

An Essay on the Zeroes of an L-function

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We set the following notation.

- K a global field
- K_ν a local field, completion of K at the place ν of K
- A_K the adèle ring of K
- C_K the idele class group $GL_1(A_K)/K^*$
- \hat{C}_K the dual group of C_K .

0.

We will summarize the spectral interpretation of critical zeros of $L(\chi, s)$ associated χ of C_K by Alain Connes. Let h be a test function. The Weil explicit formula says

$$\sum_{\nu} \int_{K_{\nu}^*} \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu = \hat{h}(0) + \hat{h}(1) - \sum_{L(\chi, \rho) = 0} \hat{h}(\chi, \rho).$$

Suppose that there exists a representation U of C_K , and that

$$\mathrm{tr} U(h) = \sum_{\nu} \int_{K_{\nu}^*} \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu$$

is satisfied. We see that

$$\mathrm{tr} U(h) = \hat{h}(0) + \hat{h}(1) - \sum_{L(\chi, \rho) = 0} \hat{h}(\chi, \rho)$$

holds. We can say that critical zeros of $L(\chi, s)$ appear as the spectra of the operator U . It is just *the spectral interpretation of critical zeros of $L(\chi, s)$* .

Let

$$X = A_K/K^*.$$

The left regular representation U of C_K on $L^2_\delta(X)$ which is a weighted L^2 space can be used to accomplish our task. Namely, it holds that

$$\text{tr}U(h) = \hat{h}(0) + \hat{h}(1) - \sum_{\substack{L(\chi, \rho) = 0 \\ \text{Re } \rho = 1/2}} \hat{h}(\chi, \rho) + \infty h(1).$$

However we will not try to treat the representation $(U, L^2_\delta(X))$ directly. Instead of the representation $(U, L^2_\delta(X))$, we will think of the operator $Q_\Lambda U$ where U is the left regular representation of C_K on $L^2(X)$. Because, firstly there is a possibility of using some results to compute $\text{Trace } Q_\Lambda U$, secondly we can eliminate the parameter δ of $L^2_\delta(X)$. Now, we can show that

$$\text{Trace } Q_\Lambda U(h) = \hat{h}(0) + \hat{h}(1) - \sum_{\substack{L(\chi, \rho) = 0 \\ \text{Re } \rho = 1/2}} \hat{h}(\chi, \rho) + \infty h(1) \quad \Lambda \rightarrow \infty$$

for the function h which belongs to *Bruhat–Shwartz space* $\mathcal{S}(C_K)$ of functions on C_K .

We try to compute $\text{Trace } Q_\Lambda U(h)$. This has the relationship to the validity of the Riemann Hypothesis. Suppose that we can compute as follows;

$$\text{Trace } Q_\Lambda U(h) = 2h(1)\log' \Lambda + \sum_v \int_{K_v^*}' \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu + o(1) \quad \Lambda \rightarrow \infty$$

where $2\log' \Lambda = \int_{\lambda \in C_K, |\lambda| \in [\Lambda^{-1}, \Lambda]} d^* \lambda$. We obtain a trace formula:

$$\hat{h}(0) + \hat{h}(1) - \sum_{\substack{L(\chi, \rho) = 0 \\ \text{Re } \rho = 1/2}} \hat{h}(\chi, \rho) + \infty h(1) = 2h(1)\log' \Lambda + \sum_v \int_{K_v^*}' \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu + o(1) \quad \Lambda \rightarrow \infty.$$

The left side is spectral and the right side is geometrical. From the Weil explicit formula, we have seen that

$$\sum_v \int_{K_v^*}' \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu = \hat{h}(0) + \hat{h}(1) - \sum_{L(\chi, \rho) = 0} \hat{h}(\chi, \rho).$$

Therefore, one obtains that

$$\sum_{L(\chi, \rho) = 0} \hat{h}(\chi, \rho) = \sum_{\substack{L(\chi, \rho) = 0 \\ \operatorname{Re} \rho = 1/2}} \hat{h}(\chi, \rho).$$

It means the validity of the Riemann Hypothesis. Conversely, the validity of the Riemann Hypothesis implies that

$$\operatorname{Trace} Q_{\Lambda} U(h) = 2h(1) \log' \Lambda + \sum_{\nu} \int_{K_{\nu}^*}' \frac{h(\mu^{-1})}{|1 - \mu|} d^* \mu + o(1) \quad \Lambda \rightarrow \infty.$$

1.

We try to characterize L-functions from the view of the representation theory.

We will begin with the local case. Denote the set of the irreducible representations of K_ν^* by $\text{Irr}(K_\nu^*)$. Let (π_ν, V_{π_ν}) be an irreducible representation of K_ν^* . Put

$$\pi_\nu(f)v = \int_{K_\nu^*} f(g)\pi_\nu(g)v d^*g, \quad f \in \mathcal{S}(K_\nu).$$

Suppose that $\text{tr}\pi_\nu(f)$ can be defined, namely $\pi_\nu(f)$ is a trace class operator. So we may think that there exists a character $\text{tr}\pi_\nu$ of K_ν^* , and

$$\text{tr}\pi_\nu(f) = \int_{K_\nu^*} f(g)\text{tr}\pi_\nu(g) d^*g.$$

Define the local zeta function as

$$Z(s, \chi, \Phi) = \int_{K_\nu^*} \Phi(g)\chi(g)|g|^s d^*g.$$

Here $s \in \mathbb{C}$, χ is a character of K_ν^* and $\Phi \in \mathcal{S}(K_\nu)$. The integral converges at $\text{Re}(s) > 0$. The L-factor $L(s, \chi)$ is defined as $Z(s, \chi, \Phi)/L(s, \chi)$ being entire. We will see that the local zeta function associated with π_ν can be

$$Z(s, \text{tr}\pi_\nu, \Phi) = \int_{K_\nu^*} \Phi(g)\text{tr}\pi_\nu(g)|g|^s d^*g.$$

The L-factor $L(s, \pi_\nu)$ is defined as $Z(s, \text{tr}\pi_\nu, \Phi)/L(s, \pi_\nu)$ being entire.

Next, we will think of the global case. It is performed on the adèle ring of K . Set

$$\pi = \otimes_{\nu} \pi_{\nu}, \quad V_{\pi} = \otimes_{\nu} V_{\pi_{\nu}}.$$

We can obtain an irreducible representation (π, V_{π}) of A_K^* . Denote the set of the irreducible representations of A_K^* by $\text{Irr}(A_K^*)$. Suppose that $\pi(f)$ where $f \in \mathcal{S}(A_K)$ is a trace class operator. Then $\text{tr}\pi$ is given as a character of A_K^* . We also obtain the global zeta function

$$Z(s, \text{tr}\pi, \Phi) = \prod_{\nu} Z(s, \text{tr}\pi_{\nu}, \Phi).$$

Here $\Phi \in \mathcal{S}(A_K)$. We define the L-function associated with π as follows;

$$L(s, \pi) = \prod_{\nu} L(s, \pi_{\nu}).$$

Each L-factor $L(s, \pi_\nu)$ gives the Euler factor of $L(s, \pi)$, namely $L(s, \pi)$ has the Euler product. The $L(s, \pi)$ satisfies the functional equation which is given by the functional equation of the global zeta function. Thus, $L(s, \pi)$ is analytically continued to the function which is meromorphic in the whole plain \mathbb{C} .

We shall consider an irreducible representation (π, V_π) of C_K . Let \mathcal{H}_π be a suitable completion of V_π with a certain inner product. One obtains a unitary representation (π, \mathcal{H}_π) , which is a left regular representation of C_K on \mathcal{H}_π . We may say that if $\pi \in \text{Irr}(C_K)$ then $\pi \in \hat{C}_K$. Thus,

$$\mathcal{H} = \bigoplus_{\pi \in \hat{C}_K} \mathcal{H}_\pi, \quad \mathcal{H}_\pi = \left\{ \xi \mid \xi(g^{-1}x) = \pi(g)\xi(x), \forall g \in C_K \right\}.$$

We know that $\text{tr}\pi$ is a character of C_K . We frequently use χ to denote a character of C_K . Then, $\text{tr}\pi = \chi$. Correspondingly, $L(s, \pi) = L(s, \chi)$.

