Undulating Relativity

Author: Alfredo Dimas Moreira Garcia
E-mail: avaliac@sjc.sp.gov.br

ABSTRACT

The Special Theory of Relativity takes us to two results that presently are considered “inexplicable” to many renowned scientists, to know:

-The dilatation of time, and
-The contraction of the Lorentz Length.

The solution to these have driven the author to the development of the Undulating Relativity (UR) theory, where the Temporal variation is due to the differences on the route of the light propagation and the lengths are constants between two landmarks in uniform relative movement.

The Undulating Relativity provides transformations between the two landmarks that differs from the transformations of Lorentz for: Space (x,y,z), Time (t), Speed ($\vec{u}$), Acceleration ($\vec{a}$), Energy (E), Momentum ($\vec{p}$), Force ($\vec{F}$), Electrical Field ($\vec{E}$), Magnetic Field ($\vec{B}$), Light Frequency (y), Electrical Current ($\vec{J}$) and “Electrical Charge” (ρ).

From the analysis of the development of the Undulating Relativity, the following can be synthesized:

- It is a theory with principles completely on physics;
- The transformations are linear;
- Keeps untouched the Euclidian principles;
- Considers the Galileo’s transformation distinct on each referential;
- Ties the Speed of Light and Time to a unique phenomenon;
- The Lorentz force can be attained by two distinct types of Filed Forces, and
- With the absence of the spatial contraction of Lorentz, to reach the same classical results of the special relativity rounding is not necessary as concluded on the Doppler effect.

Both, the Undulating Relativity and the Special Relativity of Albert Einstein explain the experience of Michel-Morley, the longitudinal and transversal Doppler effect, and supplies exactly identical formulation to:

Aberration of zenith $\Rightarrow$ $tg\alpha = \frac{v}{c} / \sqrt{1 - \frac{v^2}{c^2}}$.

Fresnel’s formula $\Rightarrow$ $c' = \frac{c}{n} + v(1 - \frac{1}{n^2})$.

Mass (m) with velocity (v) = [resting mass (mo)]/ $\sqrt{1 - \frac{v^2}{c^2}}$.

$E = m.c^2$.

Momentum $\Rightarrow \vec{p} = \frac{mo.\vec{v}}{\sqrt{1 - \frac{v^2}{c^2}}}$.

Relation between momentum (p) and Energy (E) $\Rightarrow E = c.\sqrt{mo^2.c^2 + p^2}$.

Relation between the electric field ($\vec{E}$) and the magnetic field ($\vec{B}$) $\Rightarrow$ $\vec{B} = \frac{\vec{V}}{c^2} \times \vec{E}$.

Biot-Savant’s formula $\Rightarrow \vec{B} = \frac{\mu o.I}{2.\pi.R}\vec{\mu}$.

Louis De Broglie’s wave equation $\Rightarrow \psi(x,t) = a.sin\left[2\pi y\left(t - \frac{x}{u}\right)\right]; u = \frac{c^2}{v}$

Along with the equations of transformations between two references of the UR, we get the invariance of shape to Maxwell's equations, such as:

$$\Rightarrow div\vec{E} = \frac{\rho}{\varepsilon o}; \Rightarrow div\vec{E} = 0.$$
$$\Rightarrow div\vec{B} = 0.$$
$$\Rightarrow Rot\vec{E} = \frac{-\partial\vec{B}}{\partial t}.$$
$$\Rightarrow Rot\vec{B} = \mu o.\vec{j} + \varepsilon o.\mu o.\frac{\partial\vec{E}}{\partial t}; \Rightarrow Rot\vec{B} = \varepsilon o.\mu o.\frac{\partial\vec{E}}{\partial t}.$$

We also get the invariance of shape to the equation of wave and equation of continuity under differential shape:

$$\Rightarrow \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} - \frac{1}{c^2}\frac{\partial^2}{\partial t^2} = 0$$
$$\Rightarrow \frac{\partial\rho}{\partial t} + \nabla.\vec{J} = 0.$$

Other Works:

§9 Explaining the Sagnac Effect with the Undulating Relativity

§10 Explaining the experience of Ives-Stilwell with the Undulating Relativity

§11 Transformation of the power of a luminous ray between two referencials in the Special Theory of Relativity

§12 Linearity

§13 Richard C. Tolman.

§14 Velocities composition

§15 Invariance

§16 Time and Frequency

§17 Transformation of H. Lorentz

§18 The Michelson & Morley experience

§19 Regression of the perihelion of Mercury of 7,13”

§§19 Advance of Mercury's perihelion of 42.79”

§20 Inertia

§21 Advance of Mercury's perihelion of 42.79” calculated with the Undulating Relativity

§22 Spatial Deformation

§23 Space and Time Bend

§24 Variational Principle

Author: Alfredo Dimas Moreira Garcia.
e-mail: avaliac@sjc.sp.gov.br

Translator: Rodolfo Marcos Venâncio
e-mail: sooter@terra.com.br

Undulating Relativity

§ 1 Transformation to space and time

The Undulating Relativity (UR) keep the principle of the relativity and the principle of Constancy of light speed, exactly like Albert Einstein's Special Relativity Theory defined:

a) The laws, under which the state of physics systems are changed are the same, either when referred to a determined system of coordinates or to any other that has uniform translation movement in relation to the first.

b) Any ray of light moves in the resting coordinates system with a determined velocity c, that is the same, whatever this ray is emitted by a resting body or by a body in movement (which explains the experience of Michel-Morley).

Let's imagine first that two observers O and O' (in vacuum), moving in uniform translation movement in relation to each other, that is, the observer don't rotate relatively to each other. In this way, the observer O together with the axis x, y, and z of a system of a rectangle Cartesian coordinates, sees the observer O' move with velocity v, on the positive axis x, with the respective parallel axis and sliding along with the x axis while the O', together with the x', y' and z' axis of a system of a rectangle Cartesian coordinates sees O moving with velocity –v', in negative direction towards the x' axis with the respective parallel axis and sliding along with the x' axis. The observer O measures the time t and the O' observer measures the time t' ($t \neq t'$). Let's admit that both observers set their clocks in such a way that, when the coincidence of the origin of the coordinated system happens t = t' = zero.

In the instant that t = t' = 0, a ray of light is projected from the common origin to both observers. After the time interval t the observer O will notice that his ray of light had simultaneously hit the coordinates point A (x, y, z) with the ray of the O' observer with velocity c and that the origin of the system of the O' observer has run the distance v t along the positive way of the x axis, concluding that:

$$x^2 + y^2 + z^2 - c^2 t^2 = 0 \qquad 1.1$$

$$x' = x - v\,t. \qquad 1.2$$

The same way after the time interval t' the O' observer will notice that his ray of light simultaneously hit with the observer O the coordinate point A (x', y', z') with velocity c and that the origin of the system for the observer O has run the distance v't' on the negative way of the axis x', concluding that:

$$x'^2 + y'^2 + z'^2 - c^2 t'^2 = 0 \qquad 1.3$$

$$x = x' + v'\,t'. \qquad 1.4$$

Making 1.1 equal to 1.3 we have

$$x^2 + y^2 + z^2 - c^2 t^2 = x'^2 + y'^2 + z'^2 - c^2 t'^2. \qquad 1.5$$

Because of the symmetry y = y' end z = z', that simplify 1.5 in

$$x^2 - c^2 t^2 = x'^2 - c^2 t'^2. \qquad 1.6$$

To the observer O x' = x – v t (1.2) that applied in 1.6 supplies

$x^2 - c^2 t^2 = (x - v\,t)^2 - c^2 t'^2$ from where

$$t' = t\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2 t}}\ . \qquad 1.7$$

To the observer O' x = x' + v' t' (1.4) that applied in 1.6 supplies

$(x' + v'\,t')^2 - c^2 t^2 = x'^2 - c^2 t'^2$ from where

$$t = t'\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}}\,. \qquad 1.8$$

Table I, transformations to the space and time

x' = x – v t	1.2	x = x' + v' t'	1.4
y' = y	1.2.1	y = y'	1.4.1
z' = z	1.2.2	z = z'	1.4.2
$t'=t\sqrt{1+\frac{v^2}{c^2}-\frac{2vx}{c^2t}}$	1.7	$t=t'\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}}$	1.8

From the equation system formed by 1.2 and 1.4 we find

v t = v' t' or $|v|t=|v'|t'$ (considering t>o e t'>0) 1.9

what demonstrates the invariance of the space in the Undulatory Relatitivy.

From the equation system formed by 1.7 and 1.8 we find

$$\sqrt{1+\frac{v^2}{c^2}-\frac{2vx}{c^2t}}\,.\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}} = 1. \qquad 1.10$$

If in 1.2 x' = 0 then x = v t, that applied in 1.10 supplies,

$$\sqrt{1-\frac{v^2}{c^2}}\,.\sqrt{1+\frac{v'^2}{c^2}} = 1. \qquad 1.11$$

If in 1.10 x = ct and x' = c t' then

$$\left(1-\frac{v}{c}\right).\left(1+\frac{v'}{c}\right)=1. \qquad 1.12$$

To the observer O the principle of light speed constancy guarantees that the components ux, uy and uz of the light speed are also constant along its axis, thus

$$\frac{x}{t}=\frac{dx}{dt}=ux,\frac{y}{t}=\frac{dy}{dt}=uy,\frac{z}{t}=\frac{dz}{dt}=uz \qquad 1.13$$

and then we can write

$$\sqrt{1+\frac{v^2}{c^2}-\frac{2vx}{c^2t}} = \sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}\,. \qquad 1.14$$

With the use of 1.7 and 1.9 and 1.14 we can write

$$\frac{|v|}{|v'|}=\frac{t'}{t} = \sqrt{1+\frac{v^2}{c^2}-\frac{2vx}{c^2t}} = \sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}\,. \qquad 1.15$$

Differentiating 1.9 with constant v and v', or else, only the time varying we have

$$|v|dt=|v'|dt' \text{ or } \frac{|v|}{|v'|}=\frac{dt'}{dt}, \qquad 1.16$$

but from 1.15 $\frac{|v|}{|v'|}=\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}$ then $dt'=dt\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}$. 1.17

Being v and v' constants, the reazons $\frac{|v|}{|v'|}$ and $\frac{t'}{t}$ in 1.15 must also be constant because fo this the differential of $\sqrt{1+\frac{v^2}{c^2}-\frac{2vx}{c^2t}}$ must be equal to zero from where we conclude $\frac{x}{t}=\frac{dx}{dt}=ux$, that is exactly the same as 1.13.

To the observer O' the principle of Constancy of velocity of light guarantees that the components u'x', u'y', and u'z' of velocity of light are also constant alongside its axis, thus

$$\frac{x'}{t'}=\frac{dx'}{dt'}=u'\,x'\,,\frac{y'}{t'}=\frac{dy'}{dt'}=u'\,y'\,,\frac{z'}{t'}=\frac{dz'}{dt'}=u'\,z'\,, \qquad 1.18$$

and with this we can write ,

$$\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}} = \sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}\,. \qquad 1.19$$

With the use of 1.8, 1.9, and 1.19 we can write

$$\frac{|v'|}{|v|}=\frac{t}{t'} = \sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}} = \sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}\,. \qquad 1.20$$

Differentiating 1.9 with v' and v constant, that is, only the time varying we have

$|v'|dt'=|v|dt$ or $\frac{|v'|}{|v|}=\frac{dt}{dt'}$, 1.21

but from 1.20 $\frac{|v'|}{|v|}=\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}$ then $dt=dt'\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}$. 1.22

Being v' and v constant the divisions $\frac{|v'|}{|v|}$ and $\frac{t}{t'}$ in 1.20 also have to be constant because of this the differential of $\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}}$ must be equal to zero from where we conclude $\frac{x'}{t'}=\frac{dx'}{dt'}=u'x'$, that is exactly like to 1.18.

Replacing 1.14 and 1.19 in 1.10 we have

$$\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}\,.\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}} = 1. \qquad 1.23$$

To the observer O the vector position of the point A of coordinates (x,y,z) is

$\vec{R}=x\vec{i}+y\vec{j}+z\vec{k}$, 1.24

and the vector position of the origin of the system of the observer O' is

$\vec{R}o'=vt\vec{i}+0\vec{j}+0\vec{k}\Rightarrow\vec{R}o'=vt\vec{i}$. 1.25

To the observer O', the vector position of the point A of coordinates (x',y',z') is

$\vec{R}'=x'\vec{i}+y'\vec{j}+z'\vec{k}$, 1.26

and the vector position of the origin of the system of the observer O is

$\vec{R}'o=-v't'\vec{i}+0\vec{j}+0\vec{k}\Rightarrow\vec{R}'o=-v't'\vec{i}$. 1.27

Due to 1.9, 1.25, and 1.27 we have, $\vec{R}o'=-\vec{R}'o$. 1.28

As 1.24 is equal to 1.25 plus 1.26 we have

$$\vec{R}=\vec{R}o'+\vec{R}'\Rightarrow\vec{R}'=\vec{R}-\vec{R}o'. \qquad 1.29$$

Applying 1.28 in 1.29 we have, $\vec{R}=\vec{R}'-\vec{R}'o$. 1.30

To the observer O the vector velocity of the origin of the system of the observer O' is

$$\vec{v} = \frac{d\vec{R}o'}{dt} = v\vec{i} + 0\vec{j} + 0\vec{k} \Rightarrow \vec{v} = v\vec{i} \ . \qquad 1.31$$

To the observer O' the vector velocity of the origin of the system of the observer O is

$$\vec{v}' = \frac{d\vec{R}'o}{dt'} = -v'\vec{i} + 0\vec{j} + 0\vec{k} \Rightarrow \vec{v}' = -v'\vec{i} \ . \qquad 1.32$$

From 1.15, 1.20, 1.31, and 1.32 we find the following relations between $\vec{v}$ and $\vec{v}'$

$$\vec{v} = \frac{-\vec{v}'}{\sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'u'x'}{c^2}}} \qquad 1.33$$

$$\vec{v}' = \frac{-\vec{v}}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}}} \ . \qquad 1.34$$

Observation: in the table I the formulas 1.2, 1.2.1, and 1.2.2 are the components of the vector 1.29 and the formulas 1.4, 1.4.1, and 1.4.2 are the components of the vector 1.30.

§2 Law of velocity transformations $\vec{u}$ and $\vec{u}'$

Differentiating 1.29 and dividing it by 1.17 we have

$$\frac{d\vec{R}'}{dt'} = \frac{d\vec{R} - d\vec{R}o'}{dt\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}}} \Rightarrow \vec{u}' = \frac{\vec{u} - \vec{v}}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}}} = \frac{\vec{u} - \vec{v}}{\sqrt{K}} \ . \qquad 2.1$$

Differentiating 1.30 and dividing it by 1.22 we have

$$\frac{d\vec{R}}{dt} = \frac{d\vec{R}' - d\vec{R}'o}{dt'\sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'u'x'}{c^2}}} \Rightarrow \vec{u} = \frac{\vec{u}' - \vec{v}'}{\sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'u'x'}{c^2}}} = \frac{\vec{u}' - \vec{v}'}{\sqrt{K'}} \ . \qquad 2.2$$

Table 2, Law of velocity transformations $\vec{u}$ and $\vec{u}'$

$\vec{u}' = \frac{\vec{u} - \vec{v}}{\sqrt{K}}$	2.1	$\vec{u} = \frac{\vec{u}' - \vec{v}'}{\sqrt{K'}}$	2.2
$u'x' = \frac{ux - v}{\sqrt{K}}$	2.3	$ux = \frac{u'x' + v'}{\sqrt{K'}}$	2.4
$u'y' = \frac{uy}{\sqrt{K}}$	2.3.1	$uy = \frac{u'y'}{\sqrt{K'}}$	2.4.1
$u'z' = \frac{uz}{\sqrt{K}}$	2.3.2	$uz = \frac{u'z'}{\sqrt{K'}}$	2.4.2
$\lvert v' \rvert = \frac{\lvert v \rvert}{\sqrt{K}}$	1.15	$v = \frac{\lvert v' \rvert}{\sqrt{K'}}$	1.20
$\sqrt{K} = \sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}}$	2.5	$\sqrt{K'} = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'u'x'}{c^2}}$	2.6

Multiplying 2.1 by itself we have

$$u'=\frac{u\sqrt{1+\frac{v^2}{u^2}-\frac{2vux}{u^2}}}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}}. \qquad 2.7$$

If in 2.7 we make u = c then u’ = c as it is required by the principle of constancy of velocity of light.
Multiplying 2.2 by itself we have

$$u=\frac{u'\sqrt{1+\frac{v'^2}{u'^2}+\frac{2v'u'x'}{u'^2}}}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}}. \qquad 2.8$$

If in 2.8 we make u’ = c then u = c as it is required by the principle of constancy of velocity of light.

If in 2.3 we make ux = c then $u'x'=\frac{c-v}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vc}{c^2}}}=c$ as it is required by the principle of constancy of velocity of light.

If in 2.4 we make u’x’ = c then $ux=\frac{c+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'c}{c^2}}}=c$ as it is required by the principle of constancy of velocity of light.

Remodeling 2.7 and 2.8 we have

$$\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}} = \frac{\sqrt{1-\frac{u^2}{c^2}}}{\sqrt{1-\frac{u'^2}{c^2}}}. \qquad 2.9$$

$$\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}} = \frac{\sqrt{1-\frac{u'^2}{c^2}}}{\sqrt{1-\frac{u^2}{c^2}}}. \qquad 2.10$$

The direct relations between the times and velocities of two points in space can be obtained with the equalities $\vec{u}'=0 \Rightarrow u'x'=0 \Rightarrow ux=v$ coming from 2.1, that applied in 1.17, 1.22, 1.20, and 1.15 supply

$$dt'=dt\sqrt{1+\frac{v^2}{c^2}-\frac{2vv}{c^2}} \Rightarrow dt=\frac{dt'}{\sqrt{1-\frac{v^2}{c^2}}}, \qquad 2.11$$

$$dt=dt'\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'0}{c^2}} \Rightarrow dt'=\frac{dt}{\sqrt{1+\frac{v'^2}{c^2}}}, \qquad 2.12$$

$$|v|=\frac{|v'|}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'0}{c^2}}} \Rightarrow |v|=\frac{|v'|}{\sqrt{1+\frac{v'^2}{c^2}}}, \qquad 2.13$$

$$|v'|=\frac{|v|}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vv}{c^2}}} \Rightarrow |v'|=\frac{|v|}{\sqrt{1-\frac{v^2}{c^2}}}. \qquad 2.14$$

Aberration of the zenith

To the observer O' along with the star u'x' = 0, u'y' = c and u'z' = 0, and to the observer O along with the Earth we have the conjunct 2.3

$$0=\frac{ux-v}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}}\Rightarrow ux=v\,,\,c=\frac{uy}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vv}{c^2}}}\Rightarrow uy=c\sqrt{1-\frac{v^2}{c^2}}$$, uz = 0,

$u=\sqrt{ux^2+uy^2+uz^2}=\sqrt{v^2+\left(c\sqrt{1-\frac{v^2}{c^2}}\right)^2+0^2}=c$ exactly as foreseen by the principle of relativity.

To the observer O the light propagates in a direction that makes an angle with the vertical axis y given by

$$tang\alpha=\frac{ux}{uy}=\frac{v}{c.\sqrt{1-\frac{v^2}{c^2}}}=\frac{v/c}{\sqrt{1-\frac{v^2}{c^2}}} \tag{2.15}$$

that is the aberration formula of the zenith in the special relativity .
If we inverted the observers we would have the conjunct 2.4

$$0=\frac{u'x'+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}}\Rightarrow u'x'=-v'\,,\,c=\frac{u'y'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'(-v')}{c^2}}}\Rightarrow u'y'=c\sqrt{1-\frac{v'^2}{c^2}}$$, u'z'=0,

$$u'=\sqrt{u'x'^2+u'y'^2+u'z'^2}=\sqrt{(-v')^2+\left(c\sqrt{1-\frac{v'^2}{c^2}}\right)^2+0^2}=c$$

$$tang\alpha=\frac{u'x'}{u'y'}=\frac{-v'}{c.\sqrt{1-\frac{v'^2}{c^2}}}=\frac{-v'/c}{\sqrt{1-\frac{v'^2}{c^2}}} \tag{2.16}$$

that is equal to 2.15, with the negative sign indicating the contrary direction of the angles.

Fresnel's formula

Considering in 2.4, $u'x'=c/n$ the velocity of light relativily to the water, $v'=v$ the velocity of water in relation to the apparatus then $ux=c'$ will be the velocity of light relatively to the laboratory

$$c'=\frac{c/n+v}{\sqrt{1+\frac{v^2}{c^2}+\frac{2vc/n}{c^2}}}=\frac{c/n+v}{\sqrt{1+\frac{v^2}{c^2}+\frac{2v}{nc}}}=\left(\frac{c}{n}+v\right)\left(1+\frac{v^2}{c^2}+\frac{2v}{nc}\right)^{-\frac{1}{2}}\cong\left(\frac{c}{n}+v\right)\left[1-\frac{1}{2}\left(\frac{v^2}{c^2}+\frac{2v}{nc}\right)\right]$$

Ignoring the term v^2/c^2 we have

$$c'\cong\left(\frac{c}{n}+v\right)\left(1-\frac{v}{nc}\right)\cong\frac{c}{n}+v-\frac{v}{n^2}-\frac{v^2}{nc}$$

and ignoring the term v^2/nc we have the Fresnel's formula

$$c'=\frac{c}{n}+v-\frac{v}{n^2}=\frac{c}{n}+v\left(1-\frac{1}{n^2}\right). \tag{2.17}$$

Doppler effect

Making $r^2 = x^2 + y^2 + z^2$ and $r'^2 = x'^2 + y'^2 + z'^2$ in 1.5 we have $r^2 - c^2t^2 = r'^2 - c^2t'^2$ or $(r-ct) = (r'-ct')\dfrac{(r'+ct')}{(r+ct)}$ replacing then $r = ct$, $r' = ct'$ and 1.7 we find $(r-ct) = (r'-ct')\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vx}{c^2t}}$ as $c = \dfrac{w}{k} = \dfrac{w'}{k'}$ then $\dfrac{1}{k}(kr-wt) = \dfrac{1}{k'}(k'r'-w't')\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vx}{c^2t}}$ where to attend the principle of relativity we will define $k' = k\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vx}{c^2t}}$ 2.18

Resulting in the expression $(kr-wt) = (k'r'-w't')$ symmetric and invariable between the observers.

To the observer O an expression in the formula of $\psi(r,t) = f(kr-wt)$ 2.19

represents a curve that propagates in the direction of $\vec{R}$. To the observer O' an expression in the formula of $\psi'(r',t') = f'(k'r'-w't')$ 2.20

represents a curve that propates in the direction of $\vec{R}'$.

Applying in 2.18 $k = \dfrac{2\pi}{\lambda}$, $k' = \dfrac{2\pi}{\lambda'}$, 1.14, 1.19, 1.23, 2.5, and 2.6 we have

$\lambda' = \dfrac{\lambda}{\sqrt{K}}$ e $\lambda = \dfrac{\lambda'}{\sqrt{K'}}$, 2.21

that applied in $c = y\lambda = y'\lambda'$ supply, $y' = y\sqrt{K}$ and $y = y'\sqrt{K'}$. 2.22

Considering the relation of Planck-Einstein between energy (E) and frequency (y), we have to the observer O $E = hy$ and to the observer O' $E' = hy'$ that replaced in 2.22 supply

$E' = E\sqrt{K}$ and $E = E'\sqrt{K'}$. 2.23

If the observer O that sees the observer O' moving with velocity v in a positive way to the axis x, emits waves of frequency y and velocity c in a positive way to the axis x then, according to 2.22 and $ux = c$ the observer O' will measure the waves with velocity c and frequency $y' = y\left(1-\dfrac{v}{c}\right)$, 2.24

that is exactly the classic formula of the longitudinal Doppler effect.

If the observer O' that sees the observer O moving with velocity –v' in the negative way of the axis x', emits waves of frequency y' and velocity c, then the observer O according to 2.22 and $u'x' = -v'$ will measure waves of frequency y and velocity c in a perpendicular plane to the movement of O' given by

$\gamma = \gamma'\sqrt{1-\dfrac{v'^2}{c^2}}$, 2.25

that is exactly the formula of the transversal Doppler effect in the Special Relativity.

§3 Transformations of the accelerations $\vec{a}$ and $\vec{a}'$

Differentiating 2.1 and dividing it by 1.17 we have

$$\frac{d\vec{u}'}{dt'} = \frac{d\vec{u}/\sqrt{K}}{dt\sqrt{K}} + (\vec{u}-\vec{v})\frac{v}{c^2}\frac{dux/K\sqrt{K}}{dt\sqrt{K}} \Rightarrow \vec{a}' = \frac{\vec{a}}{K} + (\vec{u}-\vec{v})\frac{v}{c^2}\frac{ax}{K^2}. \quad 3.1$$

Differentiating 2.2 and dividing it by 1.22 we have

$$\frac{d\vec{u}}{dt} = \frac{d\vec{u}'/\sqrt{K'}}{dt'\sqrt{K'}} - (\vec{u}'-\vec{v}')\frac{v'}{c^2}\frac{du'x'/K'\sqrt{K'}}{dt'\sqrt{K'}} \Rightarrow \vec{a} = \frac{\vec{a}'}{K'} - (\vec{u}'-\vec{v}')\frac{v'}{c^2}\frac{a'x'}{K'^2}. \quad 3.2$$

Table 3, transformations of the accelerations $\vec{a}$ and $\vec{a}'$

$\vec{a}'=\frac{\vec{a}}{K}+(\vec{u}-\vec{v})\frac{v}{c^2}\frac{ax}{K^2}$	3.1	$\vec{a}=\frac{\vec{a}'}{K'}-(\vec{u}'-\vec{v}')\frac{v'}{c^2}\frac{a'x'}{K'^2}$	3.2
$a'x'=\frac{ax}{K}+(ux-v)\frac{v}{c^2}\frac{ax}{K^2}$	3.3	$ax=\frac{a'x'}{K'}-(u'x'+v')\frac{v'}{c^2}\frac{a'x'}{K'^2}$	3.4
$a'y'=\frac{ay}{K}+uy\frac{v}{c^2}\frac{ax}{K^2}$	3.3.1	$ay=\frac{a'y'}{K'}-u'y'\frac{v'}{c^2}\frac{a'x'}{K'^2}$	3.4.1
$a'z'=\frac{az}{K}+uz\frac{v}{c^2}\frac{ax}{K^2}$	3.3.2	$az=\frac{a'z'}{K'}-u'z'\frac{v'}{c^2}\frac{a'x'}{K'^2}$	3.4.2
$a'=\frac{a}{K}$	3.8	$a=\frac{a'}{K'}$	3.9
$K=1+\frac{v^2}{c^2}-\frac{2vux}{c^2}$	3.5	$K'=1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}$	3.6

From the tables 2 and 3 we can conclude that if to the observer O $\vec{u}.\vec{a}=zero$ and $c^2=ux^2+uy^2+uz^2$, then it is also to the observer O' $\vec{u}'.\vec{a}'=zero$ and $c^2=u'x'^2+u'y'^2+u'z'^2$, thus $\vec{u}$ is perpendicular to $\vec{a}$ and $\vec{u}'$ is perpendicular to $\vec{a}'$ as the vectors theory requires.

Differentiating 1.9 with the velocities and the times changing we have, $tdv+vdt=t'dv'+v'dt'$, but considering 1.16 we have, $vdt=v'dt'\Rightarrow tdv=t'dv'$ 3.7

Where replacing 1.15 and dividing it by 1.17 we have, $\frac{dv'}{dt'}=\frac{dv}{dtK}$ or $a'=\frac{a}{K}$. 3.8

We can also replace 1.20 in 3.7 and divide it by 1.22 deducing

$\frac{dv}{dt}=\frac{dv'}{dt'K'}$ or $a=\frac{a'}{K'}$. 3.9

The direct relations between the modules of the accelerations a and a' of two points in space can be obtained with the $\vec{u}'=0\Rightarrow u'x'=0\Rightarrow a'x'=0\Rightarrow\vec{u}=\vec{v}\Rightarrow ux=v$ coming from 2.1, that applied in 3.8 and 3.9 supply

$a'=\frac{a}{1+\frac{v^2}{c^2}-\frac{2vv}{c^2}}=\frac{a}{1-\frac{v^2}{c^2}}$ and $a=\frac{a'}{1+\frac{v'^2}{c^2}+\frac{2v'0}{c^2}}=\frac{a'}{1+\frac{v'^2}{c^2}}$. 3.10

That can also be reduced from 3.1 and 3.2 if we use the same equalities
$\vec{u}'=0\Rightarrow u'x'=0\Rightarrow a'x'=0\Rightarrow\vec{u}=\vec{v}\Rightarrow ux=v$ coming from 2.1.

§4 Transformations of the Moments $\vec{p}$ and $\vec{p}'$

Defined as $\vec{p}=m(u)\vec{u}$ and $\vec{p}'=m'(u')\vec{u}'$, 4.1

where $m(u)$ and $m'(u')$ symbolizes the function masses of the modules of velocities $u=|\vec{u}|$ and $u'=|\vec{u}'|$.

We will have the relations between $m(u)$ and $m'(u')$ and the resting mass m_o, analyzing the elastic collision in a plane between the sphere s that for the observer o moves alongside the axis y with velocity uy = w and the sphere s' that for the observer O' moves alongside the axis y' with velocity u'y' = -w. The spheres while observed in relative resting are identical and have the mass m_o. The considered collision is symmetric in relation to a parallel line to the axis y and y' passing by the center of the spheres in the moment of. Collision.

