line



viXra:1710.0212v2 Dedicated to Dr. Abdus Salam  $10$ 











60	143	83	4.9893	4.9725
61	146	85	5.0200	5.0020
62	149	87	5.0503	5.0312
63	151	88	5.0753	5.0568
64	154	$\overline{90}$	5.1050	5.0853
65	157	92	5.1343	5.1136
66	160	94	5.1633	5.1416
67	163	96	5.1919	5.1693
68	166	98	5.2202	5.1968
69	168	99	5.2436	5.2207
70	171	101	5.2714	5.2476
71	174	103	5.2988	5.2743
72	177	105	5.3260	5.3007
73	180	107	5.3529	5.3269
74	183	109	5.3795	5.3528
75	186	111	$\frac{1}{5.4058}$	5.3785
76	189	113	5.4319	5.4040
77	192	115	5.4577	5.4293
78	195	117	5.4833	$5.\overline{4544}$
79	198	119	5.5086	5.4792
80	201	121	5.5337	5.5038
81	204	123	5.5586	5.5282
82	207	$\overline{125}$	5.5832	5.5524
83	210	127	5.6077	5.5765
84	213	129	5.6319	5.6003
85	216	131	5.6558	5.6239
86	219	133	5.6796	5.6474
87	222	135	5.7032	5.6706
88	226	138	5.7302	5.6969
89	229	140	5.7534	5.7198
90	232	142	5.7763	5.7425
91	235	144	5.7991	5.7651
92	238	146	5.8217	5.7875
93	241	148	5.8442	5.8097
94	245	$\overline{151}$	5.8698	5.8349
95	248	153	5.8919	5.8568
96	251	155	5.9138	5.8785
97	254	157	5.9355	5.9001
98	257	159	5.9570	5.9215
99	261	162	5.9817	5.9459
100	264	164	6.0029	5.9670

Table-3: To fit the mass dependent magnetic dipole moments of hot Jupiters



viXra:1710.0212v2 Dedicated to Dr. Abdus Salam  $\boxed{13}$ 



# Table-4: To fit and estimate the photo electric work functions



viXra:1710.0212v2 Dedicated to Dr. Abdus Salam  $\boxed{14}$ 