Lastly we will mention trace formulae. The trace formula which is given by a zeta function:

$$\underbrace{\dots}_{\text{Zero points}} = \underbrace{\dots}_{\text{Geometrical side}}$$

is a prototype. Selberg's trace formula is that

$$\underbrace{\dots}_{\text{Eigenvalues of Laplacian}} = \underbrace{\dots}_{\text{Geometrical side}}.$$

There exists an operator M such that it is commutative with the Laplacian of H . The operator is the integral operator which has $k(z, w)$ as an integral kernel

$$M(f)(z) = \int_H k(z, w)f(w)d\mu(w).$$

The Selberg's trace formula gives the explicit formula of Selberg's zeta function.

The trace formula given by Connes is the same type as Selberg's. It is that

$$\underbrace{\dots}_{\text{Characters}} = \underbrace{\dots}_{\text{Geometrical side}}.$$

Here $U(h): C_c^\infty(X) \rightarrow C_c^\infty(X)$

$$(U(h)\xi)(x) = \int_{C_K} h(g)(U(g)\xi)(x)d^*g.$$

The operator $U(h)$ is the integral operator which has $k_h(x, y)$ as an integral kernel

$$(U(h)\xi)(x) = \int_{C_K} k_h(x, y)\xi(y)d^*y.$$

2.

The space $\mathcal{S}(A_K)_0$ is given as the codimension 2 subspace of $\mathcal{S}(A_K)$ such that

$$f(0) = 0, \quad \int_X f(x) dx = 0.$$

Let $L^2(X)_0$ be the completion of $\mathcal{S}(A_K)_0$. We obtain an exact sequence:

$$0 \rightarrow L^2(X)_0 \rightarrow L^2(X) \rightarrow \mathbb{C} \oplus \mathbb{C}(1) \rightarrow 0$$

where $\mathbb{C} \oplus \mathbb{C}(1) \cong L^2(X)/L^2(X)_0$.

[Remark] \mathbb{C} is a trivial C_K module:

$$T(g)\lambda = \lambda \quad g \in C_K, \lambda \in \mathbb{C}.$$

$\mathbb{C}(1)$ is Tate twist:

$$T(g)\lambda = |g|\lambda \quad g \in C_K, \lambda \in \mathbb{C}.$$

Here we have to give one's attention to the space X . The space X is a *delicate quotient space*. It must be non-compact. It must be also questionable to think that X contains C_K as a subspace. However, considering the construction of $L^2(C_K)$, if we restrict the function in $L^2(X)$ to C_K then it can be a function on C_K . We can also obtain the following exact sequence:

$$0 \rightarrow L^2(X)_0 \xrightarrow{T} L^2(C_K) \rightarrow \mathcal{H} \rightarrow 0$$

where $\mathcal{H} \cong L^2(C_K)/\text{Im}(T)$. Let U be a left regular representation of C_K on $L^2(X, dx)$ and V be a left regular representation of C_K on $L^2(C_K, d^*x)$. For $f(x) \in L^2(X, dx)$, let $(Tf)(a)$ be the restriction of $f(x)$ to C_K . Then,

$$(Tf)(a) = |a|^{1/2} f(a) \quad \forall a \in C_K.$$

Since $dx = |x|d^*x$, we will understand that $(Tf)(a) \in L^2(C_K, d^*x)$. Set

$$(U(g)f)(x) = f(g^{-1}x) \quad \forall g \in C_K, x \in X.$$

It turns out that

$$\begin{aligned} T(U(g)f)(a) &= \text{the restriction of } f(g^{-1}x) \\ &= |g|^{1/2}(V(g)Tf)(a) \quad \forall a, g \in C_K. \end{aligned}$$

From this equation, it is that $|g|^{-1/2} T(U(g)f)(a) = V(g)(Tf)(a)$. For $(Tf)(a)$,

$$\begin{aligned} V(g)(Tf)(a) &= \text{the restriction of } |g|^{-1/2} f(g^{-1}x) \\ &= |a|^{1/2} |g|^{-1/2} f(g^{-1}a). \end{aligned}$$

From $f \in \mathcal{S}(A_K)$, we will see that $|g|^{-1/2} f(x) \in L^2(X)_0$, and that $|g|^{-1/2} f(g^{-1}x) \in L^2(X)_0$. Thus $V(\text{Im}(T)) \subseteq \text{Im}(T)$, namely $\text{Im}(T)$ is an invariant subspace for V . Now, we have to turn one's attention to using $L^2(C_K)$. Because C_K is abelian locally compact, we can't always decompose $L^2(C_K)$ in the direct sum of finite dimensional subspaces. This fact, $L^2(C_K)$ having no finite dimensional subrepresentation, is an obstacle to our attempt computing the trace of U .

*“The second subtle point is that since C_K is abelian and non compact, its regular representation does not contain any finite dimensional subrepresentation so that the Polya-Hilbert space cannot be a subrepresentation (or unitary quotient) of V . There is an easy way out which is to replace $L^2(C_K)$ by $L^2_\delta(C_K)$ using the polynomial weight $(\log^2|a|)^{\delta/2}$, i.e. the norm $\|\xi\|_\delta^2 = \int_{C_K} |\xi(a)|^2 (1 + \log^2|a|)^{\delta/2} d^*a$.”* in A. Connes [2].

Because $L^2_\delta(C_K)$ is a weighted L^2 space, we can decompose it in the direct sum of finite dimensional subspaces. Let the Hilbert space $L^2_\delta(X)$ ($\delta > 1$) be the space of functions on X with the square norm

$$\|f\|_\delta^2 = \int_X |f(x)|^2 (1 + (\log|x|)^2)^{\delta/2} dx.$$

The Hilbert space $L^2_\delta(C_K)$ is obtained from the space of functions with the square norm

$$\|f\|_\delta^2 = \int_{C_K} |f(g)|^2 (1 + (\log|g|)^2)^{\delta/2} d^*g$$

where we normalize the Haar measure of the multiplicative group C_K

$$\int_{|g| \in [1, \Lambda]} d^*g \sim \log \Lambda \quad \Lambda \rightarrow +\infty.$$

We understand that the representation $(V, L^2_\delta(C_K))$ isn't unitary because of the suffix $(1 + (\log|g|)^2)^{\delta/2}$. However

$$\|V(a)\|_\delta = O((\log|a|)^{\delta/2}) \quad |a| \rightarrow \infty.$$

It is also satisfied that

$$\|V(a)\|_\delta = O((\log|a|)^{\delta/2}) \quad |a| \rightarrow 0.$$

[Remark] It holds that

$$\|V(a)\|_{\delta} \leq (c \cdot (1 + (\log|a|)^2)^{\delta/2})^{1/2}.$$

Here we may say that $\|V(a)\|_{\delta} \geq 0$. We can compute as follows;

$$\|V(a)\|_{\delta}^2 \leq c \cdot (1 + (\log|a|)^2)^{\delta/2}$$

moreover,

$$\|V(a)\|_{\delta}^{4/\delta} \leq c^{4/\delta} \cdot (1 + (\log|a|)^2).$$

Thus,

$$\frac{\|V(a)\|_{\delta}^{4/\delta}}{(\log|a|)^2} \leq c^{4/\delta} \cdot \frac{1 + (\log|a|)^2}{(\log|a|)^2}.$$

It turns out that

$$\frac{\|V(a)\|_{\delta}^{4/\delta}}{(\log|a|)^2} \leq c^{4/\delta} \quad |a| \rightarrow \infty \quad \text{and} \quad \frac{\|V(a)\|_{\delta}^{4/\delta}}{(\log|a|)^2} \leq c^{4/\delta} \quad |a| \rightarrow 0.$$