Before and after the collision the spheres have velocities observed by O and O' according to the following table gotten from table 2

	Sphere	Observer O	Observer O'
Before	s	$uxs = zero$, $uys = w$	$u'x's = -v'$, $u'y's = w\sqrt{1-\frac{v'^2}{c^2}}$
Collision	s'	$uxs' = v$, $uys' = -w\sqrt{1-\frac{v^2}{c^2}}$	$u'x's' = zero$, $u'y's' = -w$
After	s	$uxs = zero$, $uys = -\overline{w}$	$u'x's = -v'$, $u'y's = -\overline{w}\sqrt{1-\frac{v'^2}{c^2}}$
Collision	s'	$uxs' = v$, $uys' = \overline{w}\sqrt{1-\frac{v^2}{c^2}}$	$u'x's' = zero$, $u'y's' = \overline{w}$

To the observer O, the principle of conservation of moments establishes that the moments $px = m(u)ux$ and $py = m(u)uy$, of the spheres s and s' in relation to the axis x and y, remain constant before and after the collision thus for the axis x we have

$$m\left(\sqrt{uxs^2+uys^2}\right)uxs+m\left(\sqrt{uxs'^2+uys'^2}\right)uxs' = m\left(\sqrt{uxs^2+uys^2}\right)uxs+m\left(\sqrt{uxs'^2+uys'^2}\right)uxs' ,$$

where replacing the values of the table we have

$$m\left(\sqrt{v^2+\left(-w\sqrt{1-\frac{v^2}{c^2}}\right)^2}\right)v = m\left(\sqrt{v^2+\left(\overline{w}\sqrt{1-\frac{v^2}{c^2}}\right)^2}\right)v$$ from where we conclude that $\overline{w} = w$,

and for the axis y

$$m\left(\sqrt{uxs^2+uys^2}\right)uys+m\left(\sqrt{uxs'^2+uys'^2}\right)uys' = m\left(\sqrt{uxs^2+uys^2}\right)uys+m\left(\sqrt{uxs'^2+uys'^2}\right)uys' ,$$

where replacing the values of the table we have

$$m(w)w-m\left(\sqrt{v^2+\left(-w\sqrt{1-\frac{v^2}{c^2}}\right)^2}\right)w\sqrt{1-\frac{v^2}{c^2}} = -m(\overline{w})\overline{w}+m\left(\sqrt{v^2+\left(\overline{w}\sqrt{1-\frac{v^2}{c^2}}\right)^2}\right)\overline{w}\sqrt{1-\frac{v^2}{c^2}} ,$$

simplifying we have

$$m(w) = m\left(\sqrt{v^2+w^2\left(1-\frac{v^2}{c^2}\right)}\right)\sqrt{1-\frac{v^2}{c^2}}$$, where when $w \to 0$ becomes

$$m(0) = m\left(\sqrt{v^2+0^2\left(1-\frac{v^2}{c^2}\right)}\right)\sqrt{1-\frac{v^2}{c^2}} \Rightarrow m(0) = m(v)\sqrt{1-\frac{v^2}{c^2}} \Rightarrow m(v) = \frac{m(0)}{\sqrt{1-\frac{v^2}{c^2}}} ,$$

but $m(0)$ is equal to the resting mass m_o thus

$$m(v) = \frac{m_0}{\sqrt{1-\frac{v^2}{c^2}}}$$, with a relative velocity $v = u \Rightarrow m(u) = \frac{m_0}{\sqrt{1-\frac{u^2}{c^2}}}$ 4.2

that applied in 4.1 supplies $\vec{p} = m(u)\vec{u} = \frac{m_0\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}$. 4.1

With the same procedures we would have for the O' observer

$$m'(u') = \frac{m_0}{\sqrt{1-\frac{u'^2}{c^2}}}$$ 4.3

and $\vec{p}' = m'(u')\vec{u}' = \frac{m_0\vec{u}'}{\sqrt{1-\frac{u'^2}{c^2}}}$. 4.1

Simplifying the simbology we will adopt $m = m(u) = \frac{m_0}{\sqrt{1-\frac{u^2}{c^2}}}$ 4.2

and $m' = m'(u') = \frac{m_0}{\sqrt{1-\frac{u'^2}{c^2}}}$ 4.3

that simplify the moments in $\vec{p} = m\vec{u}$ and $\vec{p}' = m'\vec{u}'$. 4.1

Applying 4.2 and 4.3 in 2.9 and 2.10 we have

$$m = m'\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}} \Rightarrow m = m'\sqrt{K'} \text{ and } m' = m\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}} \Rightarrow m' = m\sqrt{K}\ . \qquad 4.4$$

Defining force as Newton we have $\vec{F} = \frac{d\vec{p}}{dt} = \frac{d(m\vec{u})}{dt}$ and $\vec{F}' = \frac{d\vec{p}'}{dt'} = \frac{d(m'\vec{u}')}{dt'}$, with this we can define then kinetic energy (E_k, E'_k) as

$$E_k = \int_0^u \vec{F}.d\vec{R} = \int_0^u \frac{d(m\vec{u})}{dt}.d\vec{R} = \int_0^u d(m\vec{u}).\vec{u} = \int_0^u (u^2 dm + mudu),$$

and $E'_k = \int_0^{u'} \vec{F}'.d\vec{R}' = \int_0^{u'} \frac{d(m'\vec{u}')}{dt'}.d\vec{R}' = \int_0^{u'} d(m'\vec{u}').\vec{u}' = \int_0^{u'} (u'^2\, dm' + m'u'du')$.

Remodeling 4.2 and 4.3 and differentiating we have $m^2c^2 - m^2u^2 = m_0^{\ 2}c^2 \Rightarrow u^2dm + mudu = c^2dm$ and $m'^2c^2 - m'^2u'^2 = m_0^{\ 2}c^2 \Rightarrow u'^2\,dm' + m'u'du' = c^2dm'$, that applied in the formulas of kinetic energy supplies $E_k = \int_{m_0}^{m} c^2dm = mc^2 - m_0c^2 = E - E_0$ and $E'_k = \int_{m_0}^{m'} c^2dm' = m'c^2 - m_0c^2 = E' - E_0$, 4.5

where $E = mc^2$ and $E' = m'c^2$ 4.6

are the total energies as in the special relativity and $E_o = m_oc^2$ 4.7
the resting energy.

Applying 4.6 in 4.4 we have exactly 2.23.

From 4.6, 4.2, 4.3, and 4.1 we find

$$E = c\sqrt{m_o^{\ 2}c^2 + p^2} \text{ and } E' = c\sqrt{m_o^{\ 2}c^2 + p'^2} \qquad 4.8$$

identical relations to the Special Relativity.

Multiplying 2.1 and 2.2 by m_o we get

$$\frac{m_o\vec{u}'}{\sqrt{1-\frac{u'^2}{c^2}}}=\frac{m_o\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}-\frac{m_o\vec{v}}{\sqrt{1-\frac{u^2}{c^2}}}\Rightarrow m'\vec{u}'=m\vec{u}-m\vec{v}\Rightarrow \vec{p}'=\vec{p}-\frac{E}{c^2}\vec{v} \qquad 4.9$$

and $$\frac{m_o\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}=\frac{m_o\vec{u}'}{\sqrt{1-\frac{u'^2}{c^2}}}-\frac{m_o\vec{v}'}{\sqrt{1-\frac{u'^2}{c^2}}}\Rightarrow m\vec{u}=m'\vec{u}'-m'\vec{v}'\Rightarrow \vec{p}=\vec{p}'-\frac{E'}{c^2}\vec{v}'\ . \qquad 4.10$$

Table 4, transformations of moments $\vec{p}$ and $\vec{p}'$

$\vec{p}'=\vec{p}-\frac{E}{c^2}\vec{v}$	4.9	$\vec{p}=\vec{p}'-\frac{E'}{c^2}\vec{v}'$	4.10
$p'x'=px-\frac{E}{c^2}v$	4.11	$px=p'x'+\frac{E'}{c^2}v'$	4.12
$p'y'=py$	4.11.1	$py=p'y'$	4.12.1
$p'z'=pz$	4.11.2	$pz=p'z'$	4.12.2
$E'=E\sqrt{K}$	2.23	$E=E'\sqrt{K'}$	2.23
$m=m(u)=\frac{m_0}{\sqrt{1-\frac{u^2}{c^2}}}$	4.2	$m'=m'(u')=\frac{m_0}{\sqrt{1-\frac{u'^2}{c^2}}}$	4.3
$m'=m\sqrt{K}$	4.4	$m=m'\sqrt{K'}$	4.4
$E_k=E-E_o$	4.5	$E'_k=E'-E_o$	4.5
$E=mc^2$	4.6	$E'=m'c^2$	4.6
$E_o=m_oc^2$	4.7	$E_o=m_oc^2$	4.7
$E=c\sqrt{m_o^2c^2+p^2}$	4.8	$E'=c\sqrt{m_o^2c^2+p'^2}$	4.8

Wave equation of Louis de Broglie

The observer O' associates to a resting particle in its origin the following properties:

-Resting mass m_o

-Time $t'=t_o$

-Resting Energy $E_o=m_oc^2$

-Frequency $y_o=\frac{E_o}{h}=\frac{m_oc^2}{h}$

-Wave function $\psi_o=asen2\pi y_o t_o$ with a = constant.

The observer O associates to a particle with velocity v the following:

-Mass $m=m(v)=\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}$ (from 4.2 where $u=v$)

-Time $t=\frac{t_o}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vv}{c^2}}}=\frac{t_o}{\sqrt{1-\frac{v^2}{c^2}}}$ (from 1.7 with $ux=v$ and $t'=t_o$)

-Energy $E=\frac{E_o}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_oc^2}{\sqrt{1-\frac{v^2}{c^2}}}$ (from 2.23 with $ux=v$ and $E'=E_o$)

-Frequency $y=\dfrac{y_o}{\sqrt{1-\dfrac{v^2}{c^2}}}=\dfrac{m_o c^2/h}{\sqrt{1-\dfrac{v^2}{c^2}}}$ (from 2.22 with $ux=v$ and $y'=y_o$)

-Distance x = vt (from 1.2 with x' = 0)

-Wave function $\psi = asen2\pi y_o t_o = asen2\pi y\sqrt{1-\dfrac{v^2}{c^2}}t\sqrt{1-\dfrac{v^2}{c^2}} = asen2\pi y\left(t-\dfrac{x}{u}\right)$ with $u=\dfrac{c^2}{v}$

-Wave length $u=y\lambda=\dfrac{c^2}{v}=\dfrac{E}{p}=\dfrac{yh}{p}\Rightarrow\lambda=\dfrac{h}{p}$ (from 4.9 with $\vec{p}'=\ \vec{p}_o\ =0$)

To go back to the O' observer referential where $\vec{u}'=0\Rightarrow u'x'=0$, we will consider the following variables:

-Distance x = v't' (from 1.4 with x' = 0)

-Time $t=t'\sqrt{1+\dfrac{v'^2}{c^2}+\dfrac{2v'0}{c^2}}=t'\sqrt{1+\dfrac{v'^2}{c^2}}$ (from 1.8 with $u'x'=0$)

-Frequency $y=y'\sqrt{1+\dfrac{v'^2}{c^2}}$ (from 2.22 with $u'x'=0$)

-Velocity $v=\dfrac{v'}{\sqrt{1+\dfrac{v'^2}{c^2}}}$ (de 2.13)

that applied to the wave function supplies

$$\psi'=asen2\pi y\left(t-\frac{vx}{c^2}\right)=asen2\pi y'\sqrt{1+\frac{v'^2}{c^2}}\left(t'\sqrt{1+\frac{v'^2}{c^2}}-\frac{v'^2t'}{c^2\sqrt{1+\dfrac{v'^2}{c^2}}}\right)=asen2\pi y't',$$

but as $t'=t_o$ and $y'=y_o$ then $\psi'=\psi_o$.

§5 Transformations of the Forces $\vec{F}$ and $\vec{F}'$

Differentiating 4.9 and dividing by 1.17 we have

$$\frac{d\vec{p}'}{dt'}=\frac{d\vec{p}}{dt\sqrt{K}}-\frac{dE}{dt\sqrt{K}}\frac{\vec{v}}{c^2}\Rightarrow\vec{F}'=\frac{1}{\sqrt{K}}\left[\vec{F}-\frac{dE}{dt}\frac{\vec{v}}{c^2}\right]\Rightarrow\vec{F}'=\frac{1}{\sqrt{K}}\left[\vec{F}-\left(\vec{F}.\vec{u}\right)\frac{\vec{v}}{c^2}\right]. \qquad 5.1$$

Differentiating 4.10 and dividing by 1.22 we have

$$\frac{d\vec{p}}{dt}=\frac{d\vec{p}'}{dt'\sqrt{K'}}-\frac{dE'}{dt'\sqrt{K'}}\frac{\vec{v}'}{c^2}\Rightarrow\vec{F}=\frac{1}{\sqrt{K'}}\left[\vec{F}'-\frac{dE'}{dt'}\frac{\vec{v}'}{c^2}\right]\Rightarrow\vec{F}=\frac{1}{\sqrt{K'}}\left[\vec{F}'-\left(\vec{F}'.\vec{u}'\right)\frac{\vec{v}'}{c^2}\right]. \qquad 5.2$$

From the system formed by 5.1 and 5.2 we have

$$\frac{dE}{dt}=\frac{dE'}{dt'} \text{ or } \vec{F}.\vec{u}=\vec{F}'.\vec{u}', \qquad 5.3$$

that is an invariant between the observers in the Undulating .Relativity.

Table 5, transformations of the Forces $\vec{F}$ and $\vec{F}'$

$\vec{F}'=\frac{1}{\sqrt{K}}\left[\vec{F}-\left(\vec{F}.\vec{u}\right)\frac{\vec{v}}{c^2}\right]$	5.1	$\vec{F}=\frac{1}{\sqrt{K'}}\left[\vec{F}'-\left(\vec{F}'.\vec{u}'\right)\frac{\vec{v}'}{c^2}\right]$	5.2
$F'x'=\frac{1}{\sqrt{K}}\left[Fx-\left(\vec{F}.\vec{u}\right)\frac{v}{c^2}\right]$	5.4	$Fx=\frac{1}{\sqrt{K'}}\left[F'x'+\left(\vec{F}'.\vec{u}'\right)\frac{v'}{c^2}\right]$	5.5
$F'y'=Fy/\sqrt{K}$	5.4.1	$Fy=F'y'/\sqrt{K'}$	5.5.1
$F'z'=Fz/\sqrt{K}$	5.4.2	$Fz=F'z'/\sqrt{K'}$	5.5.2
$\frac{dE'}{dt'}=\frac{dE}{dt}$	5.3	$\vec{F}.\vec{u}=\vec{F}'.\vec{u}'$	5.3

§6 Transformations of the density of charge ρ, ρ' and density of current $\vec{J}$ and $\vec{J}'$

Multiplying 2.1 and 2.2 by the density of the resting electric charge defined as $\rho_o=\frac{dq}{dv_o}$ we have

$$\frac{\rho_o\vec{u}'}{\sqrt{1-\frac{u'^2}{c^2}}}=\frac{\rho_o\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}-\frac{\rho_o\vec{v}}{\sqrt{1-\frac{u^2}{c^2}}}\Rightarrow\rho'\vec{u}'=\rho\vec{u}-\rho\vec{v}\Rightarrow\vec{J}'=\vec{J}-\rho\vec{v} \qquad 6.1$$

and $$\frac{\rho_o\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}=\frac{\rho_o\vec{u}'}{\sqrt{1-\frac{u'^2}{c^2}}}-\frac{\rho_o\vec{v}'}{\sqrt{1-\frac{u'^2}{c^2}}}\Rightarrow\rho\vec{u}=\rho'\vec{u}'-\rho'\vec{v}'\Rightarrow\vec{J}=\vec{J}'-\rho'\vec{v}'. \qquad 6.2$$

Table 6, transformations of the density of charges ρ, ρ' and density of current $\vec{J}$ and $\vec{J}'$

$\vec{J}'=\vec{J}-\rho\vec{v}$	6.1	$\vec{J}=\vec{J}'-\rho'\vec{v}'$	6.2
$J'x'=Jx-\rho v$	6.3	$Jx=J'x'+\rho'v'$	6.4
$J'y'=Jy$	6.3.1	$Jy=J'y'$	6.4.1
$J'z'=Jz$	6.3.2	$Jz=J'z'$	6.4.2
$\vec{J}=\rho\vec{u}$	6.5	$\vec{J}'=\rho'\vec{u}'$	6.6
$\rho=\frac{\rho_o}{\sqrt{1-\frac{u^2}{c^2}}}$	6.7	$\rho'=\frac{\rho_o}{\sqrt{1-\frac{u'^2}{c^2}}}$	6.8
$\rho'=\rho\sqrt{K}$	6.9	$\rho=\rho'\sqrt{K'}$	6.10

From the system formed by 6.1 and 6.2 we had 6.9 and 6.10.

§7 Transformation of the electric fields $\vec{E}$, $\vec{E}'$ and magnetic fields $\vec{B}$, $\vec{B}'$

Applying the forces of Lorentz $\vec{F}=q\left(\vec{E}+\vec{u}\times\vec{B}\right)$ and $\vec{F}'=q\left(\vec{E}'+\vec{u}'\times\vec{B}'\right)$ in 5.1 and 5.2 we have

$$q\left(\vec{E}'+\vec{u}'\times\vec{B}'\right)=\frac{1}{\sqrt{K}}\left[q\left(\vec{E}+\vec{u}\times\vec{B}\right)-\left[q\left(\vec{E}+\vec{u}\times\vec{B}\right).\vec{u}\right]\frac{\vec{v}}{c^2}\right]$$

and $q\left(\vec{E}+\vec{u}\times\vec{B}\right)=\frac{1}{\sqrt{K'}}\left[q\left(\vec{E}'+\vec{u}'\times\vec{B}'\right)-\left[q\left(\vec{E}'+\vec{u}'\times\vec{B}'\right).\vec{u}'\right]\frac{\vec{v}'}{c^2}\right]$, that simplified become

$\left(\vec{E}'+\vec{u}'\times\vec{B}'\right)=\frac{1}{\sqrt{K}}\left[\left(\vec{E}+\vec{u}\times\vec{B}\right)-\left(\vec{E}.\vec{u}\right)\frac{\vec{v}}{c^2}\right]$ and $\left(\vec{E}+\vec{u}\times\vec{B}\right)=\frac{1}{\sqrt{K'}}\left[\left(\vec{E}'+\vec{u}'\times\vec{B}'\right)-\left(\vec{E}'.\vec{u}'\right)\frac{\vec{v}'}{c^2}\right]$ from where we get the invariance of $\vec{E}.\vec{u}=\vec{E}'\vec{u}'$ between the observers as a consequence of 5.3 and the following components of each axis

$$E'x'+u'y'B'z'-u'z'B'y'=\frac{1}{\sqrt{K}}\left[Ex+uyBz-uzBy-\frac{Exuxv}{c^2}-\frac{Eyuyv}{c^2}-\frac{Ezuzv}{c^2}\right] \quad 7.1$$

$$E'y'+u'z'B'x'-u'x'B'z'=\frac{1}{\sqrt{K}}\left[Ey+uzBx-uxBz\right] \quad 7.1.1$$

$$E'z'+u'x'B'y'-u'y'B'x'=\frac{1}{\sqrt{K}}\left[Ez+uxBy-uyBx\right] \quad 7.1.2$$

$$Ex+uyBz-uzBy=\frac{1}{\sqrt{K'}}\left[E'x'+u'y'B'z'-u'z'B'y'+\frac{E'x'u'x'v'}{c^2}+\frac{E'y'u'y'v'}{c^2}+\frac{E'z'u'z'v'}{c^2}\right] \quad 7.2$$

$$Ey+uzBx-uxBz=\frac{1}{\sqrt{K'}}\left[E'y'+u'z'B'x'-u'x'B'z'\right] \quad 7.2.1$$

$$Ez+uxBy-uyBx=\frac{1}{\sqrt{K'}}\left[E'z'+u'x'B'y'-u'y'B'x'\right] \quad 7.2.2$$

To the conjunct 7.1 and 7.2 we have two solutions described in the tables 7 and 8.

Table 7, transformations of the electric fields $\vec{E}$, $\vec{E}'$ and magnetic fields $\vec{B}$ e $\vec{B}'$

$E'x'=\frac{Ex}{\sqrt{K}}\left(1-\frac{vux}{c^2}\right)$	7.3	$Ex=\frac{E'x'}{\sqrt{K'}}\left(1+\frac{v'u'x'}{c^2}\right)$	7.4
$E'y'=\frac{Ey}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)-\frac{vBz}{\sqrt{K}}$	7.3.1	$Ey=\frac{E'y'}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)+\frac{v'B'z'}{\sqrt{K'}}$	7.4.1
$E'z'=\frac{Ez}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)+\frac{vBy}{\sqrt{K}}$	7.3.2	$Ez=\frac{E'z'}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)-\frac{v'B'y'}{\sqrt{K'}}$	7.4.2
$B'x'=Bx$	7.5	$Bx=B'x'$	7.6
$B'y'=By+\frac{v}{c^2}Ez$	7.5.1	$By=B'y'-\frac{v'}{c^2}E'z'$	7.6.1
$B'z'=Bz-\frac{v}{c^2}Ey$	7.5.2	$Bz=B'z'+\frac{v'}{c^2}E'y'$	7.6.2
$E'y'=Ey\sqrt{K}$	7.7	$Ey=E'y'\sqrt{K'}$	7.8
$E'z'=Ez\sqrt{K}$	7.7.1	$Ez=E'z'\sqrt{K'}$	7.8.1
$By=-\frac{ux}{c^2}Ez$	7.9	$B'y'=-\frac{u'x'}{c^2}E'z'$	7.10
$Bz=\frac{ux}{c^2}Ey$	7.9.1	$B'z'=\frac{u'x'}{c^2}E'y'$	7.10.1

Table 8, transformations of the electric fields $\vec{E}$, $\vec{E}'$ and magnetic fields $\vec{B}$ e $\vec{B}'$

$E'x'=\frac{1}{\sqrt{K}}\left[Ex-\left(\vec{E}.\vec{u}\right)\frac{v}{c^2}\right]$	7.11	$Ex=\frac{1}{\sqrt{K'}}\left[E'x'+\left(\vec{E}'.\vec{u}'\right)\frac{v'}{c^2}\right]$	7.12
$E'y'=\frac{1}{\sqrt{K}}\left(Ey-vBz\right)$	7.11.1	$Ey=\frac{1}{\sqrt{K'}}\left(E'y'+v'B'z'\right)$	7.12.1
$E'z'=\frac{1}{\sqrt{K}}\left(Ez+vBy\right)$	7.11.2	$Ez=\frac{1}{\sqrt{K'}}\left(E'z'-v'B'y'\right)$	7.12.2
$B'x'=Bx$	7.13	$Bx=B'x'$	7.14
$B'y'=By$	7.13.1	$By=B'y'$	7.14.1
$B'z'=Bz$	7.13.2	$Bz=B'z'$	7.14.2

Relation between the electric field and magnetic field

If an electric-magnetic field has to the observer O' the naught magnetic component $\vec{B}' = zero$ and the electric component $\vec{E}'$. To the observer O this field is represented with both components, being the magnetic field described by the conjunct 7.5 and hás as components

$$Bx = zero,\ By = -\frac{vEz}{c^2},\ Bz = \frac{vEy}{c^2}, \qquad 7.15$$

that are equivalent to $\vec{B} = \frac{1}{c^2}\vec{v} \times \vec{E}$. 7.16

Formula of Biot-Savart

The observer O' associates to a resting electric charge, uniformly distributed alongside its axis x' the following electric-magnetic properties:

-Linear density of resting electric charge $\rho_o = \frac{dq}{dx'}$

-Naught electric current $I' = zero$

-Naught magnetic field $\vec{B}' = zero \Rightarrow \vec{u}' = zero$

-Radial electrical field of module $E' = \sqrt{E'\,y'^2 + E'\,z'^2} = \frac{\rho_o}{2\pi\varepsilon_o R}$ at any point of radius $R = \sqrt{y'^2 + z'^2}$ with the component $E'\,x' = zero$.

To the observer O it relates to an electric charge uniformly distributed alongside its axis with velocity $ux = v$ to which it associates the following electric-magnetic properties:

-Linear density of the electric charge $\rho = \frac{\rho_o}{\sqrt{1-\frac{v^2}{c^2}}}$ (from 6.7 with u = v)

-Electric current $I = \rho v = \frac{\rho_o v}{\sqrt{1-\frac{v^2}{c^2}}}$

-Radial electrical field of module $E = \frac{E'}{\sqrt{1-\frac{v^2}{c^2}}}$ (according to the conjuncts 7.3 and 7.5 with $\vec{B}' = zero \Rightarrow \vec{u}' = zero$ and $ux = v$)

-Magnetic field of components $Bx = zero$, $By = -\frac{vEz}{c^2}$, $Bz = \frac{vEy}{c^2}$ and module $B = \frac{vE}{c^2} = \frac{v}{c^2}\frac{E'}{\sqrt{1-\frac{v^2}{c^2}}} = \frac{v}{c^2}\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}\frac{\rho_o}{2\pi\varepsilon_o R} = \frac{\mu_o I}{2\pi R}$ where $\mu_o = \frac{1}{\varepsilon_o c^2}$, being in the vectorial form

$$\vec{B} = \frac{\mu_o I}{2\pi R}\vec{u} \qquad 7.17$$

where $\vec{u}$ is a unitary vector perpendicular to the electrical field $\vec{E}$ and tangent to the circumference that passes by the point of radius $R = \sqrt{y^2 + z^2}$ because from the conjunct 7.4 and 7.6 $\vec{E}.\vec{B} = zero$.

§8 Transformations of the differential operators

Table 9, differential operators

$\frac{\partial}{\partial x'}=\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}$	8.1	$\frac{\partial}{\partial x}=\frac{\partial}{\partial x'}-\frac{v'}{c^2}\frac{\partial}{\partial t'}$	8.2
$\frac{\partial}{\partial y'}=\frac{\partial}{\partial y}$	8.1.1	$\frac{\partial}{\partial y}=\frac{\partial}{\partial y'}$	8.2.1
$\frac{\partial}{\partial z'}=\frac{\partial}{\partial z}$	8.1.2	$\frac{\partial}{\partial z}=\frac{\partial}{\partial z'}$	8.2.2
$\frac{\partial}{\partial t'}=\frac{v}{\sqrt{K}}\frac{\partial}{\partial x}+\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vx}{c^2 t}\right)\frac{\partial}{\partial t}$	8.3	$\frac{\partial}{\partial t}=-\frac{v'}{\sqrt{K'}}\frac{\partial}{\partial x'}+\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'x'}{c^2 t'}\right)\frac{\partial}{\partial t'}$	89.4

From the system formed by 8.1, 8.2, 8.3, and 8.4 and with 1.15 and 1.20 we only find the solutions

$$\frac{\partial}{\partial x}+\frac{x/t}{c^2}\frac{\partial}{\partial t}=o \text{ and } \frac{\partial}{\partial x'}+\frac{x'/t'}{c^2}\frac{\partial}{\partial t'}=o\,. \qquad 8.5$$

From where we conclude that only the functions ψ (2.19) and ψ' (2.20) that supply the conditions

$$\frac{\partial\psi}{\partial x}+\frac{x/t}{c^2}\frac{\partial\psi}{\partial t}=o \text{ and } \frac{\partial\psi'}{\partial x'}+\frac{x'/t'}{c^2}\frac{\partial\psi'}{\partial t'}=o\,, \qquad 8.6$$

can represent the propagation with velocity c in the Undulating Relativity indicating that the field propagates with definite velocity and without distortion being applied to 1.13 and 1.18. Because of symmetry we can also write to the other axis

$$\frac{\partial\psi}{\partial y}+\frac{y/t}{c^2}\frac{\partial\psi}{\partial t}=o\,,\ \frac{\partial\psi'}{\partial y'}+\frac{y'/t'}{c^2}\frac{\partial\psi'}{\partial t'}=o \text{ and } \frac{\partial\psi}{\partial z}+\frac{z/t}{c^2}\frac{\partial\psi}{\partial t}=o\,,\ \frac{\partial\psi'}{\partial z'}+\frac{z'/t'}{c^2}\frac{\partial\psi'}{\partial t'}=o\,. \qquad 8.7$$

From the transformations of space and time of the Undulatory Relativity we get to Jacob's theorem

$$J=\frac{\partial(x',y',z',t')}{\partial(x,y,z,t)}=\frac{1-\frac{vux}{c^2}}{\sqrt{K}} \text{ and } J'=\frac{\partial(x,y,z,t)}{\partial(x',y',z',t')}=\frac{1+\frac{v'u'x'}{c^2}}{\sqrt{K'}}\,, \qquad 8.8$$

variables with ux and u'x' as a consequence of the principle of contancy of the light velocity but are equal ais $J=J'$ and will be equal to one $J=J'=1$ when $ux=u'x'=c$.

Invariance of the wave equation

The wave equation to the observer O' is

$$\frac{\partial^2}{\partial x'^2}+\frac{\partial^2}{\partial y'^2}+\frac{\partial^2}{\partial z'^2}-\frac{1}{c^2}\frac{\partial^2}{\partial t'^2}=zero$$

where applying to the formulas of tables 9 and 1.13 we get

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)^2+\frac{\partial^2}{\partial y^2}+\frac{\partial^2}{\partial z^2}-\frac{1}{c^2}\left[\frac{v}{\sqrt{K}}\frac{\partial}{\partial x}+\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)\frac{\partial}{\partial t}\right]^2=zero$$

from where we find

$$K\frac{\partial^2}{\partial x^2}+K\frac{\partial^2}{\partial y^2}+K\frac{\partial^2}{\partial z^2}-\frac{1}{c^2}\frac{\partial^2}{\partial t^2}+\frac{2v}{c^2}\frac{\partial^2}{\partial x\partial t}+\frac{2v^3}{c^4}\frac{\partial^2}{\partial x\partial t}-\frac{4v^2ux}{c^4}\frac{\partial^2}{\partial x\partial t}+\frac{v^2}{c^4}\frac{\partial^2}{\partial t^2}+\frac{v^4}{c^6}\frac{\partial^2}{\partial t^2}-\frac{2v^3ux}{c^6}\frac{\partial^2}{\partial t^2}-$$
$$-\frac{v^2}{c^2}\frac{\partial^2}{\partial x^2}-\frac{2v}{c^2}\frac{\partial^2}{\partial x\partial t}-\frac{2v^3}{c^4}\frac{\partial^2}{\partial x\partial t}+\frac{2v^2ux}{c^4}\frac{\partial^2}{\partial x\partial t}-\frac{2v^2}{c^4}\frac{\partial^2}{\partial t^2}+\frac{2vux}{c^4}\frac{\partial^2}{\partial t^2}+\frac{2v^3ux}{c^6}\frac{\partial^2}{\partial t^2}-\frac{v^2ux^2}{c^6}\frac{\partial^2}{\partial t^2}-\frac{v^4}{c^6}\frac{\partial^2}{\partial t^2}=zero$$

that simplifying supplies
$$K\frac{\partial^2}{\partial x^2}+K\frac{\partial^2}{\partial y^2}+K\frac{\partial^2}{\partial z^2}-\frac{1}{c^2}\frac{\partial^2}{\partial t^2}-\frac{2v^2ux}{c^4}\frac{\partial^2}{\partial x\partial t}-\frac{v^2}{c^2}\frac{\partial^2}{\partial x^2}-\frac{v^2}{c^4}\frac{\partial^2}{\partial t^2}+\frac{2vux}{c^4}\frac{\partial^2}{\partial t^2}-\frac{v^2ux^2}{c^6}\frac{\partial^2}{\partial t^2}=zero$$

where reordering the terms we find
$$K\frac{\partial^2}{\partial x^2}+K\frac{\partial^2}{\partial y^2}+K\frac{\partial^2}{\partial z^2}-\left(1+\frac{v^2}{c^2}-\frac{2vux}{c^2}\right)\frac{1}{c^2}\frac{\partial^2}{\partial t^2}-\frac{v^2}{c^2}\left(\frac{\partial^2}{\partial x^2}+\frac{2ux}{c^2}\frac{\partial^2}{\partial x\partial t}+\frac{ux^2}{c^4}\frac{\partial^2}{\partial t^2}\right)=zero \qquad 8.9$$

but from 8.5 and 1.13 we have $\frac{\partial}{\partial x}+\frac{x/t}{c^2}\frac{\partial}{\partial t}=o\Rightarrow\left(\frac{\partial}{\partial x}+\frac{u\,x}{c^2}\frac{\partial}{\partial t}\right)^2=\frac{\partial^2}{\partial x^2}+\frac{2ux}{c^2}\frac{\partial^2}{\partial x\partial t}+\frac{ux^2}{c^4}\frac{\partial^2}{\partial t^2}=zero$

that applied in 8.9 supplies the wave equation to the observer O $\frac{\partial^2}{\partial x^2}+\frac{\partial^2}{\partial y^2}+\frac{\partial^2}{\partial z^2}-\frac{1}{c^2}\frac{\partial^2}{\partial t^2}=zero$. 8.10

To return to the referential of the observer O' we will apply 8.10 to the formulas of tables 9 and 1.18, getting
$$\left(\frac{\partial}{\partial x'}-\frac{v'}{c^2}\frac{\partial}{\partial t'}\right)^2+\frac{\partial^2}{\partial y'^2}+\frac{\partial^2}{\partial z'^2}-\frac{1}{c^2}\left[-\frac{v'}{\sqrt{K'}}\frac{\partial}{\partial x'}+\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)\frac{\partial}{\partial t'}\right]^2=zero$$

from where we find
$$K'\frac{\partial^2}{\partial x'^2}+K'\frac{\partial^2}{\partial y'^2}+K'\frac{\partial^2}{\partial z'^2}-\frac{1}{c^2}\frac{\partial^2}{\partial t'^2}-\frac{2v'}{c^2}\frac{\partial^2}{\partial x'\partial t'}-\frac{2v'^3}{c^4}\frac{\partial^2}{\partial x'\partial t'}-\frac{4v'^2u'x'}{c^4}\frac{\partial^2}{\partial x'\partial t'}+\frac{v'^2}{c^4}\frac{\partial^2}{\partial t'^2}+\frac{v'^4}{c^6}\frac{\partial^2}{\partial t'^2}+$$
$$+\frac{2v'^3u'x'}{c^6}\frac{\partial^2}{\partial t'^2}-\frac{v'^2}{c^2}\frac{\partial^2}{\partial x'^2}+\frac{2v'}{c^2}\frac{\partial^2}{\partial x'\partial t'}+\frac{2v'^3}{c^4}\frac{\partial^2}{\partial x'\partial t'}+\frac{2v'^2u'x'}{c^4}\frac{\partial^2}{\partial x'\partial t'}-\frac{2v'^2}{c^4}\frac{\partial^2}{\partial t'^2}-\frac{2v'u'x'}{c^4}\frac{\partial^2}{\partial t'^2}-$$
$$-\frac{v'^3u'x'}{c^6}\frac{\partial^2}{\partial t'^2}-\frac{v'^2u'x'^2}{c^6}\frac{\partial^2}{\partial t'^2}-\frac{v'^4}{c^6}\frac{\partial^2}{\partial t'^2}=zero$$

that simplifying supplies
$$K'\frac{\partial^2}{\partial x'^2}+K'\frac{\partial^2}{\partial y'^2}+K'\frac{\partial^2}{\partial z'^2}-\frac{1}{c^2}\frac{\partial^2}{\partial t'^2}-\frac{2v'^2u'x'}{c^4}\frac{\partial^2}{\partial x'\partial t'}-\frac{v'^2}{c^2}\frac{\partial^2}{\partial x'^2}-\frac{v'^2}{c^4}\frac{\partial^2}{\partial t'^2}-\frac{2v'u'x'}{c^4}\frac{\partial^2}{\partial t'^2}-\frac{v'^2u'x'^2}{c^6}\frac{\partial^2}{\partial t'^2}=zero$$

where reordering the terms we find
$$K'\frac{\partial^2}{\partial x'^2}+K'\frac{\partial^2}{\partial y'^2}+K'\frac{\partial^2}{\partial z'^2}-\left(1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}\right)\frac{1}{c^2}\frac{\partial^2}{\partial t'^2}-\frac{v'^2}{c^2}\left(\frac{\partial^2}{\partial x'^2}+\frac{2u'x'}{c^2}\frac{\partial^2}{\partial x'\partial t'}+\frac{u'x'^2}{c^4}\frac{\partial^2}{\partial t'^2}\right)=zero$$

but from 8.5 and 1.18 we have
$$\frac{\partial}{\partial x'}+\frac{x'/t'}{c^2}\frac{\partial}{\partial t'}=o\Rightarrow\left(\frac{\partial}{\partial x'}+\frac{u'x'}{c^2}\frac{\partial}{\partial t'}\right)^2=\frac{\partial^2}{\partial x'^2}+\frac{2u'x'}{c^2}\frac{\partial^2}{\partial x'\partial t'}+\frac{u'x'^2}{c^4}\frac{\partial^2}{\partial t'^2}=zero$$

that replaced in the reordered equation supplies the wave equation to the observer O'.

Invariance of the Continuity equation

The continuity equation in the differential form to the observer O' is

$$\frac{\partial\rho'}{\partial t'}+\vec{\nabla}.\vec{J}'=zero\Rightarrow\frac{\partial\rho'}{\partial t'}+\frac{\partial Jx'}{\partial x'}+\frac{\partial Jy'}{\partial y'}+\frac{\partial Jz'}{\partial z'}=zero \qquad 8.11$$

where replacing the formulas of tables 6, 9, and 1.13 we get
$$\left(\frac{v}{\sqrt{K}}\frac{\partial}{\partial x}+\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)\frac{\partial}{\partial t}\right)\rho\sqrt{K}+\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)(Jx-\rho v)+\frac{\partial Jy}{\partial y}+\frac{\partial Jz}{\partial z}=zero$$

making the operations we find

$$\frac{v\partial\rho}{\partial x}+\frac{\partial\rho}{\partial t}+\frac{v^2}{c^2}\frac{\partial\rho}{\partial t}-\frac{vux}{c^2}\frac{\partial\rho}{\partial t}+\frac{\partial Jx}{\partial x}+\frac{v}{c^2}\frac{\partial Jx}{\partial t}-\frac{v\partial\rho}{\partial x}-\frac{v^2}{c^2}\frac{\partial\rho}{\partial t}+\frac{\partial Jy}{\partial y}+\frac{\partial Jz}{\partial z}=zero$$

that simplifying supplies

$$\frac{\partial\rho}{\partial t}-\frac{vux}{c^2}\frac{\partial\rho}{\partial t}+\frac{\partial Jx}{\partial x}+\frac{v}{c^2}\frac{\partial Jx}{\partial t}+\frac{\partial Jy}{\partial y}+\frac{\partial Jz}{\partial z}=zero$$

where applying $Jx=\rho ux$ with ux constant we get

$$\frac{\partial\rho}{\partial t}-\frac{vux}{c^2}\frac{\partial\rho}{\partial t}+\frac{\partial Jx}{\partial x}+\frac{v}{c^2}\frac{\partial(\rho ux)}{\partial t}+\frac{\partial Jy}{\partial y}+\frac{\partial Jz}{\partial z}=zero\Rightarrow\frac{\partial\rho}{\partial t}+\frac{\partial Jx}{\partial x}+\frac{\partial Jy}{\partial y}+\frac{\partial Jz}{\partial z}=zero \qquad 8.12$$

that is the continuity equation in the differential form to the observer O.

To get again the continuity equation in the differential form to the observer O' we will replace the formulas of tables 6, 9, and 1.18 in 8.12 getting

$$\left(-\frac{v'}{\sqrt{K'}}\frac{\partial}{\partial x'}+\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)\frac{\partial}{\partial t'}\right)\rho'\sqrt{K'}+\left(\frac{\partial}{\partial x'}-\frac{v'}{c^2}\frac{\partial}{\partial t'}\right)(J'x'+\rho'v')+\frac{\partial J'y'}{\partial y'}+\frac{\partial J'z'}{\partial z'}=zero$$

making the operations we find

$$-\frac{v'\partial\rho'}{\partial x'}+\frac{\partial\rho'}{\partial t'}+\frac{v'^2}{c^2}\frac{\partial\rho'}{\partial t'}+\frac{v'u'x'}{c^2}\frac{\partial\rho'}{\partial t'}+\frac{\partial J'x'}{\partial x'}-\frac{v'}{c^2}\frac{\partial J'x'}{\partial t'}+\frac{v'\partial\rho'}{\partial x'}-\frac{v'^2}{c^2}\frac{\partial\rho'}{\partial t'}+\frac{\partial J'y'}{\partial y'}+\frac{\partial J'z'}{\partial z'}=zero$$

that simplifying supplies

$$\frac{\partial\rho'}{\partial t'}+\frac{v'u'x'}{c^2}\frac{\partial\rho'}{\partial t'}+\frac{\partial J'x'}{\partial x'}-\frac{v'}{c^2}\frac{\partial J'x'}{\partial t'}+\frac{\partial J'y'}{\partial y'}+\frac{\partial J'z'}{\partial z'}=zero$$

where applying $J'x'=\rho'u'x'$ with u'x' constant we get

$$\frac{\partial\rho'}{\partial t'}+\frac{v'u'x'}{c^2}\frac{\partial\rho'}{\partial t'}+\frac{\partial J'x'}{\partial x'}-\frac{v'}{c^2}\frac{\partial(\rho'u'x')}{\partial t'}+\frac{\partial J'y'}{\partial y'}+\frac{\partial J'z'}{\partial z'}=zero\Rightarrow\frac{\partial\rho'}{\partial t'}+\frac{\partial J'x'}{\partial x'}+\frac{\partial J'y'}{\partial y'}+\frac{\partial J'z'}{\partial z'}=zero$$

that is the continuity equation in the differential form to the observer O'.

Invariance of Maxwell's equations

That in the differential form are written this way

With electrical charge

To the observer O		To the observer O'	
$\frac{\partial Ex}{\partial x}+\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}=\frac{\rho}{\varepsilon_o}$	8.13	$\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z'}{\partial z'}=\frac{\rho'}{\varepsilon_o}$	8.14
$\frac{\partial Bx}{\partial x}+\frac{\partial By}{\partial y}+\frac{\partial Bz}{\partial z}=0$	8.15	$\frac{\partial B'x'}{\partial x'}+\frac{\partial B'y'}{\partial y'}+\frac{\partial B'z'}{\partial z'}=0$	8.16
$\frac{\partial Ey}{\partial x}-\frac{\partial Ex}{\partial y}=-\frac{\partial Bz}{\partial t}$	8.17	$\frac{\partial E'y'}{\partial x'}-\frac{\partial E'x'}{\partial y'}=-\frac{\partial B'z'}{\partial t'}$	8.18
$\frac{\partial Ez}{\partial y}-\frac{\partial Ey}{\partial z}=-\frac{\partial Bx}{\partial t}$	8.19	$\frac{\partial E'z'}{\partial y'}-\frac{\partial E'y'}{\partial z'}=-\frac{\partial B'x'}{\partial t'}$	8.20
$\frac{\partial Ex}{\partial z}-\frac{\partial Ez}{\partial x}=-\frac{\partial By}{\partial t}$	8.21	$\frac{\partial E'x'}{\partial z'}-\frac{\partial E'z'}{\partial x'}=-\frac{\partial B'y'}{\partial t'}$	8.22
$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\mu_o Jz+\varepsilon_o\mu_o\frac{\partial Ez}{\partial t}$	8.23	$\frac{\partial B'y'}{\partial x'}-\frac{\partial B'x'}{\partial y'}=\mu_o J'z'+\varepsilon_o\mu_o\frac{\partial E'z'}{\partial t'}$	8.24

$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_o Jx+\varepsilon_o\mu_o\frac{\partial Ex}{\partial t}$	8.25	$\frac{\partial B'z'}{\partial y'}-\frac{\partial B'y'}{\partial z'}=\mu_o J'x'+\varepsilon_o\mu_o\frac{\partial E'x'}{\partial t'}$	8.26
$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\mu_o Jy+\varepsilon_o\mu_o\frac{\partial Ey}{\partial t}$	8.27	$\frac{\partial B'x'}{\partial z'}-\frac{\partial B'z'}{\partial x'}=\mu_o J'y'+\varepsilon_o\mu_o\frac{\partial E'y'}{\partial t'}$	8.28

Without electrical charge $\rho=\rho'=zero$ and $\vec{J}=\vec{J}'=zero$

To the observer O		To the observer O'	
$\frac{\partial Ex}{\partial x}+\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}=0$	8.29	$\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z'}{\partial z'}=0$	8.30
$\frac{\partial Bx}{\partial x}+\frac{\partial By}{\partial y}+\frac{\partial Bz}{\partial z}=0$	8.31	$\frac{\partial B'x'}{\partial x'}+\frac{\partial B'y'}{\partial y'}+\frac{\partial B'z'}{\partial z'}=0$	8.32
$\frac{\partial Ey}{\partial x}-\frac{\partial Ex}{\partial y}=-\frac{\partial Bz}{\partial t}$	8.33	$\frac{\partial E'y'}{\partial x'}-\frac{\partial E'x'}{\partial y'}=-\frac{\partial B'z'}{\partial t'}$	8.34
$\frac{\partial Ez}{\partial y}-\frac{\partial Ey}{\partial z}=-\frac{\partial Bx}{\partial t}$	8.35	$\frac{\partial E'z'}{\partial y'}-\frac{\partial E'y'}{\partial z'}=-\frac{\partial B'x'}{\partial t'}$	8.36
$\frac{\partial Ex}{\partial z}-\frac{\partial Ez}{\partial x}=-\frac{\partial By}{\partial t}$	8.37	$\frac{\partial E'x'}{\partial z'}-\frac{\partial E'z'}{\partial x'}=-\frac{\partial B'y'}{\partial t'}$	8.38
$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\varepsilon_o\mu_o\frac{\partial Ez}{\partial t}$	8.39	$\frac{\partial B'y'}{\partial x'}-\frac{\partial B'x'}{\partial y'}=\varepsilon_o\mu_o\frac{\partial E'z'}{\partial t'}$	8.40
$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\varepsilon_o\mu_o\frac{\partial Ex}{\partial t}$	8.41	$\frac{\partial B'z'}{\partial y'}-\frac{\partial B'y'}{\partial z'}=\varepsilon_o\mu_o\frac{\partial E'x'}{\partial t'}$	8.42
$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\varepsilon_o\mu_o\frac{\partial Ey}{\partial t}$	8.43	$\frac{\partial B'x'}{\partial z'}-\frac{\partial B'z'}{\partial x'}=\varepsilon_o\mu_o\frac{\partial E'y'}{\partial t'}$	8.44
$\varepsilon_o\mu_o=\frac{1}{c^2}$	8.45		

We demonstrate the invariance of the Law of Gauss in the differential form that for the observer O' is

$$\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z'}{\partial z'}=\frac{\rho'}{\varepsilon_o} \qquad 8.14$$

where replacing the formulas from the tables 6, 7, 9, and 1.18, and considering u'x' constant, we get

$$\left[\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right]\frac{Ex}{\sqrt{K}}\left(1-\frac{vux}{c^2}\right)+\frac{\partial}{\partial y}\left[\frac{Ey}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)-\frac{vBz}{\sqrt{K}}\right]+$$

$$+\frac{\partial}{\partial z}\left[\frac{Ez}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)+\frac{vBy}{\sqrt{K}}\right]=\frac{\rho\sqrt{K}}{\varepsilon_o}$$

making the products, summing and subtracting the term $\frac{v^2}{c^2}\frac{\partial Ex}{\partial x}$, we find

$$\frac{\partial Ex}{\partial x}+\frac{v}{c^2}\frac{\partial Ex}{\partial t}-\frac{vux}{c^2}\frac{\partial Ex}{\partial x}-\frac{v^2ux}{c^4}\frac{\partial Ex}{\partial t}+\frac{\partial Ey}{\partial y}+\frac{v^2}{c^2}\frac{\partial Ey}{\partial y}-\frac{vux}{c^2}\frac{\partial Ey}{\partial y}-\frac{v\partial Bz}{\partial y}+$$

$$+\frac{\partial Ez}{\partial z}+\frac{v^2}{c^2}\frac{\partial Ez}{\partial z}-\frac{vux}{c^2}\frac{\partial Ez}{\partial z}+\frac{v\partial By}{\partial z}+\frac{v^2}{c^2}\frac{\partial Ex}{\partial x}-\frac{v^2}{c^2}\frac{\partial Ex}{\partial x}=\frac{\rho K}{\varepsilon_o}$$

that reordering results

$$-\frac{v^2}{c^2}\left(\frac{\partial Ex}{\partial x}+\frac{ux}{c^2}\frac{\partial Ex}{\partial t}\right)-v\left(\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}-\frac{1}{c^2}\frac{\partial Ex}{\partial t}\right)+\left(\frac{\partial Ex}{\partial x}+\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}\right)\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)=\frac{\rho K}{\varepsilon_o}$$

where the first parentheses is 8.5 and because of this equal to zero , the second blank is equal to $-v(\mu_o Jx) = -v\mu_o \rho ux = -\dfrac{v\rho ux}{\varepsilon_o c^2}$ gotten from 8.25 and 8.45 resulting in

$$\left(\frac{\partial Ex}{\partial x}+\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}\right)\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)=\frac{\rho}{\varepsilon_o}\left(1+\frac{v^2}{c^2}-\frac{vux}{c^2}\right)-\frac{\rho}{\varepsilon_o}\frac{vux}{c^2}+\frac{\rho}{\varepsilon_o}\frac{vux}{c^2}$$

from where we get $\dfrac{\partial Ex}{\partial x}+\dfrac{\partial Ey}{\partial y}+\dfrac{\partial Ez}{\partial z}=\dfrac{\rho}{\varepsilon_o}$ 8.13

that is the Law of Gauss in the differential form to the observer O.

To make the inverse we will replace in 8.13 the formulas of the tables 6, 7, 9, and 1.13, and considering ux constant, we get

$$\left[\frac{\partial}{\partial x'}-\frac{v'}{c^2}\frac{\partial}{\partial t'}\right]\frac{E'x'}{\sqrt{K'}}\left(1+\frac{v'u'x'}{c^2}\right)+\frac{\partial}{\partial y'}\left[\frac{E'y'}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)+\frac{v'B'z'}{\sqrt{K'}}\right]+$$

$$+\frac{\partial}{\partial z'}\left[\frac{E'z'}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)-\frac{v'B'y'}{\sqrt{K'}}\right]=\frac{\rho'\sqrt{K'}}{\varepsilon_o}$$

making the products, adding and subtracting the term $\dfrac{v'^2}{c^2}\dfrac{\partial E'x'}{\partial x'}$, we get

$$\frac{\partial E'x'}{\partial x'}-\frac{v'}{c^2}\frac{\partial E'x'}{\partial t'}+\frac{v'u'x'}{c^2}\frac{\partial E'x'}{\partial x'}-\frac{v'^2u'x'}{c^4}\frac{\partial E'x'}{\partial t'}+\frac{\partial E'y'}{\partial y'}+\frac{v'^2}{c^2}\frac{\partial E'y'}{\partial y'}+\frac{v'u'x'}{c^2}\frac{\partial E'y'}{\partial y'}+$$

$$+\frac{v'\partial B'z'}{\partial y'}+\frac{\partial E'z'}{\partial z'}+\frac{v'^2}{c^2}\frac{\partial E'z'}{\partial z'}+\frac{v'u'x'}{c^2}\frac{\partial E'z'}{\partial z'}-\frac{v'\partial B'y'}{\partial z'}+\frac{v'^2}{c^2}\frac{\partial E'x'}{\partial x'}-\frac{v'^2}{c^2}\frac{\partial E'x'}{\partial x'}=\frac{\rho'K'}{\varepsilon_o}$$

that reordering results in

$$-\frac{v'^2}{c^2}\left(\frac{\partial E'x'}{\partial x'}+\frac{u'x'}{c^2}\frac{\partial E'x'}{\partial t'}\right)+v'\left(\frac{\partial B'z'}{\partial y'}-\frac{\partial B'y'}{\partial z'}-\frac{1}{c^2}\frac{\partial E'x'}{\partial t'}\right)+$$

$$+\left(\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z''}{\partial z'}\right)\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)=\frac{\rho'K'}{\varepsilon_o}$$

where the first blank is 8.5 and because of this equals to zero, the second blank is equal to $v'(\mu_o J'x') = v'\mu_o\rho'u'x' = \dfrac{v'\rho'u'x'}{\varepsilon_o c^2}$ gotten from 8.26 and 8.45 resulting in

$$\left(\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z'}{\partial z'}\right)\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)=\frac{\rho'}{\varepsilon_o}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'}{c^2}\right)+\frac{\rho}{\varepsilon_o}\frac{v'u'x'}{c^2}-\frac{\rho'}{\varepsilon_o}\frac{v'u'x'}{c^2}$$

from where we get $\dfrac{\partial E''x'}{\partial x'}+\dfrac{\partial E'y'}{\partial y'}+\dfrac{\partial E'z'}{\partial z'}=\dfrac{\rho'}{\varepsilon_o}$ that is the Law of Gauss in the differential form to the O' observer.

Proceeding this way we can prove the invariance of form for all the other equations of Maxwell.

§9 Explaining the Sagnac Effect with the Undulating Relativity

We must transform the straight movement of the two observers O and O' used in the deduction of the Undulating Relativity in a plain circular movement with a constant radius. Let's imagine that the observer O sees the observer O' turning around with a tangential speed v in a clockwise way (C) equals to the positive course of the axis x of UR and that the observer O' sees the observer O turning around with a tangecial speed v' in a unclockwise way (U) equals to the negative course of the axis x of the UR.

In the moment t = t' = zero, the observer O emits two rays of light from the common origin to both observers, one in a unclockwise way of arc ct_U and another in a clockwise way of arc ct_C, therefore $ct_U = ct_C$ and $t_U = t_C$, because c is the speed of the constant light, and t_U and t_C the time.

In the moment t = t' = zero the observer O' also emits two rays of light from the common origin to both observers, one in a unclockwise way (useless) of arc ct'_U and another one in a clockwise way of arc ct'_C, thus $ct'_U = ct'_C$ and $t'_U = t'_C$ because c is the speed of the constant light, and t'_U and t'_C the time.

Rewriting the equations 1.15 and 1.20 of the Undulating Relativity (UR):

$$\frac{|v|}{|v'|} = \frac{t'}{t} = \sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}} . \qquad 1.15$$

$$\frac{|v'|}{|v|} = \frac{t}{t'} = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'u'x'}{c^2}} . \qquad 1.20$$

Making ux = u'x' = c (ray of light projected alongside the positive axis x) and splitting the equations we have:

$$t' = t\left(1 - \frac{v}{c}\right) \qquad 9.1 \qquad\qquad t = t'\left(1 + \frac{v'}{c}\right) \qquad 9.2$$

$$v' = \frac{v}{\left(1 - \frac{v}{c}\right)} \qquad 9.3 \qquad\qquad v = \frac{v'}{\left(1 + \frac{v'}{c}\right)} \qquad 9.4$$

When the origin of the observer O' detects the unclockwise ray of the observer O, will be at the distance $vt_C = v't'_U$ of the observer O and simultaneously will detect its clockwise ray of light at the same point of the observer O, in a symmetric position to the diameter that goes through the observer O because $ct_U = ct_C \Rightarrow t_U = t_C$ and $ct'_U = ct'_C \Rightarrow t'_U = t'_C$, following the four equations above we have:

$$ct_U + vt_C = 2\pi R \Rightarrow t_C = \frac{2\pi R}{c + v} \qquad 9.5$$

$$ct'_C + 2v't'_U = 2\pi R \Rightarrow t'_C = \frac{2\pi R}{c + 2v'} \qquad 9.6$$

When the origin of the observer O' detects the clockwise ray of the observer O, simultaneously will detect its own clockwise ray and will be at the distance $vt_{2C} = v't'_{2U}$ of the observer O, then following the equations 1,2,3 and 4 above we have:

$$ct_{2C} = 2\pi R + vt_{2C} \Rightarrow t_{2C} = \frac{2\pi R}{c - v} \qquad 9.7$$

$$ct'_{2C} = 2\pi R \Rightarrow t'_{2C} = \frac{2\pi R}{c} \qquad 9.8$$

The time difference to the observer O is:

$$\Delta t = t_{2C} - t_C = \frac{2\pi R}{c - v} - \frac{2\pi R}{c + v} = \frac{4\pi R v}{c^2 - v^2} \qquad 9.9$$

The time difference to the observer O' is:

$$\Delta t' = t'_{2C} - t'_C = \frac{2\pi R}{c} - \frac{2\pi R}{c + 2v'} = \frac{4\pi R v'}{(c + 2v')c} \qquad 9.10$$

Replacing the equations 5 to 10 in 1 to 4 we prove that they confirm the transformations of the Undulating Relativity.

§10 Explaining the experience of Ives-Stilwell with the Undulating Relativity

We should rewrite the equations (2.21) to the wave length in the Undulating Relativity:

$$\lambda'=\frac{\lambda}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}} \text{ and } \lambda=\frac{\lambda'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}}, \qquad 2.21$$

Making ux = u'x' = c (Ray of light projected alongside the positive axis x), we have the equations:

$$\lambda'=\frac{\lambda}{\left(1-\frac{v}{c}\right)} \text{ and } \lambda=\frac{\lambda'}{\left(1+\frac{v'}{c}\right)}, \qquad 10.1$$

If the observer O, who sees the observer O' going away with the velocity v in the positive way of the axis x, emits waves, provenient of a resting source in its origin with velocity c and wave length λ_F in the positive way of the axis x, then according to the equation 10.1 the observer O' will measure the waves with velocity c and the wave length λ'_D according to the formulas:

$$\lambda'_D=\frac{\lambda_F}{\left(1-\frac{v}{c}\right)} \text{ and } \lambda_F=\frac{\lambda'_D}{\left(1+\frac{v'}{c}\right)}, \qquad 10.2$$

If the observer O', who sees the obsesrver O going away with velocity v' in the negative way of the axis x, emits waves, provenient of a resting source in its origin with velocity c and the wave length λ'_F in the positive way of the axis x, then according to the equation 10.1 the observer O will measure waves with velocity c and wave lenght λ_A according to the formulas:

$$\lambda'_F=\frac{\lambda_A}{\left(1-\frac{v}{c}\right)} \text{ and } \lambda_A=\frac{\lambda'_F}{\left(1+\frac{v'}{c}\right)}, \qquad 10.3$$

The resting sources in the origin of the observers O and O' are identical thus $\lambda_F=\lambda'_F$.

We calculate the average wave length $\overline{\lambda}$ of the measured waves (λ_A,λ'_D) using the equations 10.2 and 10.3, the left side in each equation:

$$\overline{\lambda}=\frac{\lambda'_D+\lambda_A}{2}=\frac{1}{2}\left[\frac{\lambda_F}{\left(1-\frac{v}{c}\right)}+\lambda'_F\left(1-\frac{v}{c}\right)\right] \Rightarrow \overline{\lambda}=\frac{\lambda'_D+\lambda_A}{2}=\frac{\lambda_F}{2\left(1-\frac{v}{c}\right)}\left[1+\left(1-\frac{v}{c}\right)^2\right]$$

We calculate the diffrence between the average wave length $\overline{\lambda}$ and the emited wave length by the sources $\Delta\overline{\lambda}=\overline{\lambda}-\lambda_F$:

$$\Delta\overline{\lambda}=\overline{\lambda}-\lambda_F=\frac{\lambda_F}{2\left(1-\frac{v}{c}\right)}\left[1+\left(1-\frac{v}{c}\right)^2\right]-\lambda_F$$

$$\Delta\overline{\lambda}=\frac{\lambda_F}{2\left(1-\frac{v}{c}\right)}\left[1+\left(1-\frac{v}{c}\right)^2\right]-\lambda_F\frac{2\left(1-\frac{v}{c}\right)}{2\left(1-\frac{v}{c}\right)}$$

$$\Delta\overline{\lambda}=\frac{\lambda_F}{2\left(1-\frac{v}{c}\right)}\left[1+\left(1-\frac{v}{c}\right)^2-2\left(1-\frac{v}{c}\right)\right]$$

$$\Delta\bar{\lambda}=\frac{\lambda_F}{2\left(1-\frac{v}{c}\right)}\left[1+1-2\frac{v}{c}+\frac{v^2}{c^2}-2+2\frac{v}{c}\right]$$

$$\Delta\bar{\lambda}=\frac{1}{\left(1-\frac{v}{c}\right)}\frac{\lambda_F}{2}\frac{v^2}{c^2} \qquad 10.4$$

Reference
http://www.wbabin.net/physics/faraj7.htm

§10 Ives-Stilwell (continuation)

The Doppler's effect transversal to the Undulating Relativity was obtained in the §2 as follows:

If the observer O', that sees the observer O, moves with the speed –v' in a negative way to the axis x', emits waves with the frequency y' and the speed c then the observer O according to 2.22 and $u'x'=-v'$ will measure waves of frequency y and speed c in a perpendicular plane to the movement of O' given by

$$y=y'\sqrt{1-\frac{v'^2}{c^2}} \qquad 2.25$$

For $u'x'=-v'$ we will have $ux=zero$ and $\sqrt{1-\frac{v'^2}{c^2}}\sqrt{1+\frac{v^2}{c^2}}=1$ with this we can write the relation between the transversal frequency $y=y_t$ and the source frequency $y'=y'_F$ like this

$$y_t=\frac{y'_F}{\sqrt{1+\frac{v^2}{c^2}}} \qquad 10.5$$

With $c=y_t\lambda_t=y'_F\lambda'_F$ we have the relation between the length of the transversal wave λ_t and the length of the source wave λ'_F

$$\lambda_t=\lambda'_F\sqrt{1+\frac{v^2}{c^2}} \qquad 10.6$$

The variation of the length of the transversal wave in the relation to the length of the source wave is:

$$\Delta\lambda_t=\lambda_t-\lambda'_F=\lambda'_F\sqrt{1+\frac{v^2}{c^2}}-\lambda'_F=\lambda'_F\left(\sqrt{1+\frac{v^2}{c^2}}-1\right)\cong\lambda'_F\left(1+\frac{v^2}{2c^2}-1\right)\cong\frac{\lambda'_F}{2}\frac{v^2}{c^2} \qquad 10.7$$

that is the same value gotten in the Theory of Special Relativity.

Applying 10.7 in 10.4 we have

$$\Delta\bar{\lambda}=\frac{\Delta\lambda_t}{\left(1-\frac{v}{c}\right)} \qquad 10.8$$

With the equations 10.2 and 10.3 we can get the relations 10.9, 10.10, and 10.11 described as follows

$$\lambda_A=\lambda'_D\left(1-\frac{v}{c}\right)^2 \qquad 10.9$$

And from this we have the formula of speed $\frac{v}{c}=1-\sqrt{\frac{\lambda_A}{\lambda'_D}}$ 10.10

$$\lambda_F=\lambda'_F=\sqrt{\lambda_A\lambda'_D} \qquad 10.11$$

Applying 10.10 and 10.11 in 10.6 we have

$$\lambda_t=\sqrt{\lambda_A\lambda'_D}\sqrt{1+\left(1-\sqrt{\frac{\lambda_A}{\lambda'_D}}\right)^2} \qquad 10.12$$

From 10.8 and 10.12 we conclude that $\lambda_A \le \lambda_F \le \lambda_t \le \overline{\lambda} \le \lambda'_D$. 10.13

So that we the values of λ_A and λ'_D obtained from the Ives-Stiwell experience we can evaluate λ_t, λ_F, $\frac{v}{c}$ and conclude whether there is or not the space deformation predicted in the Theory of Special Relativity.

§11 Transformation of the power of a luminous ray between two referencials in the Special Theory of Relativity

The relationship within the power developed by the forces between two referencials is written in the Special Theory of the Relativity in the following way:

$$\vec{F}'.\vec{u}' = \frac{\vec{F}.\vec{u} - vFx}{\left(1 - \frac{vux}{c^2}\right)} \qquad 11.1$$

The definition of the component of the force along the axis x is:

$$Fx = \frac{dpx}{dt} = \frac{d(mux)}{dt} = \frac{dm}{dt}ux + m\frac{dux}{dt} \qquad 11.2$$

For a luminous ray, the principle of light speed constancy guarantees that the component ux of the light speed is also constant along its axis, thus

$\frac{x}{t} = \frac{dx}{dt} = ux =$ constant, demonstrating that in two $\frac{dux}{dt} = zero$ and $Fx = \frac{dm}{dt}ux$ 11.3

The formula of energy is $E = mc^2$ from where we have $\frac{dm}{dt} = \frac{1}{c^2}\frac{dE}{dt}$ 11.4

From the definition of energy we have $\frac{dE}{dt} = \vec{F}.\vec{u}$ that applying in 4 and 3 we have $Fx = \vec{F}.\vec{u}\frac{ux}{c^2}$ 11.5

Applying 5 in 1 we heve:

$$\vec{F}'.\vec{u}' = \frac{\vec{F}.\vec{u} - \left(\vec{F}.\vec{u}\right)\frac{vux}{c^2}}{\left(1 - \frac{vux}{c^2}\right)}$$

From where we find that $\vec{F}'.\vec{u}' = \vec{F}.\vec{u}$ or $\frac{dE'}{dt'} = \frac{dE}{dt}$ 11.6

A result equal to 5.3 of the Undulating Relativity that can be experimentally proven, considering the ‘Sun’ as the source.