We can show that

$$\frac{\|V(a)\|_{\delta}^{4/\delta}}{(\log|a|)^2} = \left(\frac{\|V(a)\|_{\delta}}{(\log|a|)^{\delta/2}} \right)^{4/\delta}.$$

Therefore,

$$\frac{\|V(a)\|_{\delta}}{(\log|a|)^{\delta/2}} \leq c \quad |a| \rightarrow \infty \quad \text{and} \quad \frac{\|V(a)\|_{\delta}}{(\log|a|)^{\delta/2}} \leq c \quad |a| \rightarrow 0.$$

We have a following decomposition:

$$C_K \cong C_{K,1} \times N.$$

Here $C_{K,1}$ is the maximal compact subgroup: $\{g \in C_K \mid |g| = 1\}$ and $N = \mathbb{R}_{>0}^*$. Let χ_0 be a character of $C_{K,1}$. We use $\tilde{\chi}_0$ to denote an extension of χ_0 as a character of C_K . Namely, $\tilde{\chi}_0(g) = \chi_0(g); \forall g \in C_{K,1}$. Here $\tilde{\chi}_0$ has the form $\tilde{\chi}_0 = \chi_0 |\cdot|^{\rho}$, $\rho \in i\mathbb{R}$. Restrict V to $C_{K,1}$, one decompose $L^2_{\delta}(C_K)$ in the direct sum of the finite dimensional subspaces,

$$L^2_{\delta, \chi_0} = \left\{ \xi \in L^2_{\delta}(C_K) \mid \xi(a^{-1}g) = \chi_0(a)\xi(g) \quad \forall g \in C_K \quad \forall a \in C_{K,1} \right\}.$$

The dual space $(L^2_{\delta}(C_K))^*$ of $L^2_{\delta}(C_K)$ can be identified with $L^2_{-\delta}(C_K)$. It is also decomposed in the direct sum of the subspaces,

$$L^2_{-\delta, \chi_0} = \left\{ \xi \in L^2_{-\delta}(C_K) \mid \xi(ag) = \chi_0(a)\xi(g) \quad \forall g \in C_K \quad \forall a \in C_{K,1} \right\}.$$

Here, we use the transposed of V

$$(V^\tau(a)\eta)(x) = \eta(ax); \quad \eta(x) \in (L^2_\delta(C_K))^*.$$

The pairing between $L^2_\delta(C_K)$ and its dual $(L^2_\delta(C_K))^* = L^2_{-\delta}(C_K)$ is given by

$$\langle f, \eta \rangle = \int_{C_K} f(x)\eta(x)d^*x.$$

We can obtain the following exact sequences:

$$0 \rightarrow L^2_\delta(X)_0 \xrightarrow{\text{T}} L^2_\delta(C_K) \rightarrow \mathcal{H} \rightarrow 0.$$

Let

$$\text{Im}(\text{T})^0 = \left\{ \eta \in (L^2_\delta(C_K))^* \mid \langle \text{T}f, \eta \rangle = 0, \quad \forall f \in \mathcal{S}(A_K)_0 \right\}.$$

It holds that

$$\eta(x) \in \text{Im}(\text{T})^0 \iff \int_{C_K} \text{T}f(a)\eta(a)d^*a = 0, \quad \forall f \in \mathcal{S}(A_K)_0.$$

For $\eta(x) \in L^2_{-\delta, \chi_0}$, we may think that it has the form:

$$\eta(x) = \tilde{\chi}_0(x)\Psi(|x|).$$

Now

$$\Psi(|x|) = \int_{-\infty}^{\infty} \hat{\Psi}(t)|x|^t dt$$

where $\hat{\Psi}(t) = \int_{C_K} \Psi(a)|a|^t d^*a$. Thus,

$$\eta(x) = \int_{-\infty}^{\infty} \eta(x; t)dt; \quad \eta(x; t) = \tilde{\chi}_0(x)|x|^t \hat{\Psi}(t).$$

Then,

$$\begin{aligned} \eta(x) \in \text{Im}(\text{T})^0 &\iff \langle \text{T}f, \eta \rangle = 0 \\ &\iff \int_{C_K} \text{T}f(a) \int_{-\infty}^{\infty} \tilde{\chi}_0(a)|a|^t \hat{\Psi}(t) dt d^*a = 0 \\ &\iff \int_{-\infty}^{\infty} \int_{C_K} \text{T}f(a) \tilde{\chi}_0(a)|a|^t \hat{\Psi}(t) d^*a dt = 0, \quad \forall f \in \mathcal{S}(A_K)_0. \end{aligned}$$

As the consequence of Tate's work,

Lemma 2.1. For $\text{Re}(s) > 0$, and any character χ_0 of C_K ,

$$\int_{C_K} \mathbb{T}f(a)\chi_0(a)|a|^{s-1/2} d^*a = L(\chi_0, s)D'_s(f), \quad \forall f \in \mathcal{S}(A_K)_0.$$

Here, $D'_s(f)$ is a holomorphic function of s ($\text{Re}(s) > 0$).

From this lemma, we can say that

$$\eta(x) \in \text{Im}(\mathbb{T})^0 \iff L(\tilde{\chi}_0, 1/2 + \rho) = 0; \quad \rho \in i\mathbb{R}.$$

Here $\mathcal{H} \cong L^2_{\mathcal{S}}(C_K)/\text{Im}(\mathbb{T})$. Think of the left regular representation W of C_K on \mathcal{H} : (W, \mathcal{H}) , where one deduces W from V . Restrict W to $C_{K,1}$, one decompose \mathcal{H} in the direct sum of the subspaces,

$$\mathcal{H} = \bigoplus_{\chi_0 \in \hat{C}_{K,1}} d(\chi_0)\mathcal{H}_{\chi_0} \quad d(\chi_0) < \infty$$

where $\mathcal{H}_{\chi_0} = \{ \xi \mid W(a)\xi = \chi_0(a)\xi, \quad \forall a \in C_{K,1} \}$ and we denote the dimension of \mathcal{H}_{χ_0} by $d(\chi_0)$. We will also consider its dual. We obtain the transposition W^τ of C_K on \mathcal{H}^* : (W^τ, \mathcal{H}^*) , where one deduces W^τ from V^τ . Now, let h be a test function on C_K and set

$$W(h) = \int_{C_K} h(g)W(g) d^*g.$$

Denote h 's Fourier transform by \hat{h} :

$$\hat{h}(\chi, z) = \int_{C_K} h(\mu)\chi(\mu)|\mu|^z d^*\mu.$$

Recall

$$\mathcal{H}^* \cong (L^2_{\mathcal{S}}(C_K)/\text{Im}(\mathbb{T}))^* \cong \text{Im}(\mathbb{T})^0,$$

moreover

$$\text{tr}W = \text{tr}W^\tau.$$

We can compute

$$\begin{aligned} \int_{C_K} h(g)(V^\tau(g)\eta)(x) d^*g &= \int_{C_K} h(g)(V^\tau(g) \int_{-\infty}^{\infty} \eta(\cdot; t) dt)(x) d^*g \\ &= \int_{C_K} \int_{-\infty}^{\infty} h(g)\tilde{\chi}_0(g)|g|^t \tilde{\chi}_0(x)|x|^t \hat{\Psi}(t) dt d^*g \end{aligned}$$

$$= \int_{-\infty}^{\infty} \hat{h}(\tilde{\chi}_0, it) \tilde{\chi}_0(x) |x|^{it} \hat{\Psi}(t) dt,$$

thus

$$\begin{aligned} \langle Tf, (V^\tau(h)\eta)(x) \rangle &= \int_{C_K} Tf(a) \int_{-\infty}^{\infty} \hat{h}(\tilde{\chi}_0, it) \tilde{\chi}_0(a) |a|^{it} \hat{\Psi}(t) dt d^*a \\ &= \int_{-\infty}^{\infty} \int_{C_K} Tf(a) \tilde{\chi}_0(a) |a|^{it} \hat{h}(\tilde{\chi}_0, it) \hat{\Psi}(t) d^*a dt. \end{aligned}$$

If $\eta(x) \in \text{Im}(T)^0$ then $\langle Tf, (V^\tau(h)\eta)(x) \rangle = 0, \forall f \in \mathcal{S}(A_K)_0$. Therefore, $(V^\tau(h)\eta)(x) \in \text{Im}(T)^0$. The above computation shows that

$$\hat{h}(\tilde{\chi}_0, \rho) L(\tilde{\chi}_0, 1/2 + \rho) = 0; \rho \in i\mathbb{R}.$$

So, we see that

$$\text{tr}W(h) = \sum_{\substack{L(\tilde{\chi}_0, 1/2+\rho)=0 \\ \rho \in i\mathbb{R}}} \hat{h}(\tilde{\chi}_0, \rho).$$

Let $\chi_0 \in \hat{C}_{K,1}$. Recall $\tilde{\chi}_0 = \chi_0 |\cdot|^\rho (\rho \in \mathbb{C})$. The action of C_K on \mathcal{H}_{χ_0} can be $W(g)\xi = \tilde{\chi}_0(g)\xi$, and it turns out that $W(g)\xi = |g|^\rho \xi; g \in N$. So it is satisfied that

$$|g|^{\text{Re}(\rho)} \leq \|W(g)\|_\delta, \quad g \in C_K.$$

Let $W_{\chi_0} = W|_{\mathcal{H}_{\chi_0}}$ and $e^t = g (g \in N)$. We will rewrite the action of N on \mathcal{H}_{χ_0} as

$$W_{\chi_0}(e^t): \mathbb{R} \rightarrow \mathcal{H}_{\chi_0}.$$

The following things

- (a) $W_{\chi_0}(e^0) = 1,$
- (b) $W_{\chi_0}(e^{t+s}) = W_{\chi_0}(e^t)W_{\chi_0}(e^s)$

are satisfied. Thus $W_{\chi_0}(e^t)$ is a semi-group. From the theory of semi-group, we can say that

$$W_{\chi_0}(e^t) = e^{tD_{\chi_0}}$$

where

$$D_{\chi_0}\xi = \lim_{t \rightarrow 0^+} \frac{W_{\chi_0}(e^t)\xi - W_{\chi_0}(e^0)\xi}{t} \quad \xi \in \mathcal{H}_{\chi_0}.$$

The operator D_{χ_0} has discrete spectra. We may think that the discrete spectrum is given by the element ξ which belongs to $\text{Im}(T)^0$.

Let $\tilde{\chi}_0$ be the unique extension of $\chi_0 \in \hat{C}_{K,1}$ to C_K which is equal to 1 on N . We see that $\chi = \tilde{\chi}_0 \cdot |\cdot|^{it_0}$ ($t_0 \in \mathbb{R}$) for $\chi \in \hat{C}_K$. Then $L(\chi, 1/2+it) = L(\tilde{\chi}_0, 1/2+i(t_0+t))$. Thus, as the extension of χ_0 , we will use the above unique extension $\tilde{\chi}_0$.

Theorem 2.1. $\chi_0 \in \hat{C}_{K,1}$, $\delta > 1$. Then D_{χ_0} has discrete spectra, $\text{sp}D_{\chi_0} \subset i\mathbb{R}$ is the set of imaginary parts of zeros of the L function with Grossencharacter $\tilde{\chi}_0$ which have real part equal to $1/2$;

$$\rho \in \text{sp}D \iff L(\tilde{\chi}_0, 1/2+\rho) = 0 \text{ and } \rho \in i\mathbb{R}, \text{ where } \tilde{\chi}_0 \text{ is the unique extension of } \chi_0 \text{ to } C_K \text{ which is equal to 1 on } N.$$

Moreover the multiplicity of ρ in $\text{sp}D$ is equal to the largest integer of $n < \frac{1+\delta}{2}$, $n \leq$ multiplicity of $1/2+\rho$ as a zero of L .

The action of N is that

$$W_{\chi_0}(e^t)\xi = |e^t|^\rho \xi = e^{\rho t} \xi.$$

Then,

$$D_{\chi_0} \xi = \left. \frac{dW_{\chi_0}(e^t)\xi}{dt} \right|_{t=0} = \rho \xi.$$

Therefore, ρ is the spectrum of D_{χ_0} . Consider

$$\lim_{|g| \rightarrow \infty} \frac{|g|^\alpha}{\log|g|} = \infty \quad (\alpha > 0) \quad \text{and} \quad \lim_{|g| \rightarrow 0} \frac{|g|^\alpha}{\log|g|} = \infty \quad (\alpha < 0).$$

Because $|g|^{\text{Re}(\rho)} \leq \|W(g)\|_\delta$; $g \in C_K$, if $\text{Re}(\rho) > 0$ or $\text{Re}(\rho) < 0$ then each of them conflicts with

$$\|V(a)\|_\delta = O((\log|a|)^{\delta/2}) \quad |a| \rightarrow \infty$$

or

$$\|V(a)\|_\delta = O((\log|a|)^{\delta/2}) \quad |a| \rightarrow 0$$

Therefore, it is that $\rho = it$ ($t \in \mathbb{R}$). Thus,

$$\tilde{\chi}_0 = \chi_0 \cdot |\cdot|^{it} \quad t \in \mathbb{R}.$$

We see that D_{χ_0} has a purely imaginary spectrum, so we obtain the following corollary.

Corollary 2.2. For any Shwarz function $h \in \mathcal{S}(C_K)$ the operator $\int_{C_K} h(g)W(g) d^*g$ in \mathcal{H} is of trace class, and its trace is given by

$$\text{tr}W(h) = \sum_{\substack{L(\tilde{\chi}_0, 1/2+\rho)=0 \\ \rho \in i\mathbb{R}}} \hat{h}(\tilde{\chi}_0, \rho)$$

where the multiplicity is counted as in Theorem 2.1. and where Fourier transform \hat{h} of h is defined by $\hat{h}(\chi, z) = \int_{C_K} h(\mu)\chi(\mu)|\mu|^z d^*\mu$.

We can obtain the following exact sequences:

$$0 \rightarrow L^2_{\delta}(X)_0 \rightarrow L^2_{\delta}(X) \rightarrow \mathbb{C} \oplus \mathbb{C}(1) \rightarrow 0$$

and

$$0 \rightarrow L^2_{\delta}(X)_0 \xrightarrow{\text{T}} L^2_{\delta}(C_K) \rightarrow \mathcal{H} \rightarrow 0 .$$

We will compute $\text{tr}U(h)$ for $(U, L^2_{\delta}(X))$ from spectral side. From the above first sequence, considering *Lefchetz formula*, we will see that

$$A = \text{tr}U|_{L^2_{\delta}(X)_0} - \text{tr}U|_{L^2_{\delta}(X)} + \text{tr}U|_{\mathbb{C} \oplus \mathbb{C}(1)} .$$

From the second sequence, we will obtain

$$A' = \text{tr}U|_{L^2_{\delta}(X)_0} - \text{tr}U|_{L^2_{\delta}(C_K)} + \text{tr}U|_{\mathcal{H}} .$$

Therefore, it is satisfied that

$$\text{tr}U|_{L^2_{\delta}(X)} = \text{tr}U|_{\mathbb{C} \oplus \mathbb{C}(1)} - \text{tr}U|_{\mathcal{H}} + \text{tr}U|_{L^2_{\delta}(C_K)} + A' - A .$$

We try to compute $\text{tr}U(h)$ spectrally. Here,

$$U(h) = \int_{C_K} h(g)U(g) d^*g .$$

The first term $\text{tr}U|_{\mathbb{C} \oplus \mathbb{C}(1)}$ gives

$$\hat{h}(0) + \hat{h}(1) .$$

Considering that $T(U(g)\xi)(a) = |g|^{1/2}(V(g)T\xi)(a)$,

$$U|_{L^2_\delta(C_K)} \text{ is } (|\cdot|^{1/2}V, L^2_\delta(C_K)) \text{ and } U|_{\mathcal{H}} \text{ is } (|\cdot|^{1/2}V, \text{Im}(T)^0).$$

So we will understand that the second term gives

$$\sum_{\substack{L(\tilde{\chi}_0, \rho) = 0 \\ \text{Re } \rho = 1/2}} \hat{h}(\tilde{\chi}_0, \rho).$$

Finally, the term $\text{tr}U|_{L^2_\delta(C_K)} + A' - A$ gives $\infty h(1)$. Thus,

$$\text{tr}U(h) = \hat{h}(0) + \hat{h}(1) - \sum_{\substack{L(\tilde{\chi}_0, \rho) = 0 \\ \text{Re } \rho = 1/2}} \hat{h}(\tilde{\chi}_0, \rho) + \infty h(1).$$

3.