§12 Linearity

The Theory of Undulating Relativity has as its fundamental axiom the necessity that inertial referentials be named exclusively as those ones in which a ray of light emitted in any direction from its origin spreads in a straight line, what is mathematically described by the formulae (1.13, 1.18, 8.6 e 8.7) of the Undulating Relativity:

$$\frac{x}{t} = \frac{dx}{dt} = ux, \frac{y}{t} = \frac{dy}{dt} = uy, \frac{z}{t} = \frac{dz}{dt} = uz \qquad 1.13$$

$$\frac{x'}{t'} = \frac{dx'}{dt'} = u'x', \frac{y'}{t'} = \frac{dy'}{dt'} = u'y', \frac{z'}{t'} = \frac{dz'}{dt'} = u'z' \qquad 1.18$$

Woldemar Voigt wrote in 1.887 the linear transformation between the referentials os the observers O e O’ in the following way:

$$x = Ax' + Bt' \qquad 12.1$$

$$t = Ex' + Ft' \qquad 12.2$$

With the respective inverted equations:

$$x' = \frac{F}{AF - BE}x + \frac{-B}{AF - BE}t \qquad 12.3$$

$$t' = \frac{-E}{AF - BE}x + \frac{A}{AF - BE}t \qquad 12.4$$

Where A, B, E and F are constants and because of the symmetry we don't consider the terms with y, z and y', z'.
We know that x and x' are projections of the two rays of lights ct and ct' that spread with Constant speed c (due to the constancy principle of the Ray of light), emited in any direction from the origin of the respective inertials referential at the moment in which the origins are coincident and at the moment where:

t = t' = zero 12.5

because of this in the equation 12.2 at the moment where t' = zero we must have E = zero so that we also have t = zero, we can't assume that when t' = zero, x' also be equal to zero, because if the spreading happens in the plane y'z' we will have x' = zero plus $t' \neq zero$.

We should rewrite the corrected equations (E = zero):

$$x = Ax' + Bt' \qquad 12.6$$

$$t = Ft' \qquad 12.7$$

With the respective corrected inverted equations:

$$x' = \frac{x}{A} - \frac{Bt}{AF} \qquad 12.8$$

$$t' = \frac{t}{F} \qquad 12.9$$

If the spreading happens in the plane y' z' we have x' = zero and dividing 12.6 by 12.7 we have:

$$\frac{x}{t} = \frac{B}{F} = v \qquad 12.10$$

where v is the module of the speed in which the observer O sees the referential of the observer O' moving alongside the x axis in the positive way because the sign of the equation is positive.

If the spreading happens in the plane y z we have x = zero and dividing 12.8 by 12.9 we have:

$$\frac{x'}{t'} = -\frac{B}{A} = -v' \text{ or } \frac{B}{A} = v' \qquad 12.11$$

where v' is the module of the speed in which the observer O' sees the referential of the observer O moving alongside the x' axis in the negative way because the signal of the equation is negative.

The equation 1.6 describes the constancy principle of the speed of light that must be assumed by the equations 12.6 to 12.9:

$$x^2 - c^2t^2 = x'^2 - c^2t'^2 \qquad 1.6$$

Applying 12.6 and 12.7 in 1.6 we have:

$$(Ax' + Bt')^2 - c^2F^2t'^2 = x'^2 - c^2t'^2$$

From where we have:

$$\left(A^2 x'^2\right) - c^2 t'^2 \left[F^2 - \frac{B^2}{c^2} - \frac{2ABx'}{c^2 t'}\right] = x'^2 - c^2 t'^2$$

where making A^2 = 1 in the brackets in arc and $\left[F^2 - \frac{B^2}{c^2} - \frac{2ABx'}{c^2 t'}\right] = 1$ in the straight brackets we have the equality between both sides of the equal signal of the equation.

Appllying A = 1 in $\left[F^2 - \frac{B^2}{c^2} - \frac{2ABx'}{c^2 t'}\right] = 1$ we have $F^2 = 1 + \frac{B^2}{c^2} + \frac{2Bx'}{c^2 t'}$ 12.12

Appllying A = 1 in 12.11 we have $\frac{B}{A} = \frac{B}{1} = B = v'$ 12.11

That applied in 12.12 suplies:

$$F = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2 t'}} = F(x', t') \qquad 12.12$$

as F(x', t') is equal to the function F depending of the variables x' and t'.

Applying 12.8 and 12.9 in 1.6 we have:

$$x^2 - c^2 t^2 = \left(\frac{x}{A} - \frac{Bt}{AF}\right)^2 - c^2 \frac{t}{F^2}^2$$

From where we have:

$$x^2 - c^2 t^2 = \left(\frac{x^2}{A^2}\right) - c^2 t^2 \left[\frac{1}{F^2} - \frac{B^2}{A^2 c^2 F^2} + \frac{2Bx}{A^2 c^2 Ft}\right]$$

where making A^2 = 1 in the bracket in arc and $\left[\frac{1}{F^2} - \frac{B^2}{A^2 c^2 F^2} + \frac{2Bx}{A^2 c^2 Ft}\right] = 1$ in the straight bracket we have the equality between both sides of the equal signal of the equation.

Applying A = 1 and 12.10 in $\left[\frac{1}{F^2} - \frac{B^2}{A^2 c^2 F^2} + \frac{2Bx}{A^2 c^2 Ft}\right] = 1$ we have:

$$F = \frac{1}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2 t}}} = F(x, t) \qquad 12.13$$

as F(x, t) is equal to the function F depending on the variables x and t.

We must make the following naming according to 2.5 and 2.6:

$$K' = 1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2 t'} \Rightarrow F = \sqrt{K'} \qquad 12.14$$

$$K = 1 + \frac{v^2}{c^2} - \frac{2vx}{c^2 t} \Rightarrow F = \frac{1}{\sqrt{K}} \qquad 12.15$$

As the equation to F(x', t') from 12.12 and F(x, t) from 12.13 must be equal, we have:

$$F = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2t'}} = \frac{1}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2t}}} \qquad 12.16$$

Thus:

$$\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2t}} \cdot \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2t'}} = 1 \text{ or } \sqrt{K} \cdot \sqrt{K'} = 1 \qquad 12.17$$

Exactly equal to 1.10.

Rewriting the equations 12.6, 12.7, 12.8 and 12.9 according to the function of v, v' and F we have:

$$x = x' + v't' \qquad 12.6$$

$$t = Ft' \qquad 12.7$$

With the respective inverted corrected equations:

$$x' = x - vt \qquad 12.8$$

$$t' = \frac{t}{F} \qquad 12.9$$

We have the equations 12.6, 12.7, 12.8 and 12.9 finals replacing F by the corresponding formulae:

$$x = x' + v't' \qquad 12.6$$

$$t = t'\sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2t'}} \qquad 12.7$$

With the respective inverted final equations:

$$x' = x - vt \qquad 12.8$$

$$t' = t\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2t}} \qquad 12.9$$

That are exactly the equations of the table I

As $v = \frac{B}{F}$ and $v' = B$ then the relations between v and v' are $v = \frac{v'}{F}$ or $v' = v.F$ 12.18

We will transform F (12.12) function of the elements v', x', and t' for F (12.13) function of the elements v, x and t, replacing in 12.12 the equations 12.8, 12.9 and 12.18:

$$F = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2t'}} = \sqrt{1 + \frac{(vF)^2}{c^2} + \frac{2vF(x - vt)}{c^2\,\frac{t}{F}}}$$

$$F = \sqrt{1 + \frac{v^2F^2}{c^2} + \frac{2vxF^2}{c^2t} - \frac{2v^2F^2}{c^2}} = \sqrt{1 + \frac{2vxF^2}{c^2t} - \frac{v^2F^2}{c^2}}$$

$$F^2 = 1 + \frac{2vxF^2}{c^2t} - \frac{v^2F^2}{c^2} \Rightarrow F^2 + \frac{v^2F^2}{c^2} - \frac{2vxF^2}{c^2t} = 1 \Rightarrow F = \frac{1}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2t}}}$$

That is exactly the equation 12.13.

We will transform F (12.13) function of the elements v, x, and t for F (12.12) function of the elements v', x' and t', replacing in 12.13 the equations 12.6, 12.7 and 12.18:

$$F = \frac{1}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx}{c^2t}}} = \frac{1}{\sqrt{1 + \frac{1}{c^2}\left(\frac{v'}{F}\right)^2 - \frac{2v'(x'+v't')}{c^2FFt'}}} = \frac{1}{\sqrt{1 + \frac{v'^2}{c^2F^2} - \frac{2v'x'}{c^2t'F^2} - \frac{2v'^2}{c^2F^2}}}$$

$$F = \frac{1}{\sqrt{1 - \frac{v'^2}{c^2F'^2} - \frac{2v'x'}{c^2t'F'^2}}} \Rightarrow F^2\left(1 - \frac{v'^2}{c^2F2} - \frac{2v'x'}{c^2t'F^2}\right) = 1 \Rightarrow F = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'x'}{c^2t'}}$$

That is exactly the equation 12.12.

We have to calculate the total diferential of F(x', t') (12.12):

$$dF = \frac{\partial F}{\partial x'}dx' + \frac{\partial F}{\partial t'}dt'$$

as:

$$\frac{\partial F}{\partial x'} = \frac{1}{\sqrt{K'}}\frac{v'}{c^2t'} \text{ and } \frac{\partial F}{\partial t'} = -\frac{1}{\sqrt{K'}}\frac{v'}{c^2t'}\frac{x'}{t'} \qquad 12.19$$

we have:

$$dF = \frac{1}{\sqrt{K'}}\frac{v'}{c^2t'}dx' - \frac{1}{\sqrt{K'}}\frac{v'}{c^2t'}\frac{x'}{t'}dt'$$

where applying 1.18 we find:

$$dF = \frac{1}{\sqrt{K'}}\frac{v'}{c^2t'}dx' - \frac{1}{\sqrt{K'}}\frac{v'}{c^2t'}\frac{dx'}{dt'}dt' = o \qquad 12.20$$

From where we conclude that F function of x' and t' is a constant.

We have to calculate the total diferential of F(x, t) (12.13):

$$dF = \frac{\partial F}{\partial x}dx + \frac{\partial F}{\partial t}dt$$

as:

$$\frac{\partial F}{\partial x} = \frac{1}{K^{\frac{3}{2}}}\frac{v}{c^2t} \text{ and } \frac{\partial F}{\partial t} = -\frac{1}{K^{\frac{3}{2}}}\frac{v}{c^2t}\frac{x}{t} \qquad 12.21$$

we have:

$$dF = \frac{1}{K^{\frac{3}{2}}}\frac{v}{c^2t}dx - \frac{1}{K^{\frac{3}{2}}}\frac{v}{c^2t}\frac{x}{t}dt \qquad 12.22$$

where applying 1.13 we find:

$$dF = \frac{1}{K^{\frac{3}{2}}}\frac{v}{c^2 t}dx - \frac{1}{K^{\frac{3}{2}}}\frac{v}{c^2 t}\frac{dx}{dt}dt = o$$

From where we conclude that F function of x and t is a constant.

The equations 1.13 and 1.18 represent to the observers O and O' the principle of constancy of the light speed valid from infinitely small to the infinitely big and mean that in the Undulating Relativity the space and time are measure simultaneously. They shouldn't be interpreted with a dependency between space and time.

The time has its own interpretation that can be understood if we analyze to a determined observer the emission of two rays of light from the instant t=zero. If we add the times we get, for each ray of light, we will get a result without any use for the physics.

If in the instant t = t' = zero, the observer O' emits two rays of light, one alongside the axis x and the other alongside the axis y, after the interval of time t', the rays hit for the observer O', simultaneously, the points A_x and A_y to the distance ct' from the origin, although for the observer O, the points won't be hit simultaneously. For both rays of lights be simultaneous to both observers, they must hit the points that have the same radius in relation to the axis x and that provide the same time for both observers ($t_1 = t_2$ and $t'_1 = t'_2$), which means that only one ray of light is necessary to check the time between the referentials.
According to § 1, both referentials of the observers O and O' are inertial, thus the light spreads in a straight line according to what is demanded by the fundamental axiom of the Undulating Relativity § 12, because of this, the difference in velocities v and v' is due to only a difference in time between the referentials.

$$v = \frac{x - x'}{t} \qquad 1.2 \qquad\qquad v' = \frac{x - x'}{t'} \qquad 1.4$$

We can also relate na inertial referential for which the light spread in a straight line according to what is demanded by the fundamental axiom of the Undulating Relativity, with an accelerated moving referential for which the light spread in a curve line, considering that in this case the difference v and v' isn't due to only the difference of time between the referentials.

According to § 1, if the observer O at the instant t = t' = zero, emits a ray of light from the origin of its referential, after an interval of time t_1, the ray of light hits the point A_1 with coordinates (x_1, y_1, z_1, t_1) to the distance ct_1 of the origin of the observer O, then we have:

$$t'_1 = t_1\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx_1}{c^2 t_1}}$$

After hitting the point A_1 the ray of light still spread in the same direction and in the same way, after an interval of time t_2, the ray of light hits the point A_2 with coordinates ($x_1 + x_2$, $y_1 + y_2$, $z_1 + z_2$, $t_1 + t_2$) to the distance ct_2 to the point A_1, then we have:

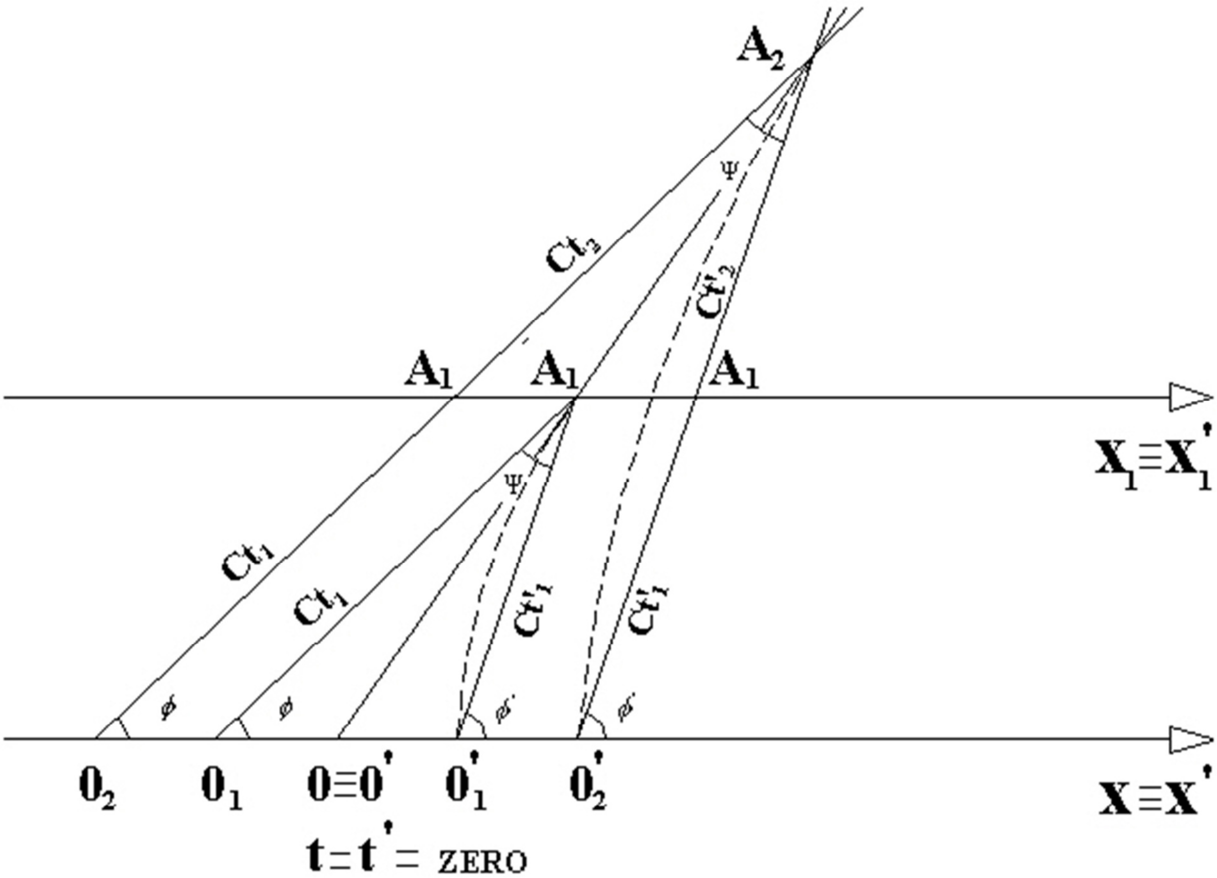
$$\frac{x}{t} = \frac{dx}{dt} = ux \Rightarrow \frac{x_1}{t_1} = \frac{x_2}{t_2} = ux \Rightarrow \sqrt{1 + \frac{v^2}{c^2} - \frac{2vx_1}{c^2 t_1}} = \sqrt{1 + \frac{v^2}{c^2} - \frac{2vx_2}{c^2 t_2}} = \sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}}$$

and with this we get:

$$t'_2 = t_2\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx_2}{c^2 t_2}} = t_2\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}}$$

$$t'_1 + t'_2 = t_1\sqrt{1 + \frac{v^2}{c^2} - \frac{2vx_1}{c^2 t_1}} + t_2\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}} = (t_1 + t_2)\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}} = (t_1 + t_2)\sqrt{1 + \frac{v^2}{c^2} - \frac{2v(x_1 + x_2)}{c^2(t_1 + t_2)}}$$

The geometry of space and time in the Undulating Relativity is summarized in the figure below that can be expanded to A_n points and several observers.



In the figure the angles have a relation $\psi=\phi'-\phi$ and are equal to the following segments:

O_1 to $O\equiv O'$ is equal to $O\equiv O'$ to O'_1 $\left(O_1\leftrightarrow O'_1=vt_1=v't'_1\right)$

O_2 to O_1 is equal to O'_1 to O'_2 $\left(O_2\leftrightarrow O'_2=v(t_1+t_2)=v'(t'_1+t'_2)\rightarrow vt_2=v't'_2=O_2\leftrightarrow O_1+O'_1\leftrightarrow O'_2\right)$
And are parallel to the following segments:

O_2 to A_2 is parallel to O_1 to A_1

O'_2 to A_2 is parallel to O'_1 to A_1

$X\equiv X'$ is parallel to $X_1\equiv X'_1$

The cosine of the angles of inclination ϕ and ϕ' to the rays for the observers O and O' according to 2.3 and 2.4 are:

$$u'x'=\frac{ux-v}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}}\Rightarrow\frac{u'x'}{c}=\frac{\frac{ux}{c}-\frac{v}{c}}{\sqrt{1+\frac{v^2}{c^2}-\frac{2vux}{c^2}}}\Rightarrow\cos\phi'=\frac{\cos\phi-v/c}{\sqrt{1+\frac{v^2}{c^2}-\frac{2v}{c}\cos\phi}}$$

$$\cos\phi'=\frac{\cos\phi-v/c}{\sqrt{K}} \qquad 12.23$$

And with this we have: $sen\phi'=\frac{sen\phi}{\sqrt{K}}$ 12.24

$$ux=\frac{u'x'+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}}\Rightarrow\frac{ux}{c}=\frac{\frac{u'x'}{c}+\frac{v'}{c}}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'}{c^2}}}\Rightarrow\cos\phi=\frac{\cos\phi'+v'/c}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'}{c}\cos\phi'}}$$

$$\cos\phi=\frac{\cos\phi'+v'/c}{\sqrt{K'}} \qquad 12.25$$

And with this we have $sen\phi = \frac{sen\phi'}{\sqrt{K'}}$ 12.26

The cosine of the angle ψ with intersection of rays equal to:

$$cos\psi = \frac{1-\frac{vux}{c^2}}{\sqrt{K}} = \frac{1+\frac{v'u'x'}{c^2}}{\sqrt{K'}} = \frac{1-\frac{v}{c}cos\phi}{\sqrt{K}} = \frac{1+\frac{v'}{c}cos\phi'}{\sqrt{K'}} \quad 12.27$$

And with this we have: $sen\psi = \frac{v}{c}\frac{sen\phi}{\sqrt{K}} = \frac{v'}{c}\frac{sen\phi'}{\sqrt{K'}}$ 12.28

The invariance of the $cos\psi$ shows the harmony of all adopted hypotheses for space and time in the Undulating Relativity.

The $cos\psi$ is equal to the Jacobians of the transformations for the space and time of the picture I, where the radicals

$\sqrt{K} = \sqrt{1+\frac{v^2}{c^2}-\frac{2vx}{c^2t}}$ and $\sqrt{K'} = \sqrt{1+\frac{v'^2}{c^2}+\frac{2v'x'}{c^2t'}}$ are considered variables and are derived.

$$cos\psi = J = \frac{\partial x'^i}{\partial x^j} = \frac{\partial(x',y',z',t')}{\partial(x,y,z,t)} = \begin{vmatrix} 1 & 0 & 0 & -v \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \frac{-v/c^2}{\sqrt{K}} & 0 & 0 & \frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vx}{c^2t}\right) \end{vmatrix} = \frac{1-\frac{vx}{c^2t}}{\sqrt{K}} = \frac{1-\frac{vux}{c^2}}{\sqrt{K}} \quad 8.8$$

$$cos\psi = J' = \frac{\partial x^k}{\partial x'^l} = \frac{\partial(x,y,z,t)}{\partial(x',y',z',t')} = \begin{vmatrix} 1 & 0 & 0 & v' \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \frac{v'/c^2}{\sqrt{K'}} & 0 & 0 & \frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'x'}{c^2t'}\right) \end{vmatrix} = \frac{1+\frac{v'x'}{c^2t'}}{\sqrt{K'}} = \frac{1+\frac{v'u'x'}{c^2}}{\sqrt{K'}} \quad 8.8$$

§13 Richard C. Tolman

The §4 Transformations of the Momenta of Undulating Relativity was developed based on the experience conducted by Lewis and Tolman, according to the reference [3]. Where the collision of two spheres preserving the principle of conservation of energy and the principle of conservation of momenta, shows that the mass is a function of the velocity according to:

$$m = \frac{m_o}{\sqrt{1-\frac{(u)^2}{c^2}}}$$

where m_o is the mass of the sphere when in resting position and $u = |\vec{u}| = \sqrt{\vec{u}\vec{u}}$ the module of its speed.

Analyzing the collision between two identical spheres when in relative resting position, that for the observer O' are named S'$_1$ and S'$_2$ are moving along the axis x' in the contrary way with the following velocities before the collision:

Table 1

Esphere S'$_1$	Esphere S'$_2$
$u'x'_1 = v'$	$u'x_2' = -v'$
$u'y'_1 = zero$	$u'y'_2 = zero$
$u'z'_1 = zero$	$u'z'_2 = zero$

For the observer O the same spheres are named S$_1$ and S$_2$ and have the velocities $(ux_1,\ ux_2,\ uy_i = uz_i = zero)$ before the collision calculated according to the table 2 as follows:

The velocity ux_1 of the sphere S_1 is equals to:

$$ux_1=\frac{u'x'_1+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'_1}{c^2}}}=\frac{v'+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'v'}{c^2}}}=\frac{2v'}{\sqrt{1+\frac{3v'^2}{c^2}}}.$$

The transformation from v' to v according to 1.20 from Table 2 is:

$$v=\frac{v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'_1}{c^2}}}=\frac{v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'v'}{c^2}}}=\frac{v'}{\sqrt{1+\frac{3v'^2}{c^2}}}.$$

That applied in ux_1 supplies:

$$ux_1=2\left(\frac{v'}{\sqrt{1+\frac{3v'^2}{c^2}}}\right)=2v$$

The velocity ux_2 of the sphere S_2 is equal to:

$$ux_2=\frac{u'x'_2+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'_2}{c^2}}}=\frac{-v'+v'}{\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'(-v')}{c^2}}}=zero$$

Table 2

Sphere S_1	Sphere S_2
$ux_1=\frac{2v'}{\sqrt{1+\frac{3v'}{c^2}}}=2v$	$ux_2=zero$
$uy_1=zero$	$uy_2=zero$
$uz_1=zero$	$uz_2=zero$

For the observers O and O' the two spheres have the same mass when in relative resting position. And for the observer O' the two spheres collide with velocities of equal module and opposite direction because of this the momenta $(p'_1=p'_2)$ null themselves during the collision, forming for a brief time $(\Delta t')$ only one body of mass

$$m_0=m'_1+m'_2.$$

According to the principle of conservation of momenta for the observer O we will have to impose that the momenta before the collision are equal to the momenta after the collision, thus:

$$m_1ux_1+m_2ux_2=(m_1+m_2)w$$

Where for the observer O, w is the arbitrary velocity that supposedly for a brief time (Δt) will also see the masses united $(m=m_1+m_2)$ moving. As the masses m_i have different velocities and the masses vary according to their own velocities, this equation cannot be simplified algebraically, having this variation of masses:

To the left side of the equal sign in the equation we have:

$$u=ux_1=2v$$

$$m_1 = \frac{m_o}{\sqrt{1-\frac{(u)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{(ux_1)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{(2v)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{4v^2}{c^2}}}$$

$$u = ux_2 = zero$$

$$m_2 = \frac{m_o}{\sqrt{1-\frac{(u)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{(ux_2)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{(zero)^2}{c^2}}} = m_0$$

To the right side of the equal sign in the equation we have:

$$u = w$$

$$m_1 = \frac{m_o}{\sqrt{1-\frac{(u)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{(w)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{w^2}{c^2}}}$$

$$m_2 = \frac{m_o}{\sqrt{1-\frac{(u)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{(w)^2}{c^2}}} = \frac{m_o}{\sqrt{1-\frac{w^2}{c^2}}}$$

Applying in the equation of conservation of momenta we have:

$$m_1 ux_1 + m_2 ux_2 = (m_1 + m_2)w = m_1 w + m_2 w$$

$$\frac{m_0}{\sqrt{1-\frac{4v^2}{c^2}}} 2v + m_0 .0 = \frac{m_0}{\sqrt{1-\frac{w^2}{c^2}}} w + \frac{m_0}{\sqrt{1-\frac{w^2}{c^2}}} w$$

From where we have:

$$\frac{2m_0 v}{\sqrt{1-\frac{4v^2}{c^2}}} = \frac{2m_0 w}{\sqrt{1-\frac{w^2}{c^2}}} \Rightarrow \frac{v}{\sqrt{1-\frac{4v^2}{c^2}}} = \frac{w}{\sqrt{1-\frac{w^2}{c^2}}}$$

$$w = \frac{v}{\sqrt{1-\frac{3v^2}{c^2}}}$$

As $w \neq v$ for the observer O the masses united $(m = m_1 + m_2)$ wouldn't move momentarily alongside to the observer O' which is conceivable if we consider that the instants $\Delta t \neq \Delta t'$ are different where supposedly the masses would be in a resting position from the point of view of each observer and that the mass acting with velocity 2v is bigger than the mass in resting position.

If we operate with these variables in line we would have:

$$m_1 ux_1 + m_2 ux_2 = (m_1 + m_2)w = m_1 w + m_2 w$$

$$\frac{m_0}{\sqrt{1-\frac{1}{c^2}\left(\frac{2v'}{\sqrt{1+\frac{3v'}{c^2}}}\right)^2}} \frac{2v'}{\sqrt{1+\frac{3v'}{c^2}}} + m_0 .0 = \frac{m_0}{\sqrt{1-\frac{w^2}{c^2}}} w + \frac{m_0}{\sqrt{1-\frac{w^2}{c^2}}} w = \frac{2m_0 w}{\sqrt{1-\frac{w^2}{c^2}}}$$

$$\frac{2m_0 v'}{\sqrt{\left(1+\frac{3v'}{c^2}\right)\left(1-\frac{1}{c^2}\left(\frac{4v'^2}{\left(1+\frac{3v'}{c^2}\right)}\right)\right)}}=\frac{2m_0 w}{\sqrt{1-\frac{w^2}{c^2}}}$$

$$\frac{2m_0 v'}{\sqrt{1+\frac{3v'}{c^2}-\frac{4v'^2}{c^2}}}=\frac{2m_0 w}{\sqrt{1-\frac{w^2}{c^2}}}$$

$$\frac{2m_0 v'}{\sqrt{1-\frac{v'^2}{c^2}}}=\frac{2m_0 w}{\sqrt{1-\frac{w^2}{c^2}}}$$

From where we conclude that $w=v'$ which must be equal to the previous value of w, that is:

$$w=v'=\frac{v}{\sqrt{1-\frac{3v^2}{c^2}}}$$

A relation between v and v' that is obtained from Table 2 when $ux_1=2v$ that corresponds for the observer O to the velocity acting over the sphere in resting position.

§14 Velocities composition

Reference – Millennium Relativity

URL: http://www.mrelativity.net/MBriefs/VComp_Sci_Estab_Way.htm

Let's write the transformations of Hendrik A. Lorentz for space and time in the Special Theory of Relativity:

$x'=\frac{x-vt}{\sqrt{1-\frac{v^2}{c^2}}}$	14.1a	$x=\frac{x'+vt'}{\sqrt{1-\frac{v^2}{c^2}}}$	14.3a
$y'=y$	14.1b	$y=y'$	14.3b
$z'=z$	14.1c	$z=z'$	14.3c
$t'=\frac{t-\frac{vx}{c^2}}{\sqrt{1-\frac{v^2}{c^2}}}$	14.2	$t=\frac{t'+\frac{vx'}{c^2}}{\sqrt{1-\frac{v^2}{c^2}}}$	14.4

From them we obtain the equations of velocity transformation:

$u'x'=\frac{ux-v}{1-\frac{vux}{c^2}}$	14.5a	$ux=\frac{u'x'+v}{1+\frac{vu'x'}{c^2}}$	14.6a
$u'y'=\frac{uy\sqrt{1-\frac{v^2}{c^2}}}{1-\frac{vux}{c^2}}$	14.5b	$uy=\frac{u'y'\sqrt{1-\frac{v^2}{c^2}}}{1+\frac{vu'x'}{c^2}}$	14.6b
$u'z'=\frac{uz\sqrt{1-\frac{v^2}{c^2}}}{1-\frac{vux}{c^2}}$	14.5c	$uz=\frac{u'z'\sqrt{1-\frac{v^2}{c^2}}}{1+\frac{vu'x'}{c^2}}$	14.6c

Let's consider that in relation to the observer O' an object moves with velocity:

$$u'x'=1{,}5.10^5\,km/s(=0{,}50c).$$

And that the velocity of the observer O' in relation to the observer O is:

$v=1,5.10^5\,km/s(=0,50c)$.

The velocity ux of the object in relation to the observer O must be calculated by the formula 14.6a:

$$ux=\frac{u'x'+v}{1+\frac{vu'x'}{c^2}}=\frac{1,5.10^5+1,5.10^5}{1+\frac{1,5.10^5.1,5.10^5}{\left(3,0.10^5\right)^2}}=2,4.10^5\,km/s(=0,80c).$$

Where we use $c=3,0.10^5\,km/s(=1,00c)$.

Considering that the object has moved during one second in relation to the observer O $(t=1,00s)$ we can then with 14.2 calculate the time passed to the observer O':

$$t'=\frac{t-\frac{vx}{c^2}}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{t\left(1-\frac{vux}{c^2}\right)}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{1,00\left(1-\frac{1,5.10^5.2,4.10^5}{\left(3,0.10^5\right)^2}\right)}{\sqrt{1-\frac{\left(1,5.10^5\right)^2}{\left(3,0.10^5\right)^2}}}=\frac{0,60}{\sqrt{0,75}}\Rightarrow t'=0,693s.$$

To the observer O the observer O' is away the distance d given by the formula:

$d=vt=1,5.10^5.1,00=1,5.10^5\,km$.