We try to compute $\text{tr}U$ geometrically.

Let S be a finite set of places of K containing all infinite places. Set

$$A_S = \prod_{v \in S} K_v \times \prod_{v \notin S} R_v \quad \text{and} \quad J_S = \prod_{v \in S} K_v^* \times \prod_{v \notin S} R_v^*$$

where R_v is the ring of integers of K_v . The S -units of K is given by

$$\mathcal{O}_S^* = J_S \cap K^*.$$

The idele class C_K is embedded in $C_S = J_S / \mathcal{O}_S^*$ and $X_S = A_S / \mathcal{O}_S^*$ plays the same role as X . We will think of $L^2(X_S)$ which is obtained by a completion of $\mathcal{S}(A_S)$. Let

$$R_\Lambda = \hat{P}_\Lambda P_\Lambda, \quad \Lambda \in \mathbb{R}_+.$$

Here P_Λ is the orthogonal projection onto the subspace,

$$P_\Lambda = \left\{ \xi \in L^2(X_S) \mid \xi(x) = 0, \quad \forall x, |x| > \Lambda \right\}$$

while $\hat{P}_\Lambda = F P_\Lambda F^{-1}$ where F is the Fourier transform.

Theorem 3.1. For any $h \in \mathcal{S}_c(C_S)$, one has

$$\text{Trace}(R_\Lambda U(h)) = 2 \log'(\Lambda) h(1) + \sum_{v \in S} \int_{K_v^*}' \frac{h(\mu^{-1})}{|1 - \mu|} d^* \mu + o(1) \quad \Lambda \rightarrow \infty$$

where $2 \log'(\Lambda) = \int_{\lambda \in C_S, |\lambda| \in [\Lambda^{-1}, \Lambda]} d^* \lambda$.

Let χ_0 be a character of $C_{S,1}$ which is the subgroup: $\{g \in C_S \mid |g| = 1\}$. The Hilbert space $L^2(X_S)$ is decomposed in the subspace,

$$L^2_{\chi_0} = \left\{ \xi \in L^2(X_S) \mid \xi(a^{-1}x) = \chi_0(a) \xi(x) \quad \forall x \in X_S, a \in C_{S,1} \right\}.$$

Let \mathcal{U}_S be the image in C_S of the open subgroup $\prod R_v^*$. Fix a character χ of \mathcal{U}_S , and think of χ_0 whose restriction to \mathcal{U}_S is equal to χ . Set

$$L^2(X_S)_\chi = \left\{ \xi \in L^2(X_S) \mid \xi(a^{-1}x) = \chi(a)\xi(x) \quad \forall x \in X_S, a \in \mathcal{U}_S \right\}.$$

We can find $h_\chi \in \mathcal{S}(C_S)$ such that

$$\text{Supp}(h_\chi) = \mathcal{U}_S \quad h_\chi(x) = \lambda \bar{\chi}(x) \quad \forall x \in \mathcal{U}_S$$

where the constant λ is determined by corresponding normalization of the Haar measure on C_S .

Let $B_\Lambda = \text{Im}(P_\Lambda) \cap \text{Im}(\hat{P}_\Lambda)$ be the intersection of the ranges of the projection P_Λ and \hat{P}_Λ . We will think of B_Λ^χ which is the intersection of B_Λ with $L^2(X_S)_\chi$. For each character χ of \mathcal{U}_S , we can find a vector $\eta_\chi \in L^2(X_S)_\chi$ such that

$$U(g)(\eta_\chi) \in B_\Lambda \quad g \in C_S, \Lambda^{-1} \leq |g| \leq \Lambda.$$

Then B_Λ^χ is given as the linear span of $U(g)(\eta_\chi)$:

$$B_\Lambda^\chi = \sum_{g \in D_S \mid |g| \in [\Lambda^{-1}, \Lambda]} \lambda_g U(g)(\eta_\chi) \quad D_S = C_S / \mathcal{U}_S.$$

Set

$$(B_\Lambda^\chi)^0 = \text{The whole of } \sum_{g \in D_S \mid |g| \in [\Lambda^{-1}, \Lambda]: \text{finite sum}} \lambda_g U(g)(\eta_\chi).$$

It turns out that $(B_\Lambda^\chi)^0 \subseteq B_\Lambda^\chi \subseteq L^2(X_S)_\chi$. We may say that $(B_\Lambda^\chi)^0$ is dense in B_Λ^χ . So, from the compactness of $\{g \in C_S \mid \Lambda^{-1} \leq |g| \leq \Lambda\}$, we can consider that B_Λ^χ is a vector space which has a countable basis at most. It must be hard to show that $B_\Lambda^\chi = L^2(X_S)_\chi$ for sufficient large Λ . We will replace R_Λ by the orthogonal projection Q_Λ on $\text{Im}(P_\Lambda) \cap \text{Im}(\hat{P}_\Lambda)$. Suppose that $B_\Lambda^\chi = L^2(X_S)_\chi$ for sufficient large Λ . Then we can identify $\text{tr}R_\Lambda U$ with $\text{tr}Q_\Lambda U$ of $(U, L^2(X_S))$. From the Theorem 3.1., we can show the following.

Corollary. Let Q_Λ be the orthogonal projection on the subspace of $L^2(X_S)$ spanned by the $f \in \mathcal{S}(A_S)$, which vanish as well as Fourier transform for $|x| > \Lambda$. Let $h \in \mathcal{S}(C_S)$ have compact support. Then when $\Lambda \rightarrow \infty$, one has

$$\text{Trace}(Q_\Lambda U(h)) = 2h(1)\log'(\Lambda) + \sum_{v \in S} \int_{K_v^*} \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu + o(1)$$

where $2\log'(\Lambda) = \int_{\lambda \in C_S, |\lambda| \in [\Lambda^{-1}, \Lambda]} d^* \lambda$.

We can get from the above corollary an S-independent global formulation:

$$\text{Trace } Q_\Lambda U(h) = 2h(1) \log' \Lambda + \sum_v \int_{K_v^*} \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu + o(1) \quad \Lambda \rightarrow \infty$$

where $Q_\Lambda U$ is a trace class operator for $(U, L^2(X))$.

In order to obtain the identification of $\text{tr} Q_\Lambda U$ with $\text{tr} R_\Lambda U$, we have to show that $B_\Lambda^\chi = L^2(X_S)_\chi$ for sufficient large Λ . If C_S is compact then the compactness must be sufficient for us to show the equation. If C_K were compact, we could show the validity of the Riemann hypothesis.