To the observer O' the observer O is away the distance d' given by the formula:

$$d'=vt'=1,5.10^5.\frac{0,60}{\sqrt{0,75}}=1,03923.10^5\,km.$$

To the distance of the object (d_O, d'_O) in relation to the observers O and O' is given by the formulae:

$d_O=uxt=2,4.10^5.1,00=2,4.10^5\,km$.

$$d'_O=u'x't'=1,5.10^5.\frac{0,60}{\sqrt{0,75}}=1,03923.10^5\,km.$$

To the observer O the distance between the object and the observer O' is given by the formula:

$\Delta d=d_O-d=2,4.10^5-1,5.10^5=0,90.10^5\,km$.

To the observer O the velocity of the object in relation to the observer O' is given by:

$$\frac{\Delta d}{t}=\frac{0,90.10^5\,km}{1,00s}=0,90.10^5\,km/s$$ (=0,30c)

Relating the times t and t' using the formula $t'=t\sqrt{1-\frac{v^2}{c^2}}$ is only possible and exclusively when $ux=v$ and $u'x'=zero$ what isn't the case above, to make it possible to understand this we write the equations 14.2 and 14.4 in the formula below:

$t'=\frac{t\left(1-\frac{v}{c}cos\phi\right)}{\sqrt{1-\frac{v^2}{c^2}}}$	14.2	$t=\frac{t'\left(1+\frac{v}{c}cos\phi'\right)}{\sqrt{1-\frac{v^2}{c^2}}}$	14.4

Where $cos\phi=\frac{x}{ct}$ and $cos\phi'=\frac{x'}{ct'}$.

The equations above can be written as:

$$t'=f(t,\phi) \text{ e } t=f'(t',\phi') \qquad 14.7$$

In each referential of the observers O and O' the light propagation creates a sphere with radius ct and ct' that intercept each other forming a circumference that propagates with velocity c. The radius ct and ct' and the positive way of the axis x and x' form the angles ϕ and ϕ' constant between the referentials. If for the same pair of referentials te angles were variable the time would be alleatory and would become useless for the Physics. In the equation $t'=f(t,\phi)$ we have t' identical function of t and ϕ, if we have in it ϕ constant and t' varies according to t we get the common relation between the times t and t' between two referentials, however if we have t constant and t' varies according to ϕ we will have for each value of ϕ one value of t' and t between two different referentials, and this analysis is also valid for $t=f'(t',\phi')$.

Dividing 14.5a by c we have:

$$\frac{u'x'}{c}=\frac{\frac{ux}{c}-\frac{v}{c}}{1-\frac{vux}{c^2}}\Rightarrow cos\phi'=\frac{cos\phi-\frac{v}{c}}{1-\frac{v}{c}cos\phi}. \qquad 14.8$$

Where $cos\phi=\frac{x}{ct}=\frac{ux}{c}$ and $cos\phi'=\frac{x'}{ct'}=\frac{u'x'}{c}$.

Isolating the velocity we have:

$$\frac{v}{c}=\frac{(cos\phi-cos\phi')}{(1-cos\phi cos\phi')} \quad \text{or} \quad v=\frac{ux-u'x'}{1-\frac{uxu'x'}{c^2}} \qquad 14.9$$

From where we conclude that we must have angles ϕ and ϕ' constant so that we have the same velocity between the referentials.

This demand of constant angles between the referentials must solve the controversies of Herbert Dingle.

§15 Invariance

The transformations to the space and time of table I, group 1.2 plus 1.7, in the matrix form is written like this:

$$\begin{bmatrix} x' \\ y' \\ z' \\ t' \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & -v \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix}\begin{bmatrix} x \\ y \\ z \\ t \end{bmatrix} \qquad 15.1$$

That written in the form below represents the same coordinate transformations:

$$\begin{bmatrix} x' \\ y' \\ z' \\ ct' \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix}\begin{bmatrix} x \\ y \\ z \\ ct \end{bmatrix} \qquad 15.2$$

We call as:

$$x'=x'^{i}=\begin{bmatrix} x' \\ y' \\ z' \\ ct' \end{bmatrix}=\begin{bmatrix} x'^{1} \\ x'^{2} \\ x'^{3} \\ cx'^{4} \end{bmatrix},\ \alpha=\alpha_{ij}=\begin{bmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix},\ x=x^{j}=\begin{bmatrix} x \\ y \\ z \\ ct \end{bmatrix}=\begin{bmatrix} x^{1} \\ x^{2} \\ x^{3} \\ cx^{4} \end{bmatrix} \qquad 15.3$$

That are the functions $x'^i = x'^i\left(x^j\right) = x'^i\left(x^1, x^2, x^3, cx^4\right) = x'^i\left(x, y, z, ct\right)$ 15.4

That in the symbolic form is written:

$x' = \alpha . x$ or in the indexed form $x'^i = \sum_{j=1}^{4} \alpha_{ij} x^j \Rightarrow x'^i = \alpha_{ij} x^j$ 15.5

Where we use Einstein's sum convention.

The transformations to the space and time of table I, group 1.4 plus 1.8, in the matrix form is written:

$$\begin{bmatrix} x \\ y \\ z \\ t \end{bmatrix} = \begin{bmatrix} 1\,0\,0 & v' \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K'} \end{bmatrix} \begin{bmatrix} x' \\ y' \\ z' \\ t' \end{bmatrix} \quad 15.6$$

That written in the form below represents the same coordinate transformations:

$$\begin{bmatrix} x \\ y \\ z \\ ct \end{bmatrix} = \begin{bmatrix} 1\,0\,0 & v'/c \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K'} \end{bmatrix} \begin{bmatrix} x' \\ y' \\ z' \\ ct' \end{bmatrix} \quad 15.7$$

That we call as:

$$x = x^k = \begin{bmatrix} x \\ y \\ z \\ ct \end{bmatrix} = \begin{bmatrix} x^1 \\ x^2 \\ x^3 \\ cx^4 \end{bmatrix},\ \alpha' = \alpha'_{kl} = \begin{bmatrix} 1\,0\,0 & v'/c \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K'} \end{bmatrix},\ x' = x'^l = \begin{bmatrix} x' \\ y' \\ z' \\ ct' \end{bmatrix} = \begin{bmatrix} x'^1 \\ x'^2 \\ x'^3 \\ cx'^4 \end{bmatrix} \quad 15.8$$

That are the functions $x^k = x^k\left(x'^l\right) = x^k\left(x'^1, x'^2, x'^3, cx'^4\right) = x^k\left(x', y', z', ct'\right)$ 15.9

That in the symbolic form is written:

$x = \alpha' . x'$ or in the indexed form $x^k = \sum_{l=1}^{4} \alpha'_{kl} x'^l \Rightarrow x^k = \alpha'_{kl} x'^l$ 15.10

Being $\sqrt{K} = \sqrt{1 + \frac{v^2}{c^2} - \frac{2vx^1}{c^2 x^4}}$ (1.7), $\sqrt{K'} = \sqrt{1 + \frac{v'^2}{c^2} + \frac{2v' x'^1}{c^2 x'^4}}$ (1.8) and $\sqrt{K}.\sqrt{K'} = 1$ (1.10).

The transformation matrices $\alpha = \alpha_{ij}$ and $\alpha' = \alpha'_{kl}$ have the properties:

$$\alpha.\alpha' = \alpha_{ij}\alpha'_{kl} = \sum_{j=1}^{4} \alpha_{ij}\alpha'_{jl} = \begin{bmatrix} 1\,0\,0 & -v/c \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K} \end{bmatrix} \begin{bmatrix} 1\,0\,0 & v'/c \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K'} \end{bmatrix} = \begin{bmatrix} 1\,0\,0\,0 \\ 0\,1\,0\,0 \\ 0\,0\,1\,0 \\ 0\,0\,0\,1 \end{bmatrix} = I = \delta_l^i \quad 15.11$$

$$\alpha^t \alpha'^t = \alpha_{ji}\alpha'_{lk} = \sum_{i=1}^{4} \alpha_{ji}\alpha'_{ik} = \begin{bmatrix} 1 & 0\,0 & 0 \\ 0 & 1\,0 & 0 \\ 0 & 0\,1 & 0 \\ -v/c & 0\,0 & \sqrt{K} \end{bmatrix} \begin{bmatrix} 1 & 0\,0 & 0 \\ 0 & 1\,0 & 0 \\ 0 & 0\,1 & 0 \\ v'/c & 0\,0 & \sqrt{K'} \end{bmatrix} = \begin{bmatrix} 1\,0\,0\,0 \\ 0\,1\,0\,0 \\ 0\,0\,1\,0 \\ 0\,0\,0\,1 \end{bmatrix} = I = \delta_k^j \quad 15.12$$

Where $\alpha^t = \alpha_{ji}$ is the transposed matrix of $\alpha = \alpha_{ij}$ and $\alpha'^t = \alpha'_{lk}$ is the transpose matrix of $\alpha' = \alpha'_{kl}$ and δ is the Kronecker's delta.

$$\alpha'.\alpha = \alpha'_{kl}\alpha_{ij} = \sum_{l=1}^{4} \alpha'_{kl}\alpha_{lj} = \begin{bmatrix} 1\,0\,0 & v'/c \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K'} \end{bmatrix} \begin{bmatrix} 1\,0\,0 & -v/c \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & \sqrt{K} \end{bmatrix} = \begin{bmatrix} 1\,0\,0\,0 \\ 0\,1\,0\,0 \\ 0\,0\,1\,0 \\ 0\,0\,0\,1 \end{bmatrix} = I = \delta_j^k \quad 15.13$$

$$\alpha'^t \alpha^t = \alpha'_{lk}\alpha_{ji} = \sum_{k=1}^{4}\alpha'_{lk}\alpha_{ki} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ v'/c & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -v/c & 0 & 0 & \sqrt{K} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} = I = \delta_i^l \qquad 15.14$$

Where $\alpha'^t = \alpha'_{lk}$ is the transposed matrix of $\alpha' = \alpha'_{kl}$ and $\alpha^t = \alpha_{ji}$ is the transposed matrix of $\alpha = \alpha_{ij}$ and δ is the Kronecker's delta.

Observation: the matrices α_{ij} and α'_{kl} are inverse of one another but are not orthogonal, that is: $\alpha_{ji} \neq \alpha'_{kl}$ and $\alpha_{ij} \neq \alpha'_{lk}$.

The partial derivatives $\frac{\partial x'^i}{\partial x^j}$ of the total differential $dx'^i = \frac{\partial x'^i}{\partial x^j}dx^j$ of the coordinate components that correlate according to $x'^i = x'^i\left(x^j\right)$, where in the transformation matrix $\alpha = \alpha_{ij}$ the radical $\sqrt{K}$ is considered constant and equal to:

Table 10, partial derivatives of the coordinate components:

$\frac{\partial x'^i}{\partial x^j} = \frac{\partial x'^1}{\partial x^j} =$	$\frac{\partial x'^1}{\partial x^1} = 1$	$\frac{\partial x'^1}{\partial x^2} = 0$	$\frac{\partial x'^1}{\partial x^3} = 0$	$\frac{\partial x'^1}{\partial x^4} = -\frac{v}{c}$
$\frac{\partial x'^i}{\partial x^j} = \frac{\partial x'^2}{\partial x^j} =$	$\frac{\partial x'^2}{\partial x^1} = 0$	$\frac{\partial x'^2}{\partial x^2} = 1$	$\frac{\partial x'^2}{\partial x^3} = 0$	$\frac{\partial x'^2}{\partial x^4} = 0$
$\frac{\partial x'^i}{\partial x^j} = \frac{\partial x'^3}{\partial x^j} =$	$\frac{\partial x'^3}{\partial x^1} = 0$	$\frac{\partial x'^3}{\partial x^2} = 0$	$\frac{\partial x'^3}{\partial x^3} = 1$	$\frac{\partial x'^3}{\partial x^4} = 0$
$\frac{\partial x'^i}{\partial x^j} = \frac{\partial x'^4}{\partial x^j} =$	$\frac{\partial x'^4}{\partial x^1} = 0$	$\frac{\partial x'^4}{\partial x^2} = 0$	$\frac{\partial x'^4}{\partial x^3} = 0$	$\frac{\partial x'^4}{\partial x^4} = \sqrt{K}$

The total differential of the coordinates in the matrix form is equal to:

$$\begin{bmatrix} dx'^1 \\ dx'^2 \\ dx'^3 \\ cdx'^4 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix}\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} \qquad 15.15$$

That we call as:

$$dx' = dx'^i = \begin{bmatrix} dx'^1 \\ dx'^2 \\ dx'^3 \\ cdx'^4 \end{bmatrix}, \quad A = A_j^i = \frac{\partial x'^i}{\partial x^j} = \begin{bmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix}, \quad dx = dx^j = \begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} \qquad 15.16$$

Then we have $dx' = Adx \Rightarrow dx'^i = \sum_{j=1}^{4} A_j^i dx^j \Rightarrow dx'^i = \frac{\partial x'^i}{\partial x^j}dx^j$ 15.17

The partial derivatives $\frac{\partial x^k}{\partial x'^l}$ of the total differential $dx^k = \frac{\partial x^k}{\partial x'^l}dx'^l$ of the coordinate components that correlate according to $x^k = x^k\left(x'^l\right)$, where in the transformation matrix $\alpha' = \alpha'_{kl}$ the radical $\sqrt{K'}$ is considered constant and equal to:

Table 11 partial derivatives of the coordinate components:

$\frac{\partial x^k}{\partial x'^l}=\frac{\partial x^1}{\partial x'^l}=$	$\frac{\partial x^1}{\partial x'^1}=1$	$\frac{\partial x^1}{\partial x'^2}=0$	$\frac{\partial x^1}{\partial x'^3}=0$	$\frac{\partial x^1}{\partial x'^4}=\frac{v'}{c}$
$\frac{\partial x^k}{\partial x'^l}=\frac{\partial x^2}{\partial x'^l}=$	$\frac{\partial x^2}{\partial x'^1}=0$	$\frac{\partial x^2}{\partial x'^2}=1$	$\frac{\partial x^2}{\partial x'^3}=0$	$\frac{\partial x^2}{\partial x'^4}=0$
$\frac{\partial x^k}{\partial x'^l}=\frac{\partial x^3}{\partial x'^l}=$	$\frac{\partial x^3}{\partial x'^1}=0$	$\frac{\partial x^3}{\partial x'^2}=0$	$\frac{\partial x^3}{\partial x'^3}=1$	$\frac{\partial x^3}{\partial x'^4}=0$
$\frac{\partial x^k}{\partial x'^l}=\frac{\partial x^4}{\partial x'^l}=$	$\frac{\partial x^4}{\partial x'^1}=0$	$\frac{\partial x^4}{\partial x'^2}=0$	$\frac{\partial x^4}{\partial x'^3}=0$	$\frac{\partial x^4}{\partial x'^4}=\sqrt{K'}$

The total differential of the coordinates in the matrix form is equal to:

$$\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & v'/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} dx'^1 \\ dx'^2 \\ dx'^3 \\ cdx'^4 \end{bmatrix} \tag{15.18}$$

That we call as:

$$dx=dx^k=\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix},\ A'=A'^k_l=\frac{\partial x^k}{\partial x'^l}=\begin{bmatrix} 1 & 0 & 0 & v'/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix},\ dx'=dx'^l=\begin{bmatrix} dx'^1 \\ dx'^2 \\ dx'^3 \\ cdx'^4 \end{bmatrix} \tag{15.19}$$

Then we have: $dx=A'dx' \Rightarrow dx^k=\sum_{l=1}^{4} A'^k_l\, dx'^l \Rightarrow dx^k=\frac{\partial x^k}{\partial x'^l}dx'^l$ 15.20

The Jacobians of the transformations 15.15 and 15.18 are:

$$J=\frac{\partial x'^i}{\partial x^j}=\frac{\partial\left(x'^1,x'^2,x'^3,x'^4\right)}{\partial\left(x^1,x^2,x^3,x^4\right)}=\begin{vmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{vmatrix}=\sqrt{K} \tag{15.21}$$

$$J'=\frac{\partial x^k}{\partial x'^l}=\frac{\partial\left(x^1,x^2,x^3,x^4\right)}{\partial\left(x'^1,x'^2,x'^3,x'^4\right)}=\begin{vmatrix} 1 & 0 & 0 & v'/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{vmatrix}=\sqrt{K'} \tag{15.22}$$

Where $\sqrt{K}=\sqrt{1+\frac{v^2}{c^2}-\frac{2vux^1}{c^2}}$ (2.5), $\sqrt{K'}=\sqrt{1+\frac{v'^2}{c^2}+\frac{2v'u'x'^1}{c^2}}$ (2.6) and $\sqrt{K}.\sqrt{K'}=1$ (1.23).

The matrices of the transformation A and A' also have the properties 15.11, 15.12, 15.13 and 15.14 of the matrices α and α'.

From the function $\phi=\phi\left(x^k\right)=\phi'=\phi'\left[x^k\left(x'^l\right)\right]$ where the coordinates correlate in the form $x^k=x^k\left(x'^l\right)$ we have $\frac{\partial\phi}{\partial x'^l}=\frac{\partial\phi}{\partial x^k}\frac{\partial x^k}{\partial x'^l}$ described as:

$$\frac{\partial\phi}{\partial x'^1}=\frac{\partial\phi}{\partial x^k}\frac{\partial x^k}{\partial x'^1}=\frac{\partial\phi}{\partial x^1}\frac{\partial x^1}{\partial x'^1}+\frac{\partial\phi}{\partial x^2}\frac{\partial x^2}{\partial x'^1}+\frac{\partial\phi}{\partial x^3}\frac{\partial x^3}{\partial x'^1}+\frac{\partial\phi}{\partial x^4}\frac{\partial x^4}{\partial x'^1}$$
$$\frac{\partial\phi}{\partial x'^2}=\frac{\partial\phi}{\partial x^k}\frac{\partial x^k}{\partial x'^2}=\frac{\partial\phi}{\partial x^1}\frac{\partial x^1}{\partial x'^2}+\frac{\partial\phi}{\partial x^2}\frac{\partial x^2}{\partial x'^2}+\frac{\partial\phi}{\partial x^3}\frac{\partial x^3}{\partial x'^2}+\frac{\partial\phi}{\partial x^4}\frac{\partial x^4}{\partial x'^2}$$
$$\frac{\partial\phi'}{\partial x'^3}=\frac{\partial\phi}{\partial x^k}\frac{\partial x^k}{\partial x'^3}=\frac{\partial\phi}{\partial x^1}\frac{\partial x^1}{\partial x'^3}+\frac{\partial\phi}{\partial x^2}\frac{\partial x^2}{\partial x'^3}+\frac{\partial\phi}{\partial x^3}\frac{\partial x^3}{\partial x'^3}+\frac{\partial\phi}{\partial x^4}\frac{\partial x^4}{\partial x'^3}$$
$$\frac{\partial\phi}{\partial x'^4}=\frac{\partial\phi}{\partial x^k}\frac{\partial x^k}{\partial x'^4}=\frac{\partial\phi}{\partial x^1}\frac{\partial x^1}{\partial x'^4}+\frac{\partial\phi}{\partial x^2}\frac{\partial x^2}{\partial x'^4}+\frac{\partial\phi}{\partial x^3}\frac{\partial x^3}{\partial x'^4}+\frac{\partial\phi'}{\partial x^4}\frac{\partial x^4}{\partial x'^4}$$

That in the matrix form and without presenting the function ϕ becomes:

$$\frac{\partial\phi}{\partial x'^1}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{\partial x'^4}\right]=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{\partial x^4}\right]\begin{bmatrix}\frac{\partial x^1}{\partial x'^1}=1 & \frac{\partial x^1}{\partial x'^2}=0 & \frac{\partial x^1}{\partial x'^3}=0 & \frac{\partial x^1}{\partial x'^4}=v' \\ \frac{\partial x^2}{\partial x'^1}=0 & \frac{\partial x^2}{\partial x'^2}=1 & \frac{\partial x^2}{\partial x'^3}=0 & \frac{\partial x^2}{\partial x'^4}=0 \\ \frac{\partial x^3}{\partial x'^1}=0 & \frac{\partial x^3}{\partial x'^2}=0 & \frac{\partial x^3}{\partial x'^3}=1 & \frac{\partial x^3}{\partial x'^4}=0 \\ \frac{\partial x^4}{\partial x'^1}=\frac{v'}{c^2\sqrt{K'}} & \frac{\partial x^4}{\partial x'^2}=0 & \frac{\partial x^4}{\partial x'^3}=0 & \frac{\partial x^4}{\partial x'^4}=\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'^1}{c^2}\right)\end{bmatrix}$$

Where replacing the items below:

$$\frac{\partial x^4}{\partial x'^1}=\frac{v'}{c^2\sqrt{K'}}=\frac{v}{c^2}$$

$$\frac{\partial x^1}{\partial x'^4}=v'=\frac{v}{\sqrt{K}}$$

$$\frac{\partial x^4}{\partial x'^4}=\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'^1}{c^2}\right)=\frac{\partial x'^4}{\partial x^4}=\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux^1}{c^2}\right)$$

Observation: this last relation shows that the time varies in an equal form between the referentials.

We get:

$$\frac{\partial\phi}{\partial x'^1}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{\partial x'^4}\right]=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{\partial x^4}\right]\begin{bmatrix}\frac{\partial x^1}{\partial x'^1}=1 & \frac{\partial x^1}{\partial x'^2}=0 & \frac{\partial x^1}{\partial x'^3}=0 & \frac{\partial x^1}{\partial x'^4}=\frac{v}{\sqrt{K}} \\ \frac{\partial x^2}{\partial x'^1}=0 & \frac{\partial x^2}{\partial x'^2}=1 & \frac{\partial x^2}{\partial x'^3}=0 & \frac{\partial x^2}{\partial x'^4}=0 \\ \frac{\partial x^3}{\partial x'^1}=0 & \frac{\partial x^3}{\partial x'^2}=0 & \frac{\partial x^3}{\partial x'^3}=1 & \frac{\partial x^3}{\partial x'^4}=0 \\ \frac{\partial x^4}{\partial x'^1}=\frac{v}{c^2} & \frac{\partial x^4}{\partial x'^2}=0 & \frac{\partial x^4}{\partial x'^3}=0 & \frac{\partial x^4}{\partial x'^4}=\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux^1}{c^2}\right)\end{bmatrix}$$

That is the group 8.1 plus 8.3 of the table 9, differential operators, in the matrix form.

From the function $\phi'=\phi'(x'^i)=\phi=\phi[x'^i(x^j)]$ where the coordinates correlate in the form $x'^i=x'^i(x^j)$ we have $\frac{\partial\phi'}{\partial x^j}=\frac{\partial\phi'}{\partial x'^i}\frac{\partial x'^i}{\partial x^j}$ described as:

$$\frac{\partial\phi'}{\partial x^1}=\frac{\partial\phi'}{\partial x'^i}\frac{\partial x'^i}{\partial x^1}=\frac{\partial\phi'}{\partial x'^1}\frac{\partial x'^1}{\partial x^1}+\frac{\partial\phi'}{\partial x'^2}\frac{\partial x'^2}{\partial x^1}+\frac{\partial\phi'}{\partial x'^3}\frac{\partial x'^3}{\partial x^1}+\frac{\partial\phi'}{\partial x'^4}\frac{\partial x'^4}{\partial x^1}$$
$$\frac{\partial\phi'}{\partial x^2}=\frac{\partial\phi'}{\partial x'^i}\frac{\partial x'^i}{\partial x^2}=\frac{\partial\phi'}{\partial x'^1}\frac{\partial x'^1}{\partial x^2}+\frac{\partial\phi'}{\partial x'^2}\frac{\partial x'^2}{\partial x^2}+\frac{\partial\phi'}{\partial x'^3}\frac{\partial x'^3}{\partial x^2}+\frac{\partial\phi'}{\partial x'^4}\frac{\partial x'^4}{\partial x^2}$$
$$\frac{\partial\phi'}{\partial x^3}=\frac{\partial\phi'}{\partial x'^i}\frac{\partial x'^i}{\partial x^3}=\frac{\partial\phi'}{\partial x'^1}\frac{\partial x'^1}{\partial x^3}+\frac{\partial\phi'}{\partial x'^2}\frac{\partial x'^2}{\partial x^3}+\frac{\partial\phi'}{\partial x'^3}\frac{\partial x'^3}{\partial x^3}+\frac{\partial\phi'}{\partial x'^4}\frac{\partial x'^4}{\partial x^3}$$
$$\frac{\partial\phi'}{\partial x^4}=\frac{\partial\phi'}{\partial x'^i}\frac{\partial x'^i}{\partial x^4}=\frac{\partial\phi'}{\partial x'^1}\frac{\partial x'^1}{\partial x^4}+\frac{\partial\phi'}{\partial x'^2}\frac{\partial x'^2}{\partial x^4}+\frac{\partial\phi'}{\partial x'^3}\frac{\partial x'^3}{\partial x^4}+\frac{\partial\phi'}{\partial x'^4}\frac{\partial x'^4}{\partial x^4}$$

That in the matrix form and without presenting the function ϕ becomes:

$$\frac{\partial\phi'}{\partial x^j}=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{\partial x^4}\right]=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{\partial x'^4}\right]\begin{bmatrix}\frac{\partial x'^1}{\partial x^1}=1 & \frac{\partial x'^1}{\partial x^2}=0 & \frac{\partial x'^1}{\partial x^3}=0 & \frac{\partial x'^1}{\partial x^4}=-v\\ \frac{\partial x'^2}{\partial x^1}=0 & \frac{\partial x'^2}{\partial x^2}=1 & \frac{\partial x'^2}{\partial x^3}=0 & \frac{\partial x'^2}{\partial x^4}=0\\ \frac{\partial x'^3}{\partial x^1}=0 & \frac{\partial x'^3}{\partial x^2}=0 & \frac{\partial x'^3}{\partial x^3}=1 & \frac{\partial x'^3}{\partial x^4}=0\\ \frac{\partial x'^4}{\partial x^1}=\frac{-v}{c^2\sqrt{K}} & \frac{\partial x'^4}{\partial x^2}=0 & \frac{\partial x'^4}{\partial x^3}=0 & \frac{\partial x'^4}{\partial x^4}=\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux^1}{c^2}\right)\end{bmatrix}$$

Where replacing the items below:

$$\frac{\partial x'^1}{\partial x^4}=-v=\frac{-v'}{\sqrt{K'}}$$

$$\frac{\partial x'^4}{\partial x^1}=\frac{-v}{c^2\sqrt{K}}=\frac{-v'}{c^2}$$

$$\frac{\partial x'^4}{\partial x^4}=\frac{1}{\sqrt{K}}\left(1+\frac{v^2}{c^2}-\frac{vux^1}{c^2}\right)=\frac{\partial x^4}{\partial x'^4}=\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'^1}{c^2}\right)$$

Observation: this last relation shows that the time varies in an equal form between the referentials.

We get:

$$\frac{\partial\phi'}{\partial x^j}=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{\partial x^4}\right]=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{\partial x'^4}\right]\begin{bmatrix}\frac{\partial x'^1}{\partial x^1}=1 & \frac{\partial x'^1}{\partial x^2}=0 & \frac{\partial x'^1}{\partial x^3}=0 & \frac{\partial x'^1}{\partial x^4}=-\frac{v'}{\sqrt{K'}}\\ \frac{\partial x'^2}{\partial x^1}=0 & \frac{\partial x'^2}{\partial x^2}=1 & \frac{\partial x'^2}{\partial x^3}=0 & \frac{\partial x'^2}{\partial x^4}=0\\ \frac{\partial x'^3}{\partial x^1}=0 & \frac{\partial x'^3}{\partial x^2}=0 & \frac{\partial x'^3}{\partial x^3}=1 & \frac{\partial x'^3}{\partial x^4}=0\\ \frac{\partial x'^4}{\partial x^1}=\frac{-v'}{c^2} & \frac{\partial x'^4}{\partial x^2}=0 & \frac{\partial x'^4}{\partial x^3}=0 & \frac{\partial x'^4}{\partial x^4}=\frac{1}{\sqrt{K'}}\left(1+\frac{v'^2}{c^2}+\frac{v'u'x'^1}{c^2}\right)\end{bmatrix}$$

That is the group 8.2 plus 8.4 from the table 9, differential operators in the matrix form.

Applying 8.5 in 8.3 and in 8.4 we simplify these equations in the following way:

Table 9B, differential operators with the equations 8.3 and 8.4 simplified:

$\frac{\partial}{\partial x'^1}=\frac{\partial}{\partial x^1}+\frac{v}{c^2}\frac{\partial}{\partial x^4}$	8.1	$\frac{\partial}{\partial x^1}=\frac{\partial}{\partial x'^1}-\frac{v'}{c^2}\frac{\partial}{\partial x'^4}$	8.2
$\frac{\partial}{\partial x'^2}=\frac{\partial}{\partial x^2}$	8.1.1	$\frac{\partial}{\partial x^2}=\frac{\partial}{\partial x'^2}$	8.2.1
$\frac{\partial}{\partial x'^3}=\frac{\partial}{\partial x^3}$	8.1.2	$\frac{\partial}{\partial x^3}=\frac{\partial}{\partial x'^3}$	8.2.2
$\frac{-\partial}{c\partial x'^4}=\sqrt{K}\frac{-\partial}{c\partial x^4}$	8.3B	$\frac{-\partial}{c\partial x^4}=\sqrt{K'}\frac{-\partial}{c\partial x'^4}$	8.4B
$\frac{\partial}{\partial x^1}+\frac{ux^1}{c^2}\frac{\partial}{\partial x^4}=zero$	8.5	$\frac{\partial}{\partial x'^1}+\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}=zero$	8.5

The table 9B, in the matrix form becomes:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{-\partial}{c\partial x'^4}\right]=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{-\partial}{c\partial x^4}\right]\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ -v/c & 0 & 0 & \sqrt{K}\end{bmatrix} \qquad 15.23$$

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{-\partial}{c\partial x^4}\right]=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{-\partial}{c\partial x'^4}\right]\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ v'/c & 0 & 0 & \sqrt{K'}\end{bmatrix} \qquad 15.24$$

The squared matrices of the transformations above are transposed of the matrices A and A'.

Invariance of the Total Differential

In the observer O referential the total differential of a function $\phi(x^k)$ is equal to:

$$d\phi(x^k)=\frac{\partial\phi}{\partial x^k}dx^k=\frac{\partial\phi}{\partial x^1}dx^1+\frac{\partial\phi}{\partial x^2}dx^2+\frac{\partial\phi}{\partial x^3}dx^3+\frac{\partial\phi}{\partial x^4}dx^4=\left[\frac{\partial\phi}{\partial x^1}\frac{\partial\phi}{\partial x^2}\frac{\partial\phi}{\partial x^3}\frac{\partial\phi}{c\partial x^4}\right]\begin{bmatrix}dx^1\\ dx^2\\ dx^3\\ cdx^4\end{bmatrix} \qquad 15.25$$

Where the coordinates correlate with the ones from the observer O' according to $x^k=x^k(x'^l)$, replacing the transformations 15.24 and 15.18 and without presenting the function ϕ we have:

$$d\phi=\frac{\partial\phi}{\partial x^k}dx^k=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ -v'/c & 0 & 0 & \sqrt{K'}\end{bmatrix}\begin{bmatrix}1 & 0 & 0 & v'/c\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & \sqrt{K'}\end{bmatrix}\begin{bmatrix}dx'^1\\ dx'^2\\ dx'^3\\ cdx'^4\end{bmatrix} \qquad 15.26$$

The multiplication of the middle matrices supplies:

$$\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ -v'/c & 0 & 0 & \sqrt{K'}\end{bmatrix}\begin{bmatrix}1 & 0 & 0 & v'/c\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & \sqrt{K'}\end{bmatrix}=\begin{bmatrix}1 & 0 & 0 & v'/c\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ -v'/c & 0 & 0 & 1+\frac{2v'dx'^1}{c^2dx'^4}\end{bmatrix} \qquad 15.27$$

Result that can be divided in two matrices:

$$\begin{bmatrix}1 & 0 & 0 & v'/c\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ -v'/c & 0 & 0 & 1+\frac{2v'dx'^1}{c^2dx'^4}\end{bmatrix}=\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & 1\end{bmatrix}+\begin{bmatrix}0 & 0 & 0 & v'/c\\ 0 & 0 & 0 & 0\\ 0 & 0 & 0 & 0\\ -v'/c & 0 & 0 & \frac{2v'dx'^1}{c^2dx'^4}\end{bmatrix} \qquad 15.28$$

That applied to the total differential supplies:

$$d\phi = \frac{\partial \phi}{\partial x^k} dx^k = \left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&1\end{bmatrix}+\begin{bmatrix}0&0&0&v'/c\\0&0&0&0\\0&0&0&0\\-v'/c&0&0&\frac{2v'dx'^1}{c^2dx'^4}\end{bmatrix}\right]\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix} \qquad 15.29$$

Executing the operations of the second term we have:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}0&0&0&v'/c\\0&0&0&0\\0&0&0&0\\-v'/c&0&0&\frac{2v'dx'^1}{c^2dx'^4}\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix}=-\frac{v'}{c^2}\frac{\partial}{\partial x'^4}dx'^1+v'\frac{\partial}{\partial x'^1}dx'^4+\frac{2v'}{c^2}\frac{dx'^1}{dx'^4}\frac{\partial}{\partial x'^4}dx'^4$$

Where applying 8.5 we have:

$$-\frac{v'}{c^2}\frac{\partial}{\partial x'^4}dx'^1+v'\left(-\frac{1}{c^2}\frac{dx'^1}{dx'^4}\frac{\partial}{\partial x'^4}\right)dx'^4+\frac{2v'}{c^2}\frac{dx'^1}{dx'^4}\frac{\partial}{\partial x'^4}dx'^4=zero$$

Then we have:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}0&0&0&v'/c\\0&0&0&0\\0&0&0&0\\-v'/c&0&0&\frac{2v'dx'^1}{c^2dx'^4}\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix}=zero \qquad 15.30$$

With this result we have in 15.29 the invariance of the total differential:

$$d\phi = \frac{\partial \phi}{\partial x^k} dx^k = \left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&1\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix}=\frac{\partial \phi'}{\partial x'^l}dx'^l=d\phi' \qquad 15.31$$

In the observer O' referential the total differential of a function $\phi(x'^i)$ is equal to:

$$d\phi'(x'^i)=\frac{\partial \phi'}{\partial x'^i}dx'^i=\frac{\partial \phi'}{\partial x'^1}dx'^1+\frac{\partial \phi'}{\partial x'^2}dx'^2+\frac{\partial \phi'}{\partial x'^3}dx'^3+\frac{\partial \phi'}{\partial x'^4}dx'^4=\left[\frac{\partial \phi'}{\partial x'^1}\frac{\partial \phi'}{\partial x'^2}\frac{\partial \phi'}{\partial x'^3}\frac{\partial \phi'}{c\partial x'^4}\right]\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix} \qquad 15.32$$

Where the coordinates correlate with the ones from the observer O referential according to $x'^i = x'^i(x^j)$, replacing the transformations 15.23 and 15.15 and without presenting the function ϕ we have:

$$d\phi' = \frac{\partial \phi'}{\partial x'^i}dx'^i = \left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\v/c&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}1&0&0&-v/c\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}dx^1\\dx^2\\dx^3\\cdx^4\end{bmatrix} \qquad 15.33$$

The multiplication of the middle matrices supplies:

$$\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\v/c&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}1&0&0&-v/c\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K}\end{bmatrix}=\begin{bmatrix}1&0&0&-v/c\\0&1&0&0\\0&0&1&0\\v/c&0&0&1-\frac{2vdx^1}{c^2dx^4}\end{bmatrix} \qquad 15.34$$

Result that can be divided in two matrices:

$$\begin{bmatrix} 1 & 0\,0 & -v/c \\ 0 & 1\,0 & 0 \\ 0 & 0\,1 & 0 \\ v/c & 0\,0\,1 & -\frac{2vdx^1}{c^2dx^4} \end{bmatrix} = \begin{bmatrix} 1\,0\,0\,0 \\ 0\,1\,0\,0 \\ 0\,0\,1\,0 \\ 0\,0\,0\,1 \end{bmatrix} + \begin{bmatrix} 0 & 0\,0 & -v/c \\ 0 & 0\,0 & 0 \\ 0 & 0\,0 & 0 \\ v/c & 0\,0 & -\frac{2vdx^1}{c^2dx^4} \end{bmatrix} \qquad 15.35$$

That applied to the total differential supplies:

$$d\phi' = \frac{\partial\phi'}{\partial x'^i}dx'^i = \left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\left[\begin{bmatrix} 1\,0\,0\,0 \\ 0\,1\,0\,0 \\ 0\,0\,1\,0 \\ 0\,0\,0\,1 \end{bmatrix} + \begin{bmatrix} 0 & 0\,0 & -v/c \\ 0 & 0\,0 & 0 \\ 0 & 0\,0 & 0 \\ v/c & 0\,0 & -\frac{2vdx^1}{c^2dx^4} \end{bmatrix}\right]\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} \qquad 15.36$$

Executing the operations of the second term we have:

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 0 & 0\,0 & -v/c \\ 0 & 0\,0 & 0 \\ 0 & 0\,0 & 0 \\ v/c & 0\,0 & -\frac{2vdx^1}{c^2dx^4} \end{bmatrix}\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} = \frac{v}{c^2}\frac{\partial}{\partial x^4}dx^1 - v\frac{\partial}{\partial x^1}dx^4 - \frac{2v}{c^2}\frac{dx^1}{dx^4}\frac{\partial}{\partial x^4}dx^4$$

Where applying 8.5 we have:

$$\frac{v}{c^2}\frac{\partial}{\partial x^4}dx^1 - v\left(-\frac{1}{c^2}\frac{dx^1}{dx^4}\frac{\partial}{\partial x^4}\right)dx^4 - \frac{2v}{c^2}\frac{dx^1}{dx^4}\frac{\partial}{\partial x^4}dx^4 = zero$$

Then we have:

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 0 & 0\,0 & -v/c \\ 0 & 0\,0 & 0 \\ 0 & 0\,0 & 0 \\ v/c & 0\,0 & -\frac{2vdx^1}{c^2dx^4} \end{bmatrix}\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} = zero \qquad 15.37$$

With this result we have in 15.36 the invariance of the total differential:

$$d\phi' = \frac{\partial\phi'}{\partial x'^i}dx'^i = \left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 1\,0\,0\,0 \\ 0\,1\,0\,0 \\ 0\,0\,1\,0 \\ 0\,0\,0\,1 \end{bmatrix}\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} = \frac{\partial\phi}{\partial x^j}dx^j = d\phi \qquad 15.38$$

Invariance of the Wave Equation

The wave equation to the observer O is equal to:

$$\nabla^2\phi - \frac{1}{c^2}\frac{\partial\phi^2}{\partial(x^4)^2} = \frac{\partial\phi^2}{\partial(x^1)^2} + \frac{\partial\phi^2}{\partial(x^2)^2} + \frac{\partial\phi^2}{\partial(x^3)^2} - \frac{1}{c^2}\frac{\partial\phi^2}{\partial(x^4)^2} = \left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 1\,0\,0 & 0 \\ 0\,1\,0 & 0 \\ 0\,0\,1 & 0 \\ 0\,0\,0 & -1 \end{bmatrix}\begin{bmatrix} \frac{\partial^1}{\partial x^1} \\ \frac{\partial}{\partial x^2} \\ \frac{\partial}{\partial x^3} \\ \frac{\partial}{c\partial x^4} \end{bmatrix} = 0 \qquad 15.39$$

Where applying 15.24 and the transposed from 15.24 we have:

$$\nabla^2\phi-\frac{1}{c^2}\frac{\partial\phi^2}{\partial\left(x^4\right)^2}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\\frac{-v'}{c}&0&0&\sqrt{K'}\end{bmatrix}\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}1&0&0&\frac{-v'}{c}\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K'}\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x'^1}\\\frac{\partial}{\partial x'^2}\\\frac{\partial}{\partial x'^3}\\\frac{\partial}{c\partial x'^4}\end{bmatrix} \qquad 15.40$$

The multiplication of the three middle matrices supplies:

$$\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\\frac{-v'}{c}&0&0&\sqrt{K'}\end{bmatrix}\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}1&0&0&\frac{-v'}{c}\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K'}\end{bmatrix}=\begin{bmatrix}1&0&0&\frac{-v'}{c}\\0&1&0&0\\0&0&1&0\\\frac{-v'}{c}&0&0&-1\frac{-2v'u'x'^1}{c^2}\end{bmatrix} \qquad 15.41$$

Result that can be divided in two matrices:

$$\begin{bmatrix}1&0&0&\frac{-v'}{c}\\0&1&0&0\\0&0&1&0\\\frac{-v'}{c}&0&0&-1\frac{-2v'u'x'^1}{c^2}\end{bmatrix}=\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}+\begin{bmatrix}0&0&0&\frac{-v'}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v'}{c}&0&0&\frac{-2v'u'x'^1}{c^2}\end{bmatrix} \qquad 15.42$$

That applied in the wave equation supplies:

$$\nabla^2\phi-\frac{1}{c^2}\frac{\partial\phi^2}{\partial\left(x^4\right)^2}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}+\begin{bmatrix}0&0&0&\frac{-v'}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v'}{c}&0&0&\frac{-2v'u'x'^1}{c^2}\end{bmatrix}\right]\begin{bmatrix}\frac{\partial}{\partial x'^1}\\\frac{\partial}{\partial x'^2}\\\frac{\partial}{\partial x'^3}\\\frac{\partial}{c\partial x'^4}\end{bmatrix} \qquad 15.43$$

Executing the operations of the second term we have:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}0&0&0&\frac{-v'}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v'}{c}&0&0&\frac{-2v'u'x'^1}{c^2}\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x'^1}\\\frac{\partial}{\partial x'^2}\\\frac{\partial}{\partial x'^3}\\\frac{\partial}{c\partial x'^4}\end{bmatrix}=-\frac{v'}{c^2}\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^4}-\frac{v'}{c^2}\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^4}-\frac{2v'u'x'^1}{c^2}\frac{\partial^2}{c^2\,\partial\left(x'^4\right)^2}$$

Executing the operations we have:

$$-\frac{2v'}{c^2}\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^4}-\frac{2v'}{c^2}\frac{u'x'^1}{c^2}\frac{\partial^2}{\partial\left(x'^4\right)^2}$$

Where applying 8.5 we have:

$$-\frac{2v'}{c^2}\left(-\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}\right)\frac{\partial}{\partial x'^4}-\frac{2v'}{c^2}\frac{u'x'^1}{c^2}\frac{\partial^2}{\partial\left(x'^4\right)^2}=zero$$

Then we have:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}0 & 0 & 0 & \frac{-v'}{c}\\ 0 & 0 & 0 & 0\\ 0 & 0 & 0 & 0\\ \frac{-v'}{c} & 0 & 0 & \frac{-2v'u'x'^1}{c^2}\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x'^1}\\ \frac{\partial}{\partial x'^2}\\ \frac{\partial}{\partial x'^3}\\ \frac{\partial}{c\partial x'^4}\end{bmatrix}=zero \qquad 15.44$$

With this result we have in 15.43 the invariance of the wave equation:

$$\nabla^2\phi-\frac{1}{c^2}\frac{\partial\phi^2}{\partial\left(x^4\right)^2}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & -1\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x'^1}\\ \frac{\partial}{\partial x'^2}\\ \frac{\partial}{\partial x'^3}\\ \frac{\partial}{c\partial x'^4}\end{bmatrix}=\nabla^2\phi'-\frac{1}{c^2}\frac{\partial\phi'^2}{\partial\left(x'^4\right)^2} \qquad 15.45$$

The wave equation to the observer O' is equal to:

$$\nabla^2\phi'-\frac{1}{c^2}\frac{\partial\phi'^2}{\partial\left(x'^4\right)^2}=\frac{\partial\phi'^2}{\partial\left(x'^1\right)^2}+\frac{\partial\phi'^2}{\partial\left(x'^2\right)^2}+\frac{\partial\phi'^2}{\partial\left(x'^3\right)^2}-\frac{1}{c^2}\frac{\partial\phi'^2}{\partial\left(x'^4\right)^2}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & -1\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x'^1}\\ \frac{\partial}{\partial x'^2}\\ \frac{\partial}{\partial x'^3}\\ \frac{\partial}{c\partial x'^4}\end{bmatrix}=0 \qquad 15.46$$

Where applying 15.23 and the transposed from 15.23 we have:

$$\nabla^2\phi'-\frac{1}{c^2}\frac{\partial\phi'^2}{\partial\left(x'^4\right)^2}=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ \frac{v}{c} & 0 & 0 & \sqrt{K}\end{bmatrix}\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & -1\end{bmatrix}\begin{bmatrix}1 & 0 & 0 & \frac{v}{c}\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & \sqrt{K}\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x^1}\\ \frac{\partial}{\partial x^2}\\ \frac{\partial}{\partial x^3}\\ \frac{\partial}{c\partial x^4}\end{bmatrix} \qquad 15.47$$

The multiplication of the three middle matrices supplies:

$$\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ \frac{v}{c} & 0 & 0 & \sqrt{K}\end{bmatrix}\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & -1\end{bmatrix}\begin{bmatrix}1 & 0 & 0 & \frac{v}{c}\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & \sqrt{K}\end{bmatrix}=\begin{bmatrix}1 & 0 & 0 & \frac{v}{c}\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ \frac{v}{c} & 0 & 0 & -1+\frac{2vux^1}{c^2}\end{bmatrix} \qquad 15.48$$

Result that can be divided in two matrices:

$$\begin{bmatrix}1 & 0 & 0 & \frac{v}{c}\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ \frac{v}{c} & 0 & 0 & -1+\frac{2vux^1}{c^2}\end{bmatrix}=\begin{bmatrix}1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & -1\end{bmatrix}+\begin{bmatrix}0 & 0 & 0 & \frac{v}{c}\\ 0 & 0 & 0 & 0\\ 0 & 0 & 0 & 0\\ \frac{v}{c} & 0 & 0 & \frac{2vux^1}{c^2}\end{bmatrix} \qquad 15.49$$

That applied in the wave equation supplies:

$$\nabla^2\phi' - \frac{1}{c^2}\frac{\partial \phi'^2}{\partial (x'^4)^2} = \left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}+\begin{bmatrix}0&0&0&\frac{v}{c}\\0&0&0&0\\0&0&0&0\\\frac{v}{c}&0&0&\frac{2vux^1}{c^2}\end{bmatrix}\right]\begin{bmatrix}\frac{\partial}{\partial x^1}\\\frac{\partial}{\partial x^2}\\\frac{\partial}{\partial x^3}\\\frac{\partial}{c\partial x^4}\end{bmatrix} \quad 15.50$$

Executing the operations of the second term we have:

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix}0&0&0&\frac{v}{c}\\0&0&0&0\\0&0&0&0\\\frac{v}{c}&0&0&\frac{2vux^1}{c^2}\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x^1}\\\frac{\partial}{\partial x^2}\\\frac{\partial}{\partial x^3}\\\frac{\partial}{\partial x^4}\end{bmatrix} = \frac{v}{c^2}\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^4} + \frac{v}{c^2}\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^4} + \frac{2v}{c^2}\frac{ux^1}{c^2}\frac{\partial^2}{\partial (x^4)^2}$$

Executing the operations we have:

$$\frac{2v}{c^2}\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^4} + \frac{2v}{c^2}\frac{ux^1}{c^2}\frac{\partial^2}{\partial (x^4)^2}$$

Where applying 8.5 we have:

$$\frac{2v}{c^2}\left(\frac{-ux^1}{c^2}\frac{\partial}{\partial x^4}\right)\frac{\partial}{\partial x^4} + \frac{2v}{c^2}\frac{ux^1}{c^2}\frac{\partial^2}{\partial (x^4)^2} = zero$$

Then we have:

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix}0&0&0&\frac{v}{c}\\0&0&0&0\\0&0&0&0\\\frac{v}{c}&0&0&\frac{2vux^1}{c^2}\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x^1}\\\frac{\partial}{\partial x^2}\\\frac{\partial}{\partial x^3}\\\frac{\partial}{\partial x^4}\end{bmatrix} = zero \quad 15.51$$

Then in 15.50 we have the invariance of the wave equation:

$$\nabla^2\phi' - \frac{1}{c^2}\frac{\partial \phi'^2}{\partial (x'^4)^2} = \left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}\frac{\partial}{\partial x^1}\\\frac{\partial}{\partial x^2}\\\frac{\partial}{\partial x^3}\\\frac{\partial}{c\partial x^4}\end{bmatrix} = \nabla^2\phi - \frac{1}{c^2}\frac{\partial \phi^2}{\partial (x^4)^2} \quad 15.52$$

Invariance of the equations 8.5 of linear propagation

Replacing 2.4, 8.2, 8.4B in 8.5 we have:

$$\frac{\partial}{\partial x^1} + \frac{ux^1}{c^2}\frac{\partial}{\partial x^4} = \frac{\partial}{\partial x'^1} - \frac{v'}{c^2}\frac{\partial}{\partial x'^4} + \frac{1}{c^2}\frac{(u'x'^1+v')}{\sqrt{K'}}\sqrt{K'}\frac{\partial}{\partial x'^4} = zero$$

Executing the operations we have:

$$\frac{\partial}{\partial x^1} + \frac{ux^1}{c^2}\frac{\partial}{\partial x^4} = \frac{\partial}{\partial x'^1} - \frac{v'}{c^2}\frac{\partial}{\partial x'^4} + \frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4} + \frac{v'}{c^2}\frac{\partial}{\partial x'^4} = zero$$

That simplified supplies the invariance of the equation 8.5:

$$\frac{\partial}{\partial x^1}+\frac{ux^1}{c^2}\frac{\partial}{\partial x^4}=\frac{\partial}{\partial x'^1}+\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}=zero$$

Replacing 2.3, 8.1, 8.3B in 8.5 we have:

$$\frac{\partial}{\partial x'^1}+\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}=\frac{\partial}{\partial x^1}+\frac{v}{c^2}\frac{\partial}{\partial x^4}+\frac{1}{c^2}\frac{(ux^1-v)}{\sqrt{K}}\sqrt{K}\frac{\partial}{\partial x^4}=zero$$

Executing the operations we have:

$$\frac{\partial}{\partial x'^1}+\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}=\frac{\partial}{\partial x^1}+\frac{v}{c^2}\frac{\partial}{\partial x^4}+\frac{ux^1}{c^2}\frac{\partial}{\partial x^4}-\frac{v}{c^2}\frac{\partial}{\partial x^4}=zero$$

That simplified supplies the invariance of the equation 8.5:

$$\frac{\partial}{\partial x'^1}+\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}=\frac{\partial}{\partial x^1}+\frac{ux^1}{c^2}\frac{\partial}{\partial x^4}=zero$$

The table 4 in a matrix from becomes:

$$\begin{bmatrix} px'^1 \\ px'^2 \\ px'^3 \\ E'/c \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix}\begin{bmatrix} px^1 \\ px^2 \\ px^3 \\ E/c \end{bmatrix} \quad 15.53$$

$$\begin{bmatrix} px^1 \\ px^2 \\ px^3 \\ E/c \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & v'/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} px'^1 \\ px'^2 \\ px'^3 \\ E'/c \end{bmatrix} \quad 15.54$$

The table 6 in a matrix form becomes:

$$\begin{bmatrix} J'x'^1 \\ J'x'^2 \\ J'x'^3 \\ c\rho' \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & -v/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K} \end{bmatrix}\begin{bmatrix} Jx^1 \\ Jx^2 \\ Jx^3 \\ c\rho \end{bmatrix} \quad 15.55$$

$$\begin{bmatrix} Jx^1 \\ Jx^2 \\ Jx^3 \\ c\rho \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & v'/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} J'x'^1 \\ J'x'^2 \\ J'x'^3 \\ c\rho' \end{bmatrix} \quad 15.56$$

Invariance of the Continuity Equation

The continuity equation to the observer O is equal to:

$$\vec{\nabla}.\vec{J}+\frac{\partial\rho}{\partial x^4}=\frac{\partial Jx^1}{\partial x^1}+\frac{\partial Jx^2}{\partial x^2}+\frac{\partial Jx^3}{\partial x^3}+\frac{\partial\rho}{\partial x^4}=\begin{bmatrix} \frac{\partial}{\partial x^1} & \frac{\partial}{\partial x^2} & \frac{\partial}{\partial x^3} & \frac{\partial}{c\partial x^4} \end{bmatrix}\begin{bmatrix} Jx^1 \\ Jx^2 \\ Jx^3 \\ c\rho \end{bmatrix}=zero \quad 15.57$$

Where replacing 15.24 and 15.56 we have:

$$\vec{\nabla}.\vec{J}+\frac{\partial\rho}{\partial x^4}=\begin{bmatrix} \frac{\partial}{\partial x'^1} & \frac{\partial}{\partial x'^2} & \frac{\partial}{\partial x'^3} & \frac{\partial}{c\partial x'^4} \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -v'/c & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & v'/c \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} J'x'^1 \\ J'x'^2 \\ J'x'^3 \\ c\rho' \end{bmatrix}=zero \quad 15.58$$

The product of the transformation matrices is given in 15.27 and 15.28 with this:

$$\vec{\nabla}.\vec{J}+\frac{\partial\rho}{\partial x^4}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&1\end{bmatrix}+\begin{bmatrix}0&0&0&v'/c\\0&0&0&0\\0&0&0&0\\-v'/c&0&0&\frac{2v'u'x'^1}{c^2}\end{bmatrix}\right]\begin{bmatrix}J'x'^1\\J'x'^2\\J'x'^3\\c\rho'\end{bmatrix} \qquad 15.59$$

Executing the operations of the second term we have:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}0&0&0&v'/c\\0&0&0&0\\0&0&0&0\\-v'/c&0&0&\frac{2v'u'x'^1}{c^2}\end{bmatrix}\begin{bmatrix}J'x'^1\\J'x'^2\\J'x'^3\\c\rho'\end{bmatrix}=-\frac{v'}{c^2}\frac{\partial Jx'^1}{\partial x'^4}+\frac{v'\partial\rho'}{\partial x'^1}+\frac{2v'u'x'^1}{c^2}\frac{\partial\rho'}{\partial x'^4}$$

Where replacing $Jx'^1=\rho'u'x'^1$ and 8.5 we have:

$$-\frac{v'u'x'^1}{c^2}\frac{\partial\rho'}{\partial x'^4}+v'\left(-\frac{u'x'^1}{c^2}\frac{\partial}{\partial x'^4}\right)\rho'+\frac{2v'u'x'^1}{c^2}\frac{\partial\rho'}{\partial x'^4}=zero$$

Then we have:

$$\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}0&0&0&v'/c\\0&0&0&0\\0&0&0&0\\-v'/c&0&0&\frac{2v'u'x'^1}{c^2}\end{bmatrix}\begin{bmatrix}J'x'^1\\J'x'^2\\J'x'^3\\c\rho'\end{bmatrix}=zero \qquad 15.60$$

With this result we have in 15.59 the invariance of the continuity equation:

$$\vec{\nabla}.\vec{J}+\frac{\partial\rho}{\partial x^4}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&1\end{bmatrix}\begin{bmatrix}J'x'^1\\J'x'^2\\J'x'^3\\c\rho'\end{bmatrix}=\vec{\nabla}.\vec{J}'+\frac{\partial\rho'}{\partial x'^4} \qquad 15.61$$

The continuity equation to the observer O' is equal to:

$$\vec{\nabla}.\vec{J}'+\frac{\partial\rho'}{\partial x'^4}=\frac{\partial J'x'^1}{\partial x'^1}+\frac{\partial J'x'^2}{\partial x'^2}+\frac{\partial J'x'^3}{\partial x'^3}+\frac{\partial\rho'}{\partial x'^4}=\left[\frac{\partial}{\partial x'^1}\frac{\partial}{\partial x'^2}\frac{\partial}{\partial x'^3}\frac{\partial}{c\partial x'^4}\right]\begin{bmatrix}J'x'^1\\J'x'^2\\J'x'^3\\c\rho'\end{bmatrix}=zero \qquad 15.62$$

Where replacing 15.23 and 15.55 we have:

$$\vec{\nabla}.\vec{J}'+\frac{\partial\rho'}{\partial x'^4}=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\v/c&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}1&0&0&-v/c\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}Jx^1\\Jx^2\\Jx^3\\c\rho\end{bmatrix}=zero \qquad 15.63$$

The product of the transformation matrices is given in 15.34 and 15.35 then we have:

$$\vec{\nabla}.\vec{J}'+\frac{\partial\rho'}{\partial x'^4}=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&1\end{bmatrix}+\begin{bmatrix}0&0&0&-v/c\\0&0&0&0\\0&0&0&0\\v/c&0&0&-\frac{2vux^1}{c^2}\end{bmatrix}\right]\begin{bmatrix}Jx^1\\Jx^2\\Jx^3\\c\rho\end{bmatrix} \qquad 15.64$$

Executing the operations of the second term we have:

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 0 & 0 & 0 & -v/c \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ v/c & 0 & 0 & -\frac{2vux^1}{c^2} \end{bmatrix}\begin{bmatrix} Jx^1 \\ Jx^2 \\ Jx^3 \\ c\rho \end{bmatrix}=\frac{v}{c^2}\frac{\partial Jx^1}{\partial x^4}-\frac{v\partial\rho}{\partial x^1}-\frac{2vux^1}{c^2}\frac{\partial\rho}{\partial x^4}$$

Where replacing $Jx^1=\rho ux^1$ and 8.5 we have:

$$\frac{vux^1}{c^2}\frac{\partial\rho}{\partial x^4}-v\left(-\frac{ux^1}{c^2}\frac{\partial}{\partial x^4}\right)\rho-\frac{2vux^1}{c^2}\frac{\partial\rho}{\partial x^4}=zero$$

Then we have:

$$\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 0 & 0 & 0 & -v/c \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ v/c & 0 & 0 & -\frac{2vux^1}{c^2} \end{bmatrix}\begin{bmatrix} Jx^1 \\ Jx^2 \\ Jx^3 \\ c\rho \end{bmatrix}=zero \qquad 15.65$$

With this result we have in 15.64 the invariance of the continuity equation:

$$\vec{\nabla}.\vec{J}'+\frac{\partial\rho'}{\partial x'^4}=\left[\frac{\partial}{\partial x^1}\frac{\partial}{\partial x^2}\frac{\partial}{\partial x^3}\frac{\partial}{c\partial x^4}\right]\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}\begin{bmatrix} Jx^1 \\ Jx^2 \\ Jx^3 \\ c\rho \end{bmatrix}=\vec{\nabla}.\vec{J}+\frac{\partial\rho}{\partial x^4} \qquad 15.66$$

Invariance of the line differential element:

That to the observer O is written this way:

$$(ds)^2=\left(dx^1\right)^2+\left(dx^2\right)^2+\left(dx^3\right)^2-\left(cdx^4\right)^2=\left[dx^1\ dx^2\ dx^3\ cdx^4\right]\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}\begin{bmatrix} dx^1 \\ dx^2 \\ dx^3 \\ cdx^4 \end{bmatrix} \qquad 15.67$$

Where replacing 15.18 and the transposed from 15.18 we have:

$$(ds)^2=\left[dx'^1\ dx'^2\ dx'^3\ cdx'^4\right]\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \frac{v'}{c} & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & \frac{v'}{c} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} dx'^1 \\ dx'^2 \\ dx'^3 \\ cdx'^4 \end{bmatrix} \qquad 15.68$$

The multiplication of the three central matrices supplies:

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \frac{v'}{c} & 0 & 0 & \sqrt{K'} \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}\begin{bmatrix} 1 & 0 & 0 & \frac{v'}{c} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sqrt{K'} \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & \frac{v'}{c} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \frac{v'}{c} & 0 & 0 & -1\frac{-2v'dx'^1}{c^2dx'^4} \end{bmatrix} \qquad 15.69$$

Result that can be divided in two matrices:

$$\begin{bmatrix} 1 & 0 & 0 & \frac{v'}{c} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ \frac{v'}{c} & 0 & 0 & -1\frac{-2v'dx'^1}{c^2dx'^4} \end{bmatrix}=\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}+\begin{bmatrix} 0 & 0 & 0 & \frac{v'}{c} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ \frac{v'}{c} & 0 & 0 & \frac{-2v'dx'^1}{c^2dx'^4} \end{bmatrix} \qquad 15.70$$

That applied in the line differential element supplies:

$$(ds)^2=\left[dx'^1\ dx'^2\ dx'^3\ cdx'^4\right]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}+\begin{bmatrix}0&0&0&\frac{v'}{c}\\0&0&0&0\\0&0&0&0\\\frac{v'}{c}&0&0&\frac{-2v'dx'^1}{c^2dx'^4}\end{bmatrix}\right]\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix} \qquad 15.71$$

Executing the operations of the second term we have:

$$\left[dx'^1\ dx'^2\ dx'^3\ cdx'^4\right]\begin{bmatrix}0&0&0&\frac{v'}{c}\\0&0&0&0\\0&0&0&0\\\frac{v'}{c}&0&0&\frac{-2v'}{c^2}\frac{dx'^1}{dx'^4}\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix}=\frac{v'dx'^1\,cdx'^4}{c}+cdx'^4\left(\frac{v'}{c}dx'^1-\frac{2v'}{c^2}\frac{dx'^1}{dx'^4}cdx'^4\right)=zero$$

Then we have:

$$\left[dx'^1\ dx'^2\ dx'^3\ cdx'^4\right]\begin{bmatrix}0&0&0&\frac{v'}{c}\\0&0&0&0\\0&0&0&0\\\frac{v'}{c}&0&0&\frac{-2v''}{c^2}\frac{dx'^1}{dx'^4}\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix}=zero \qquad 15.72$$

With this result we have in 15.71 the invariance of the line differential element:

$$(ds)^2=\left[dx'^1\ dx'^2\ dx'^3\ cdx'^4\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix}=(dx'^1)^2+(dx'^2)^2+(dx'^3)^2-(cdx'^4)^2=(ds')^2 \qquad 15.73$$

To the observer O' the line differential element is written this way:

$$(ds')^2=(dx'^1)^2+(dx'^2)^2+(dx'^3)^2-(cdx'^4)^2=\left[dx'^1\ dx'^2\ dx'^3\ cdx'^4\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}dx'^1\\dx'^2\\dx'^3\\cdx'^4\end{bmatrix} \qquad 15.74$$

Where replacing 15.15 and the transposed from 15.15 we have:

$$(ds')^2=\left[dx^1dx^2dx^3cdx^4\right]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\\frac{-v}{c}&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}1&0&0&\frac{-v}{c}\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}dx^1\\dx^2\\dx^3\\cdx^4\end{bmatrix} \qquad 15.75$$

The multiplication of the three central matrices supplies:

$$\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\\frac{-v}{c}&0&0&\sqrt{K}\end{bmatrix}\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}1&0&0&\frac{-v}{c}\\0&1&0&0\\0&0&1&0\\0&0&0&\sqrt{K}\end{bmatrix}=\begin{bmatrix}1&0&0&\frac{-v}{c}\\0&1&0&0\\0&0&1&0\\\frac{-v}{c}&0&0&-1+\frac{2vdx^1}{c^2dx^4}\end{bmatrix} \qquad 15.76$$

Result that can be divided in two matrices:

$$\begin{bmatrix}1&0&0&\frac{-v}{c}\\0&1&0&0\\0&0&1&0\\\frac{-v}{c}&0&0&-1+\frac{2vdx^1}{c^2dx^4}\end{bmatrix}=\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}+\begin{bmatrix}0&0&0&\frac{-v}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v}{c}&0&0&\frac{2vdx^1}{c^2dx^4}\end{bmatrix} \qquad 15.77$$

That applied in the line differential element supplies:

$$(ds')^2 = [dx^1 dx^2 dx^3 cdx^4]\left[\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}+\begin{bmatrix}0&0&0&\frac{-v}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v}{c}&0&0&\frac{2vdx^1}{c^2dx^4}\end{bmatrix}\right]\begin{bmatrix}dx^1\\dx^2\\dx^3\\cdx^4\end{bmatrix} \qquad 15.78$$

Executing the operations of the second term we have:

$$[dx^1 dx^2 dx^3 cdx^4]\begin{bmatrix}0&0&0&\frac{-v}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v}{c}&0&0&\frac{2v}{c^2}\frac{dx^1}{dx^4}\end{bmatrix}\begin{bmatrix}dx^1\\dx^2\\dx^3\\cdx^4\end{bmatrix}=\frac{-vdx^1cdx^4}{c}+cdx^4\left(\frac{-v}{c}dx^1+\frac{2v}{c^2}\frac{dx^1}{dx^4}cdx^4\right)=zero$$

Then we have:

$$[dx^1 dx^2 dx^3 cdx^4]\begin{bmatrix}0&0&0&\frac{-v}{c}\\0&0&0&0\\0&0&0&0\\\frac{-v}{c}&0&0&\frac{2v}{c^2}\frac{dx^1}{dx^4}\end{bmatrix}\begin{bmatrix}dx^1\\dx^2\\dx^3\\cdx^4\end{bmatrix}=zero \qquad 15.79$$

With this result we have in 15.78 the invariance of the line differential element:

$$(ds')^2 = [dx^1 dx^2 dx^3 cdx^4]\begin{bmatrix}1&0&0&0\\0&1&0&0\\0&0&1&0\\0&0&0&-1\end{bmatrix}\begin{bmatrix}dx^1\\dx^2\\dx^3\\cdx^4\end{bmatrix}=(dx^1)^2+(dx^2)^2+(dx^3)^2-(cdx^4)^2=(ds)^2 \qquad 15.80$$

In §7 as a consequence of 5.3 we had the invariance of $\vec{E}.\vec{u}=\vec{E}'.\vec{u}'$ where now applying 7.3.1, 7.3.2, 7.4.1, 7.4.2 and the velocity transformation formulae from table 2 we have new relations between Ex and $E'x'$ distinct from 7.3 and 7.4 and with them we rewrite the table 7 in the form below:

Table 7B

$E'x'=\dfrac{Ex\sqrt{K}}{\left(1-\dfrac{v}{ux}\right)}$	7.3B	$Ex=\dfrac{E'x'\sqrt{K'}}{\left(1+\dfrac{v'}{u'x'}\right)}$	7.4B
$E'y'=Ey\sqrt{K}$	7.3.1	$Ey=E'y'\sqrt{K'}$	7.4.1
$E'z'=Ez\sqrt{K}$	7.3.2	$Ez=E'z'\sqrt{K'}$	7.4.2
$B'x'=Bx$	7.5	$Bx=B'x'$	7.6
$B'y'=By+\dfrac{v}{c^2}Ez$	7.5.1	$By=B'y'-\dfrac{v'}{c^2}E'z'$	7.6.1
$B'z'=Bz-\dfrac{v}{c^2}Ey$	7.5.2	$Bz=B'z'+\dfrac{v'}{c^2}E'y'$	7.6.2
$By=-\dfrac{ux}{c^2}Ez$	7.9	$B'y'=-\dfrac{u'x'}{c^2}E'z'$	7.10
$Bz=\dfrac{ux}{c^2}Ey$	7.9.1	$B'z'=\dfrac{u'x'}{c^2}E'y'$	7.10.1
$\left(1-\dfrac{v}{ux}\right)\left(1+\dfrac{v'}{u'x'}\right)=1$			

With the tables 7B and 9B we can have the invariance of all Maxwell's equations.