So it must be interesting to think of the compactification of C_K . With this interest, we will examine the space $Y = A_K/K$. As the same way in the case of X , we can obtain $L^2(Y)$ and $L^2(Y)_0$. We will think of the case $K = \mathbb{Q}$. It holds that

$$A_{\mathbb{Q}} = \prod_{p < \infty} \mathbb{Z}_p \times [0, 1) + \mathbb{Q} \quad \text{and} \quad A_{\mathbb{Q}}^* = \left(\prod_{p < \infty} \mathbb{Z}_p^* \times \mathbb{R}_{>0}^* \right) \cdot \mathbb{Q}^*.$$

Thus, it turns out that

$$Y = A_{\mathbb{Q}}/\mathbb{Q} \cong \prod_{p < \infty} \mathbb{Z}_p \times [0, 1] \quad \text{and} \quad C_{\mathbb{Q}} = A_{\mathbb{Q}}^*/\mathbb{Q}^* \cong \prod_{p < \infty} \mathbb{Z}_p^* \times \mathbb{R}_{>0}^*.$$

Think of $r \mapsto 2/\pi \tan(r)^{-1}$; $r \in \mathbb{R}_{>0}^*$, it must be allowed to say that

$$\mathbb{R}_{>0}^* \text{ is embedded in } [0, 1].$$

Thus,

$$C_{\mathbb{Q}} \text{ is embedded in } \prod_{p < \infty} \mathbb{Z}_p^* \times [0, 1].$$

It may be allowed to say that

$$Y = \{x \in Y \mid |x| < 1\} \cup \{x \in Y \mid |x| = 1\}$$

and that $\{x \in Y \mid |x| = 1\}$ consists of the boundary of Y . Denote it by ∂Y . It must correspond to $\prod_{p < \infty} \mathbb{Z}_p^* \times \{1\}$. Let

$$C_{\mathbb{Q}} = \prod_{p < \infty} \mathbb{Z}_p^* \times (0, 1].$$

We will think that $C_{\mathbb{Q}}$ is the compactification of $C_{\mathbb{Q}}$. We expect that $C_{\mathbb{Q}}$ fills the same role of $C_{\mathbb{Q}}$.

We can obtain an exact sequence:

$$0 \rightarrow L^2(Y)_0 \xrightarrow{T} L^2(\mathcal{C}_Q) \rightarrow \mathcal{H} \rightarrow 0$$

where $\mathcal{H} \cong L^2(\mathcal{C}_Q)/\text{Im}(T)$. Let U be a left regular representation of \mathcal{C}_Q on $L^2(Y, dx)$ and V be a left regular representation of \mathcal{C}_Q on $L^2(\mathcal{C}_Q, d^*x)$. One deduces the left regular representation W of \mathcal{C}_Q on \mathcal{H} from V . One may be allowed to say that \mathcal{C}_Q is compact because it must be complete and totally bounded. So one can decompose $L^2(\mathcal{C}_Q)$ in the direct sum of 1-dimensional subspaces,

$$L^2_{\chi_0} = \left\{ \xi \in L^2(\mathcal{C}_Q) \mid \xi(a^{-1}g) = \chi_0(a)\xi(g) \quad \forall g, a \in \mathcal{C}_Q \right\}.$$

The dual space $(L^2(\mathcal{C}_Q))^*$ of $L^2(\mathcal{C}_Q)$ can be identified with $L^2(\mathcal{C}_Q)$.

[Remark] The left regular representation U of \mathcal{C}_Q on $L^2(Y, dx)$ isn't unitary. But the left regular representation T of Y on $L^2(Y, dx)$:

$$(T(g)\xi)(x) = \xi(-g+x) \quad g, x \in Y$$

is unitary. Because Y is abelian and compact, we obtain the following decomposition:

$$L^2(Y) = \bigoplus_{\chi \in \hat{Y}} L^2_{\chi}(Y) \quad T_{\chi} = T|_{L^2_{\chi}(Y)}$$

where T_{χ} is 1-dimensional representation.

Here Y is compact. Thus the following formula:

$$\text{tr}U|_{L^2(Y)_0} = \text{tr}U|_{L^2(\mathcal{C}_Q)} - \text{tr}U|_{\mathcal{H}} + A$$

becomes meaningful.

Now, our problem is to compute $\text{tr}U|_{L^2(Y)_0}$. Basically, we may think that this problem is how to construct $L^2(Y)_0$. Set

$$\Delta = |x|^2 \frac{d^2}{dx^2}$$

which is a differential operator on Y . We shall think of the eigenvalue problems:

$$\Delta\xi - \lambda\xi = 0, \quad \xi(x) = 0 \text{ on } \partial Y$$

on the analogy of Sturm-Liouville problem. Recall that the action of \mathcal{C}_Q on the functions on Y is

$$(U(g)\phi)(x) = \phi(g^{-1}x) \quad \forall g \in C_0, x \in Y.$$

It turns out that $U(g)$ and Δ are commutative. Hence they share the same eigenspace. We try to construct the $L^2(Y)_0$ space as the space of eigenfunctions of Δ .

[Remark] One computes

$$(U(g)|x|^2 \frac{d^2}{dx^2} \phi)(x) = (U(g)|\cdot|^2 \phi'')(x) = |g^{-1}x|^2 \phi''(g^{-1}x).$$

It holds that $dgx = |g|dx$, so

$$\begin{aligned} |x|^2 \frac{d^2}{dx^2} (U(g)\phi)(x) &= |x|^2 \frac{d^2}{dx^2} \phi(g^{-1}x) = |x|^2 \frac{d^2 g^{-1}x}{dx^2} \frac{d^2}{d(g^{-1}x)^2} \phi(g^{-1}x) \\ &= |g^{-1}x|^2 \phi''(g^{-1}x). \end{aligned}$$

[Remark] The Δ becomes a differential operator on X . Since $dg^{-1}x = |g^{-1}|dx$, if one restricts U to $C_{K,1}$ then $|d(g^{-1}x)| = |dx|$. Namely, $|dx|$ is invariant under the action $U(g); \forall g \in C_{K,1}$. Thus, the Laplacian $\frac{d^2}{dx^2}$ is $C_{K,1}$ -invariant.

If we can show that the Laplacian $\frac{d^2}{dx^2}$ is C_K -invariant then we can say that $U(g); \forall g \in C_K$ is isometry, namely unitary. On the other hand it does not always mean that $U(g); \forall g \in C_K$ is unitary.

Considering $\xi'' = \lambda \frac{\xi}{|x|^2}$, it turns out that

$$(\xi' \bar{\xi})' = \xi'' \bar{\xi} + \xi' \bar{\xi}' = \lambda \frac{\xi \bar{\xi}}{|x|^2} + \xi' \bar{\xi}'.$$

One computes

$$\int_Y \xi' \bar{\xi}' dx + \lambda \int_Y \frac{\xi \bar{\xi}}{|x|^2} dx = \int_Y (\xi' \bar{\xi})' dx.$$

We can write that $\int_Y (\xi' \bar{\xi})' dx = \int_{\partial Y} \xi' \bar{\xi} dx$. From the boundary condition, it holds that

$$\int_{\partial Y} \xi' \bar{\xi} dx = 0.$$

Therefore, we obtain

$$\int_Y \xi' \bar{\xi}' dx = -\lambda \int_Y \frac{\xi \bar{\xi}}{|x|^2} dx.$$

Here $\int_Y \xi' \bar{\xi}' dx, \int_Y \frac{\xi \bar{\xi}}{|x|^2} dx \geq 0$. Thus, $\lambda \leq 0$. Write a function $\xi(x)$ on Y

$$\xi(x) = \xi(ut) \quad u \in \prod_{p < \infty} \mathbb{Z}_p, t \in [0, 1].$$

We will compute as follows.

$$(a) \quad \frac{\partial}{\partial u} x = t.$$

$$(b) \quad \frac{\partial}{\partial u} \xi(x) = \frac{\partial x}{\partial u} \frac{\partial}{\partial x} \xi(x) = t \xi'(x).$$

[Remark] From the definition, the following things are satisfied.