Invariance of the Gauss' Law for the electrical field:

$$\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z'}{\partial z'}=\frac{\rho'}{\varepsilon_0} \quad 8.14$$

Where applying the tables 6, 7B and 9B we have:

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)\frac{Ex\sqrt{K}}{(1-v/ux)}+\frac{\partial Ey\sqrt{K}}{\partial y}+\frac{Ez\sqrt{K}}{\partial z}=\frac{\rho\sqrt{K}}{\varepsilon_0}$$

Where simplifying and replacing 8.5 we have:

$$\left[\frac{\partial}{\partial x}+v\left(\frac{-1}{ux}\frac{\partial}{\partial x}\right)\right]\frac{Ex}{(1-v/ux)}+\frac{\partial Ey}{\partial y}+\frac{Ez}{\partial z}=\frac{\rho}{\varepsilon_0}$$

That reordered supplies:

$$\left[\frac{\partial}{\partial x}\left(1-\frac{v}{ux}\right)\right]\frac{Ex}{(1-v/ux)}+\frac{\partial Ey}{\partial y}+\frac{Ez}{\partial z}=\frac{\rho}{\varepsilon_0}.$$

That simplified supplies the invariance of the Gauss' Law for the electrical field.

Invariance of the Gauss' Law for the magnetic field:

$$\frac{\partial B'x'}{\partial x'}+\frac{\partial B'y'}{\partial y'}+\frac{\partial B'z'}{\partial z'}=zero \quad 8.16$$

Where applying the tables 7B and 9B we have:

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)Bx+\frac{\partial}{\partial y}\left(By+\frac{v}{c^2}Ez\right)+\frac{\partial}{\partial z}\left(Bz-\frac{v}{c^2}Ey\right)=0$$

That reordered supplies:

$$\frac{\partial Bx}{\partial x}+\frac{\partial By}{\partial y}+\frac{\partial Bz}{\partial z}+\frac{v}{c^2}\left(\frac{\partial Ez}{\partial y}-\frac{\partial Ey}{\partial z}+\frac{\partial Bx}{\partial t}\right)=0$$

Where the term in parenthesis is the Faraday-Henry's Law (8.19) that is equal to zero from where we have the invariance of the Gauss' Law for the magnetic field.

Invariance of the Faraday-Henry's Law:

$$\frac{\partial E'y'}{\partial x'}-\frac{\partial E'x'}{\partial y'}=-\frac{\partial B'z'}{\partial t'} \quad 8.18$$

Where applying the tables 7B and 9B we have:

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)Ey\sqrt{K}-\frac{\partial}{\partial y}\frac{Ex\sqrt{K}}{(1-v/ux)}=-\sqrt{K}\frac{\partial}{\partial t}\left(Bz-\frac{v}{c^2}Ey\right)$$

That simplified and multiplied by $(1-v/ux)$ we have:

$$\frac{\partial Ey}{\partial x}\left(1-\frac{v}{ux}\right)-\frac{\partial Ex}{\partial y}=-\frac{\partial Bz}{\partial t}\left(1-\frac{v}{ux}\right)$$

Where executing the products and replacing 7.9.1 we have:

$$\frac{\partial Ey}{\partial x}-\frac{\partial Ex}{\partial y}=-\frac{\partial Bz}{\partial t}+\frac{v}{ux}\left(\frac{\partial Ey}{\partial x}+\frac{ux}{c^2}\frac{\partial Ey}{\partial t}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Faraday-Henry's Law.

Invariance of the Faraday-Henry's Law:

$$\frac{\partial E'z'}{\partial y'}-\frac{\partial E'y'}{\partial z'}=-\frac{\partial B'x'}{\partial t'} \qquad 8.20$$

Where applying the tables 7B and 9B we have:

$$\frac{\partial Ez}{\partial y}\sqrt{K}-\frac{\partial Ey}{\partial z}\sqrt{K}=-\sqrt{K}\frac{\partial Bx}{\partial t}$$

That simplified supplies the invariance of the Faraday-Henry's Law.

Invariance of the Faraday-Henry's Law:

$$\frac{\partial E'x'}{\partial z'}-\frac{\partial E'z'}{\partial x'}=-\frac{\partial B'y'}{\partial t'} \qquad 8.22$$

Where applying the tables 7B and 9B we have:

$$\frac{\partial}{\partial z}\frac{Ex\sqrt{K}}{(1-v/ux)}-\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)Ez\sqrt{K}=-\sqrt{K}\frac{\partial}{\partial t}\left(By+\frac{v}{c^2}Ez\right)$$

That simplified and multiplied by $(1-v/ux)$ we have:

$$\frac{\partial Ex}{\partial z}-\frac{\partial Ez}{\partial x}\left(1-\frac{v}{ux}\right)-\frac{v}{c^2}\frac{\partial Ez}{\partial t}\left(1-\frac{v}{ux}\right)=-\frac{\partial By}{\partial t}\left(1-\frac{v}{ux}\right)-\frac{v}{c^2}\frac{\partial Ez}{\partial t}\left(1-\frac{v}{ux}\right)$$

That simplifying and making the operations we have:

$$\frac{\partial Ex}{\partial z}-\frac{\partial Ez}{\partial x}=-\frac{\partial By}{\partial t}-\frac{v}{ux}\left(\frac{\partial Ez}{\partial x}-\frac{\partial By}{\partial t}\right)$$

Where applying 7.9 we have:

$$\frac{\partial Ex}{\partial z}-\frac{\partial Ez}{\partial x}=-\frac{\partial By}{\partial t}-\frac{v}{ux}\left(\frac{\partial Ez}{\partial x}+\frac{ux}{c^2}\frac{\partial Ez}{\partial t}\right).$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Faraday-Henry's Law.

Invariance of the Ampere-Maxwell's Law:

$$\frac{\partial B'y'}{\partial x'}-\frac{\partial B'x'}{\partial y'}=\mu_o J'z'+\varepsilon_o\mu_o\frac{\partial E'z'}{\partial t'} \qquad 8.24$$

Where applying the tables 6, 7B and 9B we have:

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)\left(By+\frac{v}{c^2}Ez\right)-\frac{\partial Bx}{\partial y}=\mu_0 Jz+\varepsilon_0\mu_0\sqrt{K}\frac{\partial}{\partial t}Ez\sqrt{K}$$

That simplifying and making the operations we have:

$$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\mu_0 Jz+\varepsilon_0\mu_0\frac{\partial Ez}{\partial t}+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ez}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ez}{\partial t}-\frac{v}{c^2}\frac{\partial Ez}{\partial x}-\frac{v}{c^2}\frac{\partial By}{\partial t}-\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ez}{\partial t}$$

Where simplifying and applying 7.9 we have:

$$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\mu_0 Jz+\varepsilon_0\mu_0\frac{\partial Ez}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ez}{\partial t}-\frac{v}{c^2}\frac{\partial Ez}{\partial x}-\frac{v}{c^2}\left(\frac{-ux}{c^2}\frac{\partial Ez}{\partial t}\right)$$

That reorganized supplies

$$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\mu_0 Jz+\varepsilon_0\mu_0\frac{\partial Ez}{\partial t}-\frac{v}{c^2}\left(\frac{ux}{c^2}\frac{\partial Ez}{\partial t}+\frac{\partial Ez}{\partial x}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Ampere-Maxwell's Law:

Invariance of the Ampere-Maxwell's Law:

$$\frac{\partial B'z'}{\partial y'}-\frac{\partial B'y'}{\partial z'}=\mu_o J'x'+\varepsilon_o\mu_o\frac{\partial E'x'}{\partial t'} \qquad 8.26$$

Where applying the tables 6, 7B and 9B we have:

$$\frac{\partial}{\partial y}\left(Bz-\frac{v}{c^2}Ey\right)-\frac{\partial}{\partial z}\left(By+\frac{v}{c^2}Ez\right)=\mu_0\left(Jx-\rho v\right)+\varepsilon_0\mu_0\sqrt{K}\frac{\partial}{\partial t}\frac{Ex\sqrt{K}}{(1-v/ux)}$$

Making the operations we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\frac{v}{c^2}\left(\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}-\mu_0 c^2\rho\right)+\varepsilon_0\mu_0\left(1+\frac{v^2}{c^2}-\frac{2vux}{c^2}\right)\frac{\partial Ex}{\partial t}\frac{1}{(1-v/ux)}$$

Replacing in the first parenthesis the Gauss' Law and multiplying by $\left(1-\frac{v}{ux}\right)$ we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}+\frac{v}{ux}\left(\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}-\mu_0 Jx\right)-\frac{v}{c^2}\frac{\partial Ex}{\partial x}+\frac{v^2}{c^2}\left(\frac{1}{ux}\frac{\partial Ex}{\partial x}\right)+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ex}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

Where replacing $Jx=\rho ux$, 7.9.1, 7.9 and 8.5 we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}+\frac{v}{ux}\left(\frac{ux}{c^2}\frac{\partial Ey}{\partial y}+\frac{ux}{c^2}\frac{\partial Ez}{\partial z}-\mu_0\rho ux\right)-\frac{v}{c^2}\frac{\partial Ex}{\partial x}+\frac{v^2}{c^2}\left(\frac{-1}{c^2}\frac{\partial Ex}{\partial t}\right)+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ex}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

That simplified supplies:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}+\frac{v}{c^2}\left(\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}-\mu_0 c^2\rho\right)-\frac{v}{c^2}\frac{\partial Ex}{\partial x}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

Replacing in the first parenthesis the Gauss' Law we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}-\frac{v}{c^2}\frac{\partial Ex}{\partial x}-\frac{v}{c^2}\frac{\partial Ex}{\partial x}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

That reorganized makes:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}-\frac{2v}{c^2}\left(\frac{\partial Ex}{\partial x}+\frac{ux}{c^2}\frac{\partial Ex}{\partial t}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Ampere-Maxwell's Law:

Invariance of the Ampere-Maxwell's Law:

$$\frac{\partial B'x'}{\partial z'}-\frac{\partial B'z'}{\partial x'}=\mu_o J'y'+\varepsilon_o\mu_o\frac{\partial E'y'}{\partial t'} \qquad 8.28$$

Where applying the tables 6, 7B and 9B we have:

$$\frac{\partial Bx}{\partial z}-\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)\left(Bz-\frac{v}{c^2}Ey\right)=\mu_0 Jy+\varepsilon_0\mu_0\sqrt{K}\frac{\partial}{\partial t}Ey\sqrt{K}$$

Making the operations we have:

$$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\mu_0 Jy+\varepsilon_0\mu_0\frac{\partial Ey}{\partial t}+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ey}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ey}{\partial t}-\frac{v}{c^2}\frac{\partial Ey}{\partial x}+\frac{v}{c^2}\frac{\partial Bz}{\partial t}-\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ey}{\partial t}$$

Where simplifying and applying 7.9.1 we have:

$$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\mu_0 Jy+\varepsilon_0\mu_0\frac{\partial Ey}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ey}{\partial t}-\frac{v}{c^2}\frac{\partial Ey}{\partial x}+\frac{v}{c^2}\left(\frac{ux}{c^2}\frac{\partial Ey}{\partial t}\right)$$

That reorganized makes:

$$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\mu_0 Jy+\varepsilon_0\mu_0\frac{\partial Ey}{\partial t}-\frac{v}{c^2}\left(\frac{ux}{c^2}\frac{\partial Ey}{\partial t}+\frac{\partial Ey}{\partial x}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Ampere-Maxwell's Law:

Invariance of the Gauss' Law for the electrical field without electrical charge:

$$\frac{\partial E'x'}{\partial x'}+\frac{\partial E'y'}{\partial y'}+\frac{\partial E'z'}{\partial z'}=zero \qquad 8.30$$

Where applying the tables 7B and 9B we have:

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)\frac{Ex\sqrt{K}}{(1-v/ux)}+\frac{\partial Ey\sqrt{K}}{\partial y}+\frac{Ez\sqrt{K}}{\partial z}=zero$$

Where simplifying and replacing 8.5 we have:

$$\left[\frac{\partial}{\partial x}+v\left(\frac{-1}{ux}\frac{\partial}{\partial x}\right)\right]\frac{Ex}{(1-v/ux)}+\frac{\partial Ey}{\partial y}+\frac{Ez}{\partial z}=zero$$

That reorganized makes:

$$\left[\frac{\partial}{\partial x}\left(1-\frac{v}{ux}\right)\right]\frac{Ex}{(1-v/ux)}+\frac{\partial Ey}{\partial y}+\frac{Ez}{\partial z}=zero\,.$$

That simplified supplies the Gauss' Law for the electrical field without electrical charge.

Invariance of the Ampere-Maxwell's Law without electrical charge:

$$\frac{\partial B'y'}{\partial x'}-\frac{\partial B'x'}{\partial y'}=\varepsilon_o\mu_o\frac{\partial E'z'}{\partial t'} \qquad 8.40$$

Where applying the tables 7B and 9B we have:

$$\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)\left(By+\frac{v}{c^2}Ez\right)-\frac{\partial Bx}{\partial y}=\varepsilon_0\mu_0\sqrt{K}\frac{\partial}{\partial t}Ez\sqrt{K}$$

Making the operations we have:

$$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\varepsilon_0\mu_0\frac{\partial Ez}{\partial t}+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ez}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ez}{\partial t}-\frac{v}{c^2}\frac{\partial Ez}{\partial x}-\frac{v}{c^2}\frac{\partial By}{\partial t}-\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ez}{\partial t}$$

Where simplifying and applying 7.9 we have:

$$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\varepsilon_0\mu_0\frac{\partial Ez}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ez}{\partial t}-\frac{v}{c^2}\frac{\partial Ez}{\partial x}-\frac{v}{c^2}\left(\frac{-ux}{c^2}\frac{\partial Ez}{\partial t}\right)$$

That reorganized makes:

$$\frac{\partial By}{\partial x}-\frac{\partial Bx}{\partial y}=\varepsilon_0\mu_0\frac{\partial Ez}{\partial t}-\frac{v}{c^2}\left(\frac{ux}{c^2}\frac{\partial Ez}{\partial t}+\frac{\partial Ez}{\partial x}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Ampere-Maxwell's Law without electrical charge:

Invariance of the Ampere-Maxwell's Law without electrical charge:

$$\frac{\partial B'z'}{\partial y'}-\frac{\partial B'y'}{\partial z'}=\varepsilon_o\mu_o\frac{\partial E'x'}{\partial t'} \qquad 8.42$$

Where applying the tables 7B and 9B we have:

$$\frac{\partial}{\partial y}\left(Bz-\frac{v}{c^2}Ey\right)-\frac{\partial}{\partial z}\left(By+\frac{v}{c^2}Ez\right)=\varepsilon_0\mu_0\sqrt{K}\frac{\partial}{\partial t}\frac{Ex\sqrt{K}}{(1-v/ux)}$$

Making the operations we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\frac{v}{c^2}\left(\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}\right)+\varepsilon_0\mu_0\left(1+\frac{v^2}{c^2}-\frac{2vux}{c^2}\right)\frac{\partial Ex}{\partial t}\frac{1}{(1-v/ux)}$$

Replacing in the first parenthesis the Gauss' Law without electrical charge and multiplying by $(1-v/ux)$ we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}+\frac{v}{ux}\left(\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}\right)-\frac{v}{c^2}\frac{\partial Ex}{\partial x}+\frac{v^2}{c^2}\left(\frac{1}{ux}\frac{\partial Ex}{\partial x}\right)+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ex}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

Where replacing 7.9, 7.9.1 and 8.5 we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}+\frac{v}{ux}\left(\frac{ux}{c^2}\frac{\partial Ey}{\partial y}+\frac{ux}{c^2}\frac{\partial Ez}{\partial z}\right)-\frac{v}{c^2}\frac{\partial Ex}{\partial x}+\frac{v^2}{c^2}\left(\frac{-1}{c^2}\frac{\partial Ex}{\partial t}\right)+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ex}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

That simplified supplies:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}+\frac{v}{c^2}\left(\frac{\partial Ey}{\partial y}+\frac{\partial Ez}{\partial z}\right)-\frac{v}{c^2}\frac{\partial Ex}{\partial x}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

Replacing in the first parenthesis the Gauss' Law without electrical charge we have:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}-\frac{v}{c^2}\frac{\partial Ex}{\partial x}-\frac{v}{c^2}\frac{\partial Ex}{\partial x}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ex}{\partial t}$$

That reorganized makes:

$$\frac{\partial Bz}{\partial y}-\frac{\partial By}{\partial z}=\mu_0 Jx+\varepsilon_0\mu_0\frac{\partial Ex}{\partial t}-\frac{2v}{c^2}\left(\frac{\partial Ex}{\partial x}+\frac{ux}{c^2}\frac{\partial Ex}{\partial t}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Ampere-Maxwell's Law without electrical charge:

Invariance of the Ampere-Maxwell's Law without electrical charge:

$$\frac{\partial B'x'}{\partial z'}-\frac{\partial B'z'}{\partial x'}=\varepsilon_o\mu_o\frac{\partial E'y'}{\partial t'} \qquad 8.44$$

Where applying the tables 6, 7B and 9B we have:

$$\frac{\partial Bx}{\partial z}-\left(\frac{\partial}{\partial x}+\frac{v}{c^2}\frac{\partial}{\partial t}\right)\left(Bz-\frac{v}{c^2}Ey\right)=\varepsilon_0\mu_0\sqrt{K}\frac{\partial}{\partial t}Ey\sqrt{K}$$

Making the operations we have:

$$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\varepsilon_0\mu_0\frac{\partial Ey}{\partial t}+\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ey}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ey}{\partial t}-\frac{v}{c^2}\frac{\partial Ey}{\partial x}+\frac{v}{c^2}\frac{\partial Bz}{\partial t}-\frac{1}{c^2}\frac{v^2}{c^2}\frac{\partial Ey}{\partial t}$$

Where simplifying and applying 7.9.1 we have:

$$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\varepsilon_0\mu_0\frac{\partial Ey}{\partial t}-\frac{1}{c^2}\frac{2vux}{c^2}\frac{\partial Ey}{\partial t}-\frac{v}{c^2}\frac{\partial Ey}{\partial x}+\frac{v}{c^2}\left(\frac{ux}{c^2}\frac{\partial Ey}{\partial t}\right)$$

That reorganized makes:

$$\frac{\partial Bx}{\partial z}-\frac{\partial Bz}{\partial x}=\varepsilon_0\mu_0\frac{\partial Ey}{\partial t}-\frac{v}{c^2}\left(\frac{ux}{c^2}\frac{\partial Ey}{\partial t}+\frac{\partial Ey}{\partial x}\right)$$

As the term in parenthesis is the equation 8.5 that is equal to zero then we have the invariance of the Ampere-Maxwell's Law without electrical charge:

§15 Invariance (continuation)

A function $f(\theta)=\mathrm{f}(\mathrm{kr}-wt)$ 2.19

Where the phase is equal to $\theta=(kr-wt)$ 15.81

In order to represent an undulating movement that goes on in one arbitrary direction must comply with the wave equation and because of this we have:

$$\frac{k}{r^2}\left[3r-\frac{(x^2+y^2+z^2)}{r}\right]\frac{\partial f(\theta)}{\partial\theta}+\frac{k^2}{r^2}(x^2+y^2+z^2)\frac{\partial^2 f(\theta)}{\partial\theta^2}-k^2\frac{\partial^2 f(\theta)}{\partial\theta^2}=zero \qquad 15.82$$

That doesn't meet with the wave equation because the two last elements get nule but the first one doesn't.

In order to overcome this problem we reformulate the phase θ of the function in the following way.

A unitary vector such as

$$\vec{n} = cos\,\phi\vec{i} + cos\,\alpha\vec{j} + cos\,\beta\vec{k} \qquad 15.83$$

where $cos\phi = \frac{x}{r} = \frac{x}{ct}$, $cos\alpha = \frac{y}{r} = \frac{y}{ct}$, $cos\beta = \frac{z}{r} = \frac{z}{ct}$ 15.84

has the module equal to $n = |\vec{n}| = \sqrt{\vec{n}.\vec{n}} = \sqrt{cos^2\,\phi + cos^2\,\alpha + cos^2\,\beta} = 1$. 15.85

Making the product

$$\vec{n}.\vec{R} = \left(cos\phi\vec{i} + cos\alpha\vec{j} + cos\beta\vec{k}\right).\left(x\vec{i} + y\vec{j} + z\vec{k}\right) = cos\phi x + cos\alpha y + cos\beta z = \frac{x^2 + y^2 + z^2}{r} = \frac{r^2}{r} = r \qquad 15.86$$

we have $r = \vec{n}.\vec{R} = cos\,\phi x + cos\,\alpha y + cos\,\beta z$ that applied to the phase θ supplies a new phase

$$\Phi = (kr - wt) = \left(k\vec{n}.\vec{R} - wt\right) = \left(k\,cos\,\phi x + k\,cos\,\alpha y + k\,cos\,\beta z - wt\right) \qquad 15.87$$

with the same meaning of the previous phase $\theta = \Phi$.

Replacing $r = \vec{n}.\vec{R} = cos\,\phi x + cos\,\alpha y + cos\,\beta z$ e $k = \frac{w}{c}$ in the phase θ multiplied by –1 we also get another phase in the form

$$\Phi = (-1)(kr - wt) = (wt - kr) = \left[w\left(t - \frac{r}{c}\right)\right] = \left[w\left(t - \frac{cos\,\phi x + cos\,\alpha y + cos\,\beta z}{c}\right)\right] \qquad 15.88$$

with the same meaning of the previous phase $(-1)\theta = \Phi$.

Thus we can write a new function as:

$$f(\Phi) = \mathrm{f}\left[w\left(t - \frac{cos\phi x + cos\alpha y + cos\beta z}{c}\right)\right] \qquad 15.89$$

That replaced in the wave equation with the director cosine considered constant supplies:

$$\frac{\partial^2 f(\Phi)}{\partial\Phi^2}\frac{w^2}{c^2}cos^2\,\phi + \frac{\partial^2 f(\Phi)}{\partial\Phi^2}\frac{w^2}{c^2}cos^2\,\alpha + \frac{\partial^2 f(\Phi)}{\partial\Phi^2}\frac{w^2}{c^2}cos^2\,\beta - \frac{\partial^2 f(\Phi)}{\partial\Phi^2}\frac{w^2}{c^2} = zero \qquad 15.90$$

that simplified meets the wave equation.

The positive result of the phase Φ in the wave equation is an exclusive consequence of the director cosines being constant in the partial derivatives showing that the wave equation demands the propagation to have one steady direction in the space (plane wave).

For the observer O a source located in the origin of its referential produces in a random point located at the distance $r = ct = \sqrt{x^2 + y^2 + z^2}$ of the origin, an electrical field $\vec{E}$ described by:

$$\vec{E} = Ex\vec{i} + Ey\vec{j} + Ez\vec{k} \qquad 15.91$$

Where the components are described as:

$$\begin{aligned} Ex &= E_{xo}.f(\Phi) \\ Ey &= E_{yo}.f(\Phi) \\ Ez &= E_{zo}.f(\Phi) \end{aligned} \qquad 15.92$$

That applied in $\vec{E}$ supplies:

$$\vec{E} = E_{xo}f(\Phi)\vec{i} + E_{yo}f(\Phi)\vec{j} + E_{zo}f(\Phi)\vec{k} = \left[E_{xo}\vec{i} + E_{yo}\vec{j} + E_{zo}\right]f(\Phi) = \vec{E}_o f(\Phi). \qquad 15.93$$

with module equal to $E = \sqrt{(E_{xo})^2 + (E_{yo})^2 + (E_{zo})^2}.f(\Phi) \Rightarrow E = E_o.f(\Phi)$ 15.94

Being $\vec{E}_o = E_{xo}\vec{i} + E_{yo}\vec{j} + E_{zo}\vec{k}$ 15.95

The maximum amplitude vector Constant with the components E_{xo}, E_{yo}, E_{zo} 15.96

And module $E_o = \sqrt{(E_{xo})^2 + (E_{yo})^2 + (E_{zo})^2}$ 15.97

Being $f(\Phi)$ a function with the phase Φ equal to15.87 or 15.88.

Deriving the component E_x in relation to x and t we have:

$$\frac{\partial Ex}{\partial x} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{\partial \Phi}{\partial x} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{\partial (kr - wt)}{\partial x} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{kx}{r} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{kx}{ct} \qquad 15.98$$

$$\frac{\partial Ex}{\partial t} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{\partial \Phi}{\partial t} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{\partial (kr - wt)}{\partial t} = E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}(-w) \qquad 15.99$$

that applied in 8.5 supplies

$$\frac{\partial Ex}{\partial x} + \frac{x/t}{c^2}\frac{\partial Ex}{\partial t} = zero \Rightarrow E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{\partial \Phi}{\partial x} + \frac{x/t}{c^2}E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\frac{\partial \Phi}{\partial t} = zero \Rightarrow E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\left(\frac{\partial \Phi}{\partial x} + \frac{x/t}{c^2}\frac{\partial \Phi}{\partial t}\right) = zero$$

$$E_{xo}\frac{\partial f(\Phi)}{\partial \Phi}\left(\frac{\partial \Phi}{\partial x} + \frac{x/t}{c^2}\frac{\partial \Phi}{\partial t}\right) = zero \Rightarrow \frac{\partial \Phi}{\partial x} + \frac{x/t}{c^2}\frac{\partial \Phi}{\partial t} = zero \qquad 15.100$$

demonstrating that it is the phase Φ that must comply with 8.5.

$$\frac{\partial \Phi}{\partial x} + \frac{x/t}{c^2}\frac{\partial \Phi}{\partial t} = zero \Rightarrow \frac{\partial (kr - wt)}{\partial x} + \frac{x/t}{c^2}\frac{\partial (kr - wt)}{\partial t} = zero \Rightarrow \frac{kx}{ct} + \frac{x/t}{c^2}(-w) = zero \Rightarrow \frac{x}{ct}\left(k - \frac{w}{c}\right) = zero$$

as $k = \frac{w}{c}$ then E_x complies with 8.5.

As the phase is the same for the components E_y and E_z then they also comply with 8.5.

As the phases for the observers O and O' are equal $(kr - wt) = (k'r' - w't')$ then the components of the observer O' also comply with 8.5.

$$\frac{\partial (kr - wt)}{\partial x} + \frac{x/t}{c^2}\frac{\partial (kr - wt)}{\partial t} = \frac{\partial (k'r' - w't')}{\partial x'} + \frac{x'/t'}{c^2}\frac{\partial (k'r' - w't')}{\partial t'} = zero \qquad 15.101$$

The components relatively to the observer O of the electrical field are transformed for the referential of the observer O' according to the tables 7, 7B and 8.

Applying in 8.5 a wave function written in the form:

$$\Psi = e^{i(kx-wt)} = e^{i\Phi} = \cos(kx-wt)+i\sin(kx-wt) = \cos\Phi + i\sin\Phi \qquad 15.102$$

where $i=\sqrt{-1}$.