$$(a) \quad dx = du dt, \text{ so } \frac{dx}{du} = dt.$$

$$(b) \quad \frac{d}{du} \xi(x) = \frac{dx}{du} \frac{d}{dx} \xi(x) = \xi'(x) dt.$$

It holds that $\frac{\partial^2}{\partial u^2} = t^2 \frac{d^2}{dx^2}$, so we see that $|u|^2 \frac{\partial^2}{\partial u^2} = |u|^2 t^2 \frac{d^2}{dx^2} = |x|^2 \frac{d^2}{dx^2}$. Thus, we can identify Δ with $|u|^2 \frac{\partial^2}{\partial u^2}$. We will think of the eigenvalue problems:

$$|u|^2 \frac{\partial^2}{\partial u^2} \xi(x) - \lambda \xi(x) = 0.$$

Let

$$\eta(u) = \begin{cases} \xi(u0) & \dots \quad u \in \prod_{p < \infty} \mathbb{Z}_p^* \\ \xi(u1) & \dots \quad \text{otherwise} \end{cases} .$$

Then we can interpret the eigenvalue problems as the following problem;

$$|u|^2 \frac{\partial^2}{\partial u^2} \eta(u) - \lambda \eta(u) = 0, \quad \eta(u) = 0 \text{ on } \prod_{p < \infty} \mathbb{Z}_p^* .$$

Here we will identify $\prod_{p < \infty} \mathbb{Z}_p^*$ with ∂Y .

[Remark] Here,

$$|u|^2 \frac{\partial^2}{\partial u^2} \xi(x) = |u|^2 |t|^2 \xi''(ut).$$

From the above definition of $\eta(u)$, let

$$\eta''(u) = |0|^2 \xi''(u0) \quad \forall u \in \prod_{p < \infty} \mathbb{Z}_p^* \quad \text{and} \quad \eta''(u) = |1|^2 \xi''(u1) \quad \forall u \in \prod_{p < \infty} \mathbb{Z}_p^* .$$

Then we can say that $|u|^2 \frac{\partial^2}{\partial u^2} \xi(x)$ gives $|u|^2 \frac{\partial^2}{\partial u^2} \eta(u)$.

We can show that $\lambda \leq 0$. Here, think of *the heat equation*:

$$\left\{ \begin{array}{l} \frac{\partial}{\partial t} \xi(ut) = |u|^2 \frac{\partial^2}{\partial u^2} \xi(ut) \\ \eta(u) = \xi(u0) = 0 \quad \text{on} \quad \prod_{p < \infty} \mathbb{Z}_p^* . \end{array} \right.$$

Let $\eta_\lambda(u)$ be an eigenfunction of $|u|^2 \frac{\partial^2}{\partial u^2}$ with eigenvalue λ . Then, $e^{-\lambda t} \eta_\lambda(u)$ is a particular solution. We obtain general solutions

$$\xi(x) = \sum_{\lambda} c_{\lambda} \cdot e^{-\lambda t} \eta_{\lambda}(u).$$

Here, c_{λ} is the constant. There exists some function, called *heat kernel*, $p(t, \mu, \nu)$ on Y^2 and we can say that

$$\xi(ut) = \int_Y p(t, u, \nu) \eta(\nu) d\nu .$$

From the theory of semi-group, it holds that

$$\sum_{\lambda} e^{-\lambda t} = \int_Y p(t, u, u) du .$$

We can say that such a function $\eta(u)$ associated with $|x|^s$

$$|\cdot|^s(u) = \left\{ \begin{array}{l} |u0|^s \quad \dots \quad u \in \prod_{p < \infty} \mathbb{Z}_p^* \\ |u1|^s \quad \dots \quad \text{otherwise} \end{array} \right.$$

is an eigenfunction of Δ with eigenvalue $\lambda = s(s-1) \leq 0$. It must be allowed that

$$L^2(Y) \text{ is decomposed in the subspace } \{c | \cdot |^s \mid c \in \mathbb{C}\}.$$

Here $|0|^s = 0$, so we can say that $|\cdot|^s \in L^2(Y)_0$. Moreover, since $L^2(Y)$ is a Hilbert space, $\{|\cdot|^s\}$ is discrete. Now

$$(U(g)|\cdot|^s)(x) = |g^{-1}x|^s = |g^{-1}|^s |x|^s \quad \forall g \in \mathcal{C}_Q, x \in Y.$$

We shall think that $|g^{-1}|^s = |g|^{-s}$ is extended as a quasi-character of \mathcal{C}_Q . We may be allowed to think that the quasi-character $|g|^{-s}$ is equivalent to a quasi-character $|g|^s$. We may say that

$$\text{tr}U|_{L^2(Y)_0} \text{ extends over } \{|a|^s \mid s(s-1) = \lambda\}.$$

Moreover,

$$\text{tr}U|_{L^2(\mathcal{C}_Q)} \text{ extends over } \{\chi(a)|a|^{1/2} \mid \chi \text{ is a character of } \mathcal{C}_Q\}$$

and

$$\text{tr}U|_{\mathcal{H}} \text{ extends over}$$

$$\{\pi(a)|a|^{1/2} \mid \pi \text{ is the character of } \mathcal{C}_Q \text{ which is given by } \eta(x) \in \text{Im}(T)^0\}.$$

Recall $T(U(g)\xi)(a) = |g|^{1/2}(V(g)T\xi)(a)$,

$$\begin{aligned} T(U(g)|\cdot|^s)(a) &= |g|^{1/2}(V(g)T|\cdot|^s)(a) \\ &= |g|^{1/2}(V(g)|\cdot|^{1/2+s})(a) \\ &= |g|^{-s}|a|^{1/2+s} \end{aligned}$$

$|g|^{-s}$ being equivalent to $|g|^s$

$$= |g|^s|a|^{1/2+s}.$$

Thus we see that $\text{tr}U|_{L^2(\mathcal{C}_Q)}$ contains $\text{tr}U|_{L^2(Y)_0}$ and it also expands over such a quasi-character as $|a|^s$.

The compactness of Y guarantees to compute

$$\text{tr}U(h) = \sum_p \int_{\mathbb{Q}_p^*} \frac{h(\mu^{-1})}{|1-\mu|} d^* \mu.$$

Thus we can say that

$$\sum_{L(\chi, \rho) = 0} \hat{h}(\chi, \rho) = \sum_{\substack{L(\chi, \rho) = 0 \\ \text{Re } \rho = 1/2}} \hat{h}(\chi, \rho).$$

[Remark] In order to obtain an expected formula like Theorem 3.1., we need the evaluation of a certain *error term*. We shall compare $(U, L^2(Y))$ with $(U, L^2_\delta(X))$. For the latter, U is also trace-class so that we formally get $\text{tr}U(h) = \sum_v \int_{K_v^*} \frac{h(\mu^{-1})}{|1-\mu|} d^*\mu$. However, since $L^2_\delta(X)$ is a weighted space L^2 , we can't always obtain the expected formula. On the other hand, in the case of $L^2(Y)$, we can expect to obtain the desired formula.

On the other hand, the compactness must also guarantee $\lambda \leq -1/4$. So, we will see a certain relationship between the validity of the Riemann hypothesis and the fact that $\lambda \leq -1/4$. We shall suppose that the compactification of C_Q is equivalent to the fact that $\lambda \leq -1/4$ for the eigenvalue λ of Δ on X . Then, we may say that the validity of the Riemann hypothesis is equivalent to showing $\lambda \leq -1/4$.

4.

We will think of the case $GL_2(\mathbb{Q}) \backslash GL_2(\mathbb{A}_{\mathbb{Q}})$.

Let (π, V) be an irreducible admissible infinite dimensional representation of $GL_2(\mathbb{A}_{\mathbb{Q}})$ with central character ω . Here, ω is a quasi-character of $GL_2(\mathbb{A}_{\mathbb{Q}})$ defined by

$$\pi\left(\begin{pmatrix} a & 0 \\ 0 & a \end{pmatrix}\right) = \omega(a) \text{id}V \quad a \in \mathbb{A}_{\mathbb{Q}}^*.$$

Suppose that $W(\pi, \psi)$ is the ψ -Whittaker model. Let χ be a character of $\mathbb{A}_{\mathbb{Q}}^*/\mathbb{Q}^*$. The *Jacquet–Langlands zeta integrals* are defined by

$$Z(s, W, \chi; g) = \int_{\mathbb{A}_{\mathbb{Q}}^*} W\left(\begin{pmatrix} a & 0 \\ 0 & 1 \end{pmatrix}g\right) \chi(a) |a|^{s-1/2} d^*a$$

and

$$Z^{\vee}(s, W, \chi; g) = \int_{\mathbb{A}_{\mathbb{Q}}^*} W\left(\begin{pmatrix} a & 0 \\ 0 & 1 \end{pmatrix}g\right) \chi(a) |a|^{s-1/2} \omega^{-1}(a) d^*a.$$

There exists $s_0 \in \mathbb{R}$ such that $Z(s, W, \chi; g)$ and $Z^{\vee}(s, W, \chi; g)$ absolutely converge whenever $\text{Re}(s) > s_0$ for all $g \in GL_2(\mathbb{A}_{\mathbb{Q}})$ and $W \in W(\pi, \psi)$. There exists a unique L-function $L(s, \pi \otimes \chi)$ ($(\pi \otimes \chi)(g) = \pi(g)\chi(\det g)$) such that

$$\phi(s, W, \chi; g) = Z(s, W, \chi; g) / L(s, \pi \otimes \chi)$$

is entire in s for all $g \in GL_2(\mathbb{A}_{\mathbb{Q}})$ and $W \in W(\pi, \psi)$. Therefore, we may say

$$Z(s, W, \chi; g) = 0 \iff L(s, \pi \otimes \chi) = 0.$$

Moreover, we will see that

$$Z(s, W, \chi; \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}) = 0 \iff L(s, \pi \otimes \chi) = 0.$$

It must be instructive to compare the *Jacquet–Langlands zeta integral* with the *Tate integral*. The *Tate integral* is defined by

$$Z(s, \chi, \Phi) = \int_{\mathbb{A}_{\mathbb{Q}}^*} \Phi(a) \chi(a) |a|^s d^*a$$

where χ is a character of $\mathbb{A}_{\mathbb{Q}}^*$ and $\Phi \in \mathcal{S}(\mathbb{A}_{\mathbb{Q}})$. We will see that $\Phi(x)$ corresponds to $|x|^{-1/2} W\left(\begin{pmatrix} x & 0 \\ 0 & 1 \end{pmatrix}g\right)$.

[Remark] We may say that $W\left(\begin{smallmatrix} a & 0 \\ 0 & 1 \end{smallmatrix}\right)g \in L^2(\mathbb{A}_{\mathbb{Q}}^*/\mathbb{Q}^*, d^*x)$. Thus we see that

$$|x|^{-1/2}W\left(\begin{smallmatrix} x & 0 \\ 0 & 1 \end{smallmatrix}\right)g \in L^2(\mathbb{A}_{\mathbb{Q}}/\mathbb{Q}^*, dx).$$

The *Jacquet-Langlands zeta integral* is defined by

$$Z(s, \chi, |x|^{-1/2}W\left(\begin{smallmatrix} x & 0 \\ 0 & 1 \end{smallmatrix}\right)g) = \int_{\mathbb{A}_{\mathbb{Q}}^*} |a|^{-1/2} W\left(\begin{smallmatrix} a & 0 \\ 0 & 1 \end{smallmatrix}\right)g \chi(a) |a|^s d^*a.$$

Considering $\phi(s, W; g) = Z(s, W; g)/L(s, \pi)$, of which χ is trivial, we will understand that the L-function $L(s, \pi)$ is determined associated with $|x|^{-1/2}W\left(\begin{smallmatrix} x & 0 \\ 0 & 1 \end{smallmatrix}\right)g$. In GL_1 case, we may similarly think that there exists a unique L-function $L(s, \pi)$ which is determined associated with a certain $\Phi \in \mathcal{S}(\mathbb{A}_{\mathbb{Q}})$, and that an L-function $L(s, \pi \otimes \chi)$ is given by $Z(s, \chi, \Phi)$.

Set

$$\phi^{\vee}(s, W, \chi^{-1}; g) = Z^{\vee}(s, W, \chi^{-1}; g)/L(s, \pi^{\vee} \otimes \chi^{-1}).$$

Here π^{\vee} is the contragredient representation of π and $\pi^{\vee} = \omega(\det)^{-1}\pi$. Then there exists a unique exponential function $\varepsilon(s, \pi, \chi, \psi)$ such that

$$\phi^{\vee}(1-s, W, \chi^{-1}; \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}g) = \varepsilon(s, \pi, \chi, \psi) \phi(s, W, \chi; g).$$

We shall think of the cuspidal automorphic representation of $GL_2(\mathbb{A}_{\mathbb{Q}})$. We may think that a right regular representation $(V, L^2(GL_2(\mathbb{Q}) \backslash GL_2(\mathbb{A}_{\mathbb{Q}}), d^*g))$ is given.

Theorem 4.1. Let π be a cuspidal automorphic representation of $GL_2(\mathbb{A}_{\mathbb{Q}})$. One obtains $\pi \simeq \otimes_p \pi_p$. We will think of the case where χ is trivial.

(1) The L-function $L(s, \pi)$ has the Euler product:

$$L(s, \pi) = \prod_p L(s, \pi_p).$$

(2) There exists a exponential function $\varepsilon(s, \pi, \psi)$, and the functional equation:

$$L(s, \pi) = \varepsilon(s, \pi, \psi) L(1-s, \pi^{\vee})$$

is satisfied.

Proposition 4.2. Let π be a cuspidal automorphic representation of $GL_2(A_{\mathbb{Q}})$. It has its Whittaker model.

Recall the sequence

$$0 \rightarrow L^2(X)_0 \xrightarrow{T} L^2(C_{\mathbb{Q}}) \rightarrow \mathcal{H} \rightarrow 0.$$

We may say that

$$|x|^{-1/2} W\left(\begin{pmatrix} x & 0 \\ 0 & 1 \end{pmatrix} g\right) \in L^2(X, dx)_0.$$

Think of the pairing $\langle Tf, \eta \rangle$ for $f \in L^2(X, dx)_0$ and $\eta \in (L^2(C_{\mathbb{Q}}))^*$. Then,

$$\langle Tf, \eta \rangle = \langle W\left(\begin{pmatrix} a & 0 \\ 0 & 1 \end{pmatrix} g\right), \eta \rangle = \int_{-\infty}^{\infty} \int_{A_{\mathbb{Q}}^*} W\left(\begin{pmatrix} a & 0 \\ 0 & 1 \end{pmatrix} g\right) \chi(a) |a|^it \hat{\Psi}(t) d^* a dt.$$

This computation makes us to say that

$$\eta \in \text{Im}(T)^0 \iff Z(1/2+it, W, \chi; g) = 0 \iff L(1/2+it, \pi) = 0 \quad t \in \mathbb{R}.$$

Therefore, also in GL_2 case, we can give the same spectral interpretation of critical zeros of $L(s, \pi)$.

Expect that $\{|x|^{-1/2} W\left(\begin{pmatrix} x & 0 \\ 0 & 1 \end{pmatrix} g\right) \mid W \in W(\pi, \psi)\}$ are dense in $L^2(X, dx)_0$. Then we may say that $GL_2(\mathbb{Q}) \backslash GL_2(A_{\mathbb{Q}})$ has no complementary series representation, so if $Z(s, W, \chi; g) = 0$ then $s = 1/2+it$. This must accomplish the spectral interpretation of critical zeros of $L(\chi, s)$, and we can confirm the Riemann hypothesis.

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