Deriving we have:

$$\frac{\partial\Psi}{\partial x} = -k\,sen\,\Phi + ki\,cos\,\Phi \text{ end } \frac{\partial\Psi}{\partial t} = w\,sen\,\Phi - wi\,cos\,\Phi \qquad 15.103$$

or $\frac{\partial\Psi}{\partial x} = ke^{i\Phi}$ and $\frac{\partial\Psi}{\partial t} = -we^{i\Phi}$ 15.104

That applied in 8.5 supplies:

$$\frac{\partial\Psi}{\partial x} + \frac{x/t}{c^2}\frac{\partial\Psi}{\partial t} = zero \Rightarrow (-k\,sen\,\Phi + ki\,cos\,\Phi) + \frac{x/t}{c^2}(w\,sen\,\Phi - wi\,cos\,\Phi) = zero$$

that is equal to:

$$\left(-k+\frac{xw}{c^2t}\right)\sin\Phi + \left(ki - \frac{xwi}{c^2t}\right)\cos\Phi = zero$$

or $\frac{\partial\Psi}{\partial x} + \frac{x/t}{c^2}\frac{\partial\Psi}{\partial t} = zero \Rightarrow \left(ke^{i\Phi}\right) + \frac{x/t}{c^2}\left(-we^{i\Phi}\right) = zero$

where we must have the coefficients equal to zero so that we get na identity, then:

$$-k+\frac{xw}{c^2t} = zero \Rightarrow k = \frac{xw}{c^2t}$$

$$ki-\frac{xwi}{c^2t} = zero \Rightarrow k = \frac{xw}{c^2t}$$

$$\left(ke^{i\Phi}\right) + \frac{x/t}{c^2}\left(-we^{i\Phi}\right) = zero \Rightarrow k = \frac{xw}{c^2t}$$

Where applying $w = ck$ we have:

$$k = \frac{xw}{c^2t} = \frac{xck}{c^2t} \Rightarrow \frac{x}{t} = c$$

Then to meet with the equation 8.5 we must have a wave propagation along the axis x with the speed c.

If we apply $w = uk$ and $v = \frac{x}{t}$ we have:

$$k = \frac{xw}{c^2t} = \frac{vuk}{c^2} \Rightarrow u = \frac{c^2}{v}.$$

A result also gotten from the Louis de Broglie's wave equation.

§16 Time and Frequency

Considering the Doppler effect as a law of physics.

We can define a clock as any device that produces a frequency of identical events in a series possible to be enlisted and added in such a way that a random event *n* of a device will be identical to any event in the series of events produced by a replica of this device when the events are compared in a relative resting position.

The cyclical movement of a clock in a resting position according to the observer O referential sets the time in this referential and the cyclical movement of the arms of a clock in a resting position according to the observer O' sets the time in this referential. The formulas of time transformation 1.7 and 1.8 relate the times between the referentials in relative movement thus, relate movements in relative movement.

The relative movement between the inertial referentials produces the Doppler effect that proves that the frequency varies with velocity and as the frequency can be interpreted as being the frequency of the cyclical movement of the arms of a clock then the time varies in the same proportion that varies the frequency with the relative movement that is, it is enough to replace the time t and t' in the formulas 1.7 and 1.8 by the frequencies y and y' to get the formulas of frequency transformation, then:

$t' = t\sqrt{K} \Rightarrow y' = y\sqrt{K}$ 1.7 becomes 2.22

$t = t'\sqrt{K'} \Rightarrow y = y'\sqrt{K'}$ 1.8 becomes 2.22

The Galileo's transformation of velocities $\vec{u}' = \vec{u} - \vec{v}$ between two inertial referentials presents intrinsically three defects that can be described this way:

a) The Galileo's transformation of velocity to the axis x is $u'x' = ux - v$. In that one if we have $ux = c$ then $u'x' = c - v$ and if we have $u'x' = c$ then $ux = c + v$. As both results are not simultaneously possible or else we have $ux = c$ or $u'x' = c$ then the transformation doesn't allow that a ray of light be simultaneously observed by the observers O and O' what shows the privilege of an observer in relation to the other because each observer can only see the ray of light running in its own referential (intrinsic defect to the classic analysis of the Sagnac's effect).

b) It cannot also comply to Newton's first law of inertia because a ray of light emitted parallel to the axis x from the origin of the respective inertial referentials at the moment that the origins are coincident and at the moment in which t = t' = zero will have by the Galileo's transformation the velocity c of light altered by $\pm v$ to the referentials, on the contrary of the inertial law that wouldn't allow the existence of a variation in velocity because there is no external action acting on the ray of light and because of this both observers should see the ray of light with velocity c.

c) As it considers the time as a constant between the referentials it doesn't produce the temporal variation between the referentials in movement as it is required by the Doppler effect.

The principle of constancy of light velocity is nothing but a requirement of the Newton's first law, the inertia law.

Newton's first law, the inertia law, is introduced in Galileo's transformation when the principle of constancy of light velocity is applied in Galileo's transformation providing the equation of tables 1 and 2 of the Undulating Relativity that doesn't have the three defects described.

The time and velocity equations of tables 1 and 2 can be written as:

$$t' = t\sqrt{1 + \frac{v^2}{c^2} - \frac{2v}{c}\cos\phi} \qquad 1.7$$

$$v' = \frac{v}{\sqrt{1 + \frac{v^2}{c^2} - \frac{2v}{c}\cos\phi}} \qquad 1.15$$

$$t = t'\sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'}{c}\cos\phi'}$$ 1.8

$$v = \frac{v'}{\sqrt{1 + \frac{v'^2}{c^2} + \frac{2v'}{c}\cos\phi'}}$$ 1.20

The distance d between the referentials is equal to the product of velocity by time this way:

$$d = vt = v'\,t'$$ 1.9

It doesn't depend on the propagation angle of the ray of light, being exclusively a function of velocity and time, that is, the propagation angle of the ray of light, only alters between the inertial referential the proportion between time and velocity, keeping the distance constant in each moment, to any propagation angle.

The equations above in a function form are written as:

$$d = e(v,t) = e'(v',t')$$ 1.9

$$t' = f(v,t,\phi)$$ 1.7

$$v' = g(v,\phi)$$ 1.15

$$t = f'(v',t',\phi')$$ 1.8

$$v = g'(v',\phi')$$ 1.20

Then we have that the distance is a function of two variables, the time a function of three variables and the velocity a function of two variables.

From the definition of moment 4.1 and energy 4.6 we have:

$$\vec{p} = \frac{E}{c^2}\vec{u}$$ 16.1

The elevated to the power of two supplies:

$$\frac{u^2}{c^2} = \frac{c^2}{E^2}p^2$$ 16.2

Elevating to the power of two the energy formula we have:

$$E^2 = \left(\frac{m_0 c^2}{\sqrt{1 - \frac{u^2}{c^2}}}\right)^2 \Rightarrow E^2 - E^2\frac{u^2}{c^2} = m_0^2 c^4$$

Where applying 16.2 we have:

$$E^2 - E^2\frac{u^2}{c^2} = m_0^2 c^4 \Rightarrow E^2 - E^2\frac{c^2}{E^2}p^2 = m_0^2 c^4 \Rightarrow E = c\sqrt{p^2 + m_0^2 c^2}$$ 4.8

From where we conclude that if the mass in resting position of a particle is null $m_o = zero$ the particle energy is equal to $E = c\,p$. 16.3

That applied in 16.2 supplies:

$$\frac{u^2}{c^2}=\frac{c^2}{E^2}p^2 \Rightarrow \frac{u^2}{c^2}=\frac{c^2}{(cp)^2}p^2 \Rightarrow u=c \qquad 16.4$$

From where we conclude that the movement of a particle with a null mass in resting position $m_o = zero$ will always be at the velocity of light $u = c$.

Applying in $E = c\,p$ the relations $E = yh$ and $c = y\lambda$ we have:

$$yh = y\lambda p \Rightarrow p = \frac{h}{\lambda} \text{ and in the same way } p' = \frac{h}{\lambda'} \qquad 16.5$$

Equation that relates the moment of a particle with a null mass in resting position with its own way length.

Elevating to the power of two the formula of moment transformation (4.9) we have:

$$\vec{p}' = \vec{p} - \frac{E}{c^2}\vec{v} \Rightarrow p'^2 = p^2 + \frac{E^2}{c^4}v^2 - 2\frac{E}{c^2}vpx$$

Where applying $E = c\,p$ and $px = p\cos\phi = p\frac{ux}{c}$ we find:

$$p'^2 = p^2 + \frac{(cp)^2}{c^4}v^2 - 2\frac{cp}{c^2}vp\frac{ux}{c} \Rightarrow p' = p\sqrt{1 + \frac{v^2}{c^2} - \frac{2vux}{c^2}} \Rightarrow p' = p\sqrt{K} \qquad 16.6$$

Where applying 16.5 results in:

$$p' = p\sqrt{K} \Rightarrow \frac{h}{\lambda'} = \frac{h}{\lambda}\sqrt{K} \Rightarrow \lambda' = \frac{\lambda}{\sqrt{K}} \text{ or inverted } \lambda = \frac{\lambda'}{\sqrt{K'}} \qquad 2.21$$

Where applying $c = y\lambda$ and $c = y'\lambda'$ we have:

$$y' = y\sqrt{K} \text{ or inverted } y = y'\sqrt{K'} \qquad 2.22$$

In § 2 we have the equations 2.21 and 2.22 applying the principle of relativity to the wave phase.

17 Transformation of H. Lorentz

For two observers in a relative movement, the equation that represents the principle of constancy of light speed for a random point A is:

$$x'^2 + y'^2 + z'^2 - c^2t'^2 = x^2 + y^2 + z^2 - c^2t^2 \qquad 17.01$$

In this equation canceling the symmetric terms we have:
Nesta cancelando os termos simétricos obtemos:

$$x'^2 - c^2t'^2 = x^2 - c^2t^2 \qquad 17.02$$

That we can write as:

$$(x' - ct')(x' + ct') = (x - ct)(x + ct) \qquad 17.03$$

If in this equation we define the proportion factors η and μ as:

$$\begin{cases} (x' - ct') = \eta(x - ct) & A \\ (x' + ct') = \mu(x + ct) & B \end{cases} \qquad 17.04$$

where we must have $\eta.\mu=1$ to comply 17.03.

The equations 17.04 where first gotten by Albert Einstein.

When a ray of light moves in the plane y'z' to the observer O' we have x' = zero and x = vt and such conditions applied to the equation 17.02 supplies:

$$0-c^2t'^2=(vt)^2-c^2t^2 \Rightarrow t'=t\sqrt{1-\frac{v^2}{c^2}} \qquad 17.05$$

This result will also be supplied by the equations A and B of the group 17.04 under the same conditions:

$$\begin{cases} \left(0-ct\sqrt{1-\frac{v^2}{c^2}}\right)=\eta(vt-ct) & A \\ \left(0+ct\sqrt{1-\frac{v^2}{c^2}}\right)=\mu(vt+ct) & B \end{cases} \qquad 17.06$$

From those we have:

$$\eta=\sqrt{\frac{1+\frac{v}{c}}{1-\frac{v}{c}}} \text{ and } \mu=\sqrt{\frac{1-\frac{v}{c}}{1+\frac{v}{c}}} \qquad 17.07$$

Where we have proven that $\eta.\mu=1$.

From the group 17.04 we have the Transformations of H. Lorentz:

$$x'=\frac{(\eta+\mu)}{2}x+\frac{(\mu-\eta)}{2}ct \qquad 17.08$$

$$ct'=\frac{(\mu-\eta)}{2}x+\frac{(\eta+\mu)}{2}ct \qquad 17.09$$

$$x=\frac{(\eta+\mu)}{2}x'+\frac{(\eta-\mu)}{2}ct' \qquad 17.10$$

$$ct=\frac{(\eta-\mu)}{2}x'+\frac{(\eta+\mu)}{2}ct' \qquad 17.11$$

Indexes equations $\frac{\eta+\mu}{2}$, $\frac{\mu-\eta}{2}$ and $\frac{\eta-\mu}{2}$:

$$\eta+\mu=\sqrt{\frac{1+\frac{v}{c}}{1-\frac{v}{c}}}+\sqrt{\frac{1-\frac{v}{c}}{1+\frac{v}{c}}}=\frac{1+\frac{v}{c}+1-\frac{v}{c}}{\sqrt{1-\frac{v}{c}}\sqrt{1+\frac{v}{c}}}=\frac{2}{\sqrt{1-\frac{v^2}{c^2}}}\Rightarrow\frac{\eta+\mu}{2}=\frac{1}{\sqrt{1-\frac{v^2}{c^2}}} \qquad 17.12$$

$$\mu-\eta=\sqrt{\frac{1-\frac{v}{c}}{1+\frac{v}{c}}}-\sqrt{\frac{1+\frac{v}{c}}{1-\frac{v}{c}}}=\frac{1-\frac{v}{c}-1-\frac{v}{c}}{\sqrt{1+\frac{v}{c}}\sqrt{1-\frac{v}{c}}}=\frac{-2\frac{v}{c}}{\sqrt{1-\frac{v^2}{c^2}}}\Rightarrow\frac{\mu-\eta}{2}=\frac{-\frac{v}{c}}{\sqrt{1-\frac{v^2}{c^2}}} \qquad 17.13$$

$$\eta-\mu=\sqrt{\frac{1+\frac{v}{c}}{1-\frac{v}{c}}}-\sqrt{\frac{1-\frac{v}{c}}{1+\frac{v}{c}}}=\frac{1+\frac{v}{c}-1+\frac{v}{c}}{\sqrt{1-\frac{v}{c}}\sqrt{1+\frac{v}{c}}}=\frac{2\frac{v}{c}}{\sqrt{1-\frac{v^2}{c^2}}}\Rightarrow\frac{\eta-\mu}{2}=\frac{\frac{v}{c}}{\sqrt{1-\frac{v^2}{c^2}}} \qquad 17.14$$

Sagnac effect

When both observers' origins are equal the time is zeroed (t = t' = zero) in both referentials and two rays of light are emitted from the common origin, one in the positive direction (clockwise index c) of the axis x and x' with a wave front A_c and another in the negative direction (counter-clockwise index u) of the axis x and x' with a wave front A_u.

The propagation conditions above applied to the Lorentz equations supply the tables A and B below:

Table A

Equation	Clockwise ray (c)	Equation	Counter-clockwise ray (u)	Sum of rays
	Result		Result	
Condition	$x_c = ct_c$	Condition	$x_u = -ct_u$	
17.08	$x'_c = \mu ct_c$	17.08	$x'_u = -\eta ct_u$	
	$x'_c = \mu x_c$		$x'_u = \eta x_u$	$x'_c + x'_u = \mu x_c + \eta x_u$
17.09	$ct'_c = \mu ct_c$	17.09	$ct'_u = \eta ct_u$	$ct'_c + ct'_u = \mu ct_c + \eta ct_u$
	$x'_c = ct'_c$		$x'_u = -ct'_u$	

Table B

Equation	Clockwise ray (c)	Equation	Counter-clockwise ray (u)	Sum of rays
	Result		Result	
Condition	$x'_c = ct'_c$	Condition	$x'_u = -ct'_u$	
17.10	$x_c = \eta ct'_c$	17.10	$x_u = -\mu ct'_u$	
	$x_c = \eta x'_c$		$x_u = \mu x'_u$	$x_c + x_u = \eta x'_c + \mu x'_u$
17.11	$ct_c = \eta ct'_c$	17.11	$ct_u = \mu ct'_u$	$ct_c + ct_u = \eta ct'_c + \mu ct'_u$
	$x_c = ct_c$		$x_u = -ct_u$	

We observe that the tables A and B are inverse one to another.

When we form the group of the sum equations of the two rays from tables A and B:

$$\begin{cases} D' = ct'_c + ct'_u = \mu ct_c + \eta ct_u & A \\ D = ct_c + ct_u = \eta ct'_c + \mu ct'_u & B \end{cases} \quad 17.15$$

Where to the observer O' $D' = A_u \leftrightarrow A_c$ is the distance between the front waves A_u and A_c and where to the observer O $D = A_u \leftrightarrow A_c$ is the distance between the front waves A_u and A_c.

In the equations 17.15 above, due to the isotropy of space and time and the front waves $A_u \leftrightarrow A_c$ of the two rays of light being the same for both observers, the sum of rays of light e times must be invariable between the observers, which we can express by:

$$D' = D \Rightarrow ct'_c + ct'_u = ct_c + ct_u \Rightarrow \sum t' = \sum t \quad 17.16$$

This result that generates an equation of isotropy of space and time can be called as the conservation of space and time principle.

The three hypothesis of propagation defined as follows will be applied in 17.15 and tested to prove the conservation of space and time principle given by 17.16:

Hypothesis A:

If the space and time are isotropic and there is no movement with no privilege of one observer considered over the other in an empty space then the propagation geometry of rays of light can be given by:

$|ct_c| = |ct'_u|$ and $|ct_u| = |ct'_c|$ 17.17

This hypothesis applied to the equation A or B of the group 17.15 complies to the space and time conservation principle given by 17.16.

The hypothesis 17.17 applied to the tables A and B results in:

$$\begin{aligned}&Quadro\ A \begin{cases} ct'_c = \mu ct'_u & A \\ ct'_u = \eta ct'_c & B \end{cases} \\ &Quadro\ B \begin{cases} ct_c = \eta ct_u & C \\ ct_u = \mu ct_c & D \end{cases}\end{aligned} \tag{17.18}$$

Hypothesis B:

If the space and time are isotropic but the observer O is in an absolute resting position in an empty space then the geometry of propagation of the rays of light is given by:

$|ct_c| = |ct_u| = |ct|$ 17.19

That applied to the table A and B results in:

$$\begin{aligned}&Quadro\ A \begin{cases} ct'_c = \mu ct & A \\ ct'_u = \eta ct & B \end{cases} \\ &Quadro\ B \begin{cases} ct = \eta ct'_c & C \\ ct = \mu ct'_u & D \end{cases}\end{aligned} \tag{17.20}$$

$$\begin{cases} ct'_c = \mu^2 ct'_u & A \\ ct'_u = \eta^2 ct'_c & B \end{cases} \tag{17.21}$$

Summing A and B in 17.20 we have:

$$ct'_c + ct'_u = 2ct\left(\frac{\eta+\mu}{2}\right) \Rightarrow D' = D\left(\frac{\eta+\mu}{2}\right) \Rightarrow D' = \frac{D}{\sqrt{1-\frac{v^2}{c^2}}} \Rightarrow \sum t' = \frac{\sum t}{\sqrt{1-\frac{v^2}{c^2}}} \tag{17.22}$$

This result doesn't comply with the conservation of space and time principle given by 17.16 and as $D' \neq D$ it results in a situation of four rays of light, two to each observer, and each ray of light with its respective independent front wave from the others.

Hypothesis C:

If the space and time are isotropic but the observer O' is in an absolute resting position in an empty space then the propagation geometry of the rays of light is given:

$|ct'_c| = |ct'_u| = |ct'|$ 17.23

That applied to the tables A and B results in:

$$\begin{aligned}&Quadro\ A \begin{cases} ct' = \mu ct_c & A \\ ct' = \eta ct_u & B \end{cases} \\ &Quadro\ B \begin{cases} ct_c = \eta ct' & C \\ ct_u = \mu ct' & D \end{cases}\end{aligned} \tag{17.24}$$

$$\begin{cases} ct_c = \eta^2 ct_u & A \\ ct_u = \mu^2 ct_c & B \end{cases} \tag{17.25}$$

Summing C and D in 17.24 we have:

$$ct_c + ct_u = 2ct'\left(\frac{\eta+\mu}{2}\right) \Rightarrow D = D'\left(\frac{\eta+\mu}{2}\right) \Rightarrow D = \frac{D'}{\sqrt{1-\frac{v^2}{c^2}}} \Rightarrow \sum t = \frac{\sum t'}{\sqrt{1-\frac{v^2}{c^2}}} \qquad 17.26$$

This result doesn't comply with the conservation of space and time principle exactly the same way as hypothesis B given by 17.16 and as $D' \neq D \;\; D' \neq D$ it results in a situation of four rays of light, two to each observer and each ray of light with its respective independent front wave from the others.

Conclusion

The hypothesis A, B and C are completely compatible with the demand of isotropy of space and time as we can conclude with the geometry of propagations.

The result of hypothesis A is contrary to the result of hypothesis B and C despite of the relative movement of the observers not changing the front wave A_u relatively to the front wave A_c because the front waves have independent movement one from the other and from the observers.

The hypothesis A applied in the transformations of H. Lorentz complies with the conservation of space and time principle given by 17.16 showing the compatibility with the transformations of H. Lorentz with the hypothesis A. The application of hypothesis B and C in the transformations of H. Lorentz supplies the space and time deformations given by 17.22 and 17.26 because the transformations of H. Lorentz are not compatible with the hypothesis B and C.

For us to obtain the Sagnac effect we must consider that the observer O' is in an absolute resting position, hypothesis C above and that the path of the rays of light be of $2\pi R$:

$$ct'_c = ct'_u = ct' = 2\pi R \qquad 17.27$$

For the observer O the Sagnac effect is given by the time difference between the clockwise ray of light and the counter-clock ray of light $\Delta t = t_c - t_u$ that can be obtained using 17.24 (C-D), 17.27 and 17.14:

$$\Delta t = t_c - t_u = t'(\eta-\mu) = \frac{2\pi R}{c}\left(\frac{2\frac{v}{c}}{\sqrt{1-\frac{v^2}{c^2}}}\right) = \frac{4\pi R v}{c\sqrt{c^2-v^2}} \qquad 17.28$$

§9 The Sagnac Effect (continuation)

The moment the origins are the same the time is zeroed (t = t' = zero) at both sides of the referential and the rays of light are emitted from the common origin, one in the positive way (clockwise index c) of the axis x and x' with a wave front A_c and the other one in the negative way (counter clockwise index u) of the axis x and x' with wave front A_u.

The projected ray of light in the positive way (clockwise index c) of the axis x and x' is equationed by $x_c = ct_c$ and $x'_c = ct'_c$ that applied to the Table I supplies:

$$ct'_c = ct_c\left(1-\frac{v_c}{c}\right) \Rightarrow ct'_c = ct_c K_c \text{ (1.7) } \quad ct_c = ct'_c\left(1+\frac{v'_c}{c}\right) \Rightarrow ct_c = ct'_c K'_c \text{ (1.8)} \qquad 9.11$$

$$v'_c = \frac{v_c}{\left(1-\frac{v_c}{c}\right)} \Rightarrow v'_c = \frac{v_c}{K_c} \text{ (1.15)} \qquad v_c = \frac{v'_c}{\left(1+\frac{v'_c}{c}\right)} \Rightarrow v_c = \frac{v'_c}{K'_c} \text{ (1.20)} \qquad 9.12$$

From those we deduct that the distance between the observers is given by:

$$d_c = v_c t_c = v'_c t'_c \qquad 9.13$$

Where we have:

$$\left(1-\frac{v_c}{c}\right)\left(1+\frac{v'_c}{c}\right)=K_c K'_c=1 \qquad 9.14$$

The ray of light project in the negative way (counter clockwise index u) of the axis x and x' is equationed by $x_u=-ct_u$ and $x'_u=-ct'_u$: that applied to the Table I gives:

$$ct'_u=ct_u\left(1+\frac{v_u}{c}\right)\Rightarrow ct'_u=ct_uK_u \text{ (1.7)} \qquad ct_u=ct'_u\left(1-\frac{v'_u}{c}\right)\Rightarrow ct_u=ct'_uK'_u \text{ (1.8)} \qquad 9.15$$

$$v'_u=\frac{v_u}{\left(1+\frac{v_u}{c}\right)}\Rightarrow v'_u=\frac{v_u}{K_u} \text{ (1.15)} \qquad v_u=\frac{v'_u}{\left(1-\frac{v'_u}{c}\right)}\Rightarrow v_u=\frac{v'_u}{K'_u} \text{ (1.20)} \qquad 9.16$$

From those we deduct that the distance between the observers is given by:

$$d_u=v_ut_u=v'_ut'_u \qquad 9.17$$

Where we have:

$$\left(1+\frac{v_u}{c}\right)\left(1-\frac{v'_u}{c}\right)=K_uK'_u=1 \qquad 9.18$$

We must observe that at first there is no relationship between the equations 9.11 to 9.14 with the equations 9.15 to 9.18.

With the propagation conditions described we form the following Tables A and B:

Table A

Equation	Clockwise ray of light (c)	Equation	Counter clockwise ray of light (u)	Sum of the rays of light
	Result		Result	
Condition	$x_c=ct_c$	Condition	$x_u=-ct_u$	
1.2	$x'_c=ct_cK_c$	1.2	$x'_u=-ct_uK_u$	
	$x'_c=x_cK_c$		$x'_u=x_uK_u$	$x'_c+x'_u=x_cK_c+x_uK_u$
1.7	$ct'_c=ct_cK_c$	1.7	$ct'_u=ct_uK_u$	$ct'_c+ct'_u=ct_cK_c+ct_uK_u$
	$x'_c=ct'_c$		$x'_u=-ct'_u$	

Table B

Equation	Clockwise ray of light (c)	Equation	Counter clockwise ray of light (u)	Sum of the rays of light
	Result		Result	
Condition	$x'_c=ct'_c$	Condition	$x'_u=-ct'_u$	
1.4	$x_c=ct'_cK'_c$	1.4	$x_u=-ct'_uK'_u$	
	$x_c=x'_cK'_c$		$x_u=x'_uK'_u$	$x_c+x_u=x'_cK'_c+x'_uK'_u$
1.8	$ct_c=ct'_cK'_c$	1.8	$ct_u=ct'_uK'_u$	$ct_c+ct_u=ct'_cK'_c+ct'_uK'_u$
	$x_c=ct_c$		$x_u=-ct_u$	

We observe that for the rays of light with the same direction the Tables A and B are inverse from each other.

Forming the equations group of the sum of the rays of light of the Tables A and B:

$$\begin{cases} D'=ct'_c+ct'_u=ct_cK_c+ct_uK_u & A \\ D=ct_c+ct_u=ct'_cK'_c+ct'_uK'_u & B \end{cases} \qquad 9.19$$

Where for the observer O' $D'=A_u \leftrightarrow A_c$ is the distance between the wave fronts A_u and A_c and where for the observer O $D=A_u \leftrightarrow A_c$ is the distance between the wave fronts A_u and A_c.

In the equations above 9.19 due to the isotropy of the space and time and the wave fronts $A_u \leftrightarrow A_c$ of the rays of light being the same for both observers, the sumo of the rays of light and of times must be invariable between the observers, which is expressed by:

$$D'=D \Rightarrow ct'_c+ct'_u=ct_c+ct_u \Rightarrow \sum t'=\sum t \qquad 9.20$$

This result that equations the isotropy of space and time can be called as the space and time conservation principle.

The three hypothesis of propagations defined next will be applied in 9.19 and tested to prove the compliance of the conservation of space and time principle given by 9.20. With these hypotheses we create a bond between the equations 9.11 to 9.14 with the equations 9.15 to 9.18.

Hypothesis A:

If the space and time are isotropic and there is movement with any privilege of any observer over each other in the empty space then the propagation geometry of the rays of light is equationed by:

$$\begin{cases} ct_c=ct'_u \Rightarrow t_c=t'_u \Rightarrow v_c=v'_u \Rightarrow K_c=K'_u & A \\ ct_u=ct'_c \Rightarrow t_u=t'_c \Rightarrow v_u=v'_c \Rightarrow K_u=K'_c & B \end{cases} \qquad 9.21$$

With those we deduct that the distance between the observers is given by:

$$d_c=d_u=v_ct_c=v'_ct'_c=v_ut_u=v'_ut'_u \qquad 9.22$$

Results that applied in the equations A or B of the group 9.19 complies with the conservation of space and time principle given by 9.20, showing that the Doppler effect in the clockwise and counter clockwise rays of light are compensated in the referentials.

Hypothesis B:

If the space and time are isotropic but the observer O is in an absolute resting position in the empty space then the propagation geometry of the rays of light is equationed by:

$$\begin{cases} ct_c=ct_u=ct & A \\ v_c=v_u=v & B \\ v_ct_c=v_ut_u=vt & C \end{cases} \qquad 9.23$$

With those we deduct that the distance between the observers is given by:

$$d_c=d_u=vt=v'_ct'_c=v'_ut'_u \qquad 9.24$$

Results that applied in the equations A or B of the group 9.19 complies with the conservation of space and time principle given by 9.20, showing that the Doppler effect in the clockwise and counter clockwise rays of light are compensated in the referentials.

Hypothesis C:

If the space and time are isotropic but the observer O is in an absolute resting position in the empty space then the propagation geometry of the rays of light is equationed by:

$$\begin{cases} ct'_c = ct'_u = ct' & A \\ v'_c = v'_u = v' & B \\ v'_c t'_c = v'_u t'_u = v't' & C \end{cases} \qquad 9.25$$

With those we deduct that the distance between the observers is given by:

$$d_c = d_u = v't' = v_c t_c = v_u t_u \qquad 9.26$$

Results that applied in the equations A or B of the group 9.19 complies with the conservation of space and time principle given by 9.20, showing that the Doppler effect in the clockwise and counter clockwise rays of light are compensated in the referentials.

In order to obtain the Sagnac effect we consider that the observer O' is in an absolute resting position, hypothesis C above and that the rays of light course must be of $2\pi R$:

$$ct'_c = ct'_u = ct' = 2\pi R \qquad 9.27$$

Applying the hypothesis C in 9.11 and 9.15 we have:

$$t_c = t'_c K'_c \Rightarrow t_c = t'\left(1 + \frac{v'}{c}\right) \qquad 9.28$$

$$t_u = t'_u K'_u \Rightarrow t_u = t'\left(1 - \frac{v'}{c}\right) \qquad 9.29$$

For the observer O the Sagnac effect is given by the time difference between course of the clockwise ray of light and the counter clock ray of $\Delta t = t_c - t_u$ that can be obtained making (9.28 – 9.29) and applying 9.27 making:

$$\Delta t = t_c - t_u = t'\left(1 + \frac{v'}{c}\right) - t'\left(1 - \frac{v'}{c}\right) = \frac{2v't'}{c} = \frac{4\pi R v'}{c^2} \qquad 9.30$$

The equation $\Delta t = \frac{2v't'}{c} = \frac{2v_c t_c}{c} = \frac{2v_u t_u}{c}$ is exactly the result obtained from the geometry analysis of the propagation of the clockwise and counter clockwise rays of light in a circumference showing the coherence of the hypothesis adopted by the Undulating Relativity.

In 9.30 applying 9.12 and 9.16 we have the final result due to v_c and v_u:

$$\Delta t = t_c - t_u = \frac{2v't'}{c} = \frac{4\pi R v'}{c^2} = \frac{4\pi R v_c}{c^2 - c v_c} = \frac{4\pi R v_u}{c^2 + c v_u} \qquad 9.31$$

The classic formula of the Sagnac effect is given as:

$$\Delta t = t_c - t_u = \frac{4\pi R v}{c^2 - v^2} \qquad 9.32$$

From the propagation geometry we have:

$$\Delta t = \frac{2vt}{c} \qquad 9.33$$

The classic times would be given by:

$$t=\frac{2\pi R}{c} \qquad 9.34$$

$$t_c=\frac{2\pi R}{c-v} \qquad 9.35$$

$$t_u=\frac{2\pi R}{c+v} \qquad 9.36$$

Applying 9.34, 9.35 and 9.36 in 9.33 we have:

$$\Delta t=\frac{2v}{c}\frac{2\pi R}{c}=\frac{4\pi Rv}{c^2} \qquad 9.37$$

$$\Delta t_c=\frac{2v}{c}\frac{2\pi R}{(c-v)}=\frac{4\pi Rv}{c^2-cv} \qquad 9.38$$

$$\Delta t_u=\frac{2v}{c}\frac{2\pi R}{(c+v)}=\frac{4\pi Rv}{c^2+cv} \qquad 9.39$$

The results 9.37, 9.38 and 9.39 are completely different from 9.32.

§18 The Michelson & Morley experience

The traditional analysis that supplies the solution for the null result of this experience considers a device in a resting position at the referential of the observer O' that emits two rays of light, one horizontal in the x' direction (clockwise index c) and another vertical in the direction y'. The horizontal ray of light (clockwise index c) runs until a mirror placed in x' = L at this point the ray of light reflects (counter clockwise index u) and returns to the origin of the referential where x' = zero. The vertical ray of light runs until a mirror placed in y' = L reflects and returns to the origin of the referential where y' = zero.

In the traditional analysis according to the speed of light constancy principle for the observer O' the rays of light track is given by:

$$ct'_c=ct'_u=L \qquad 18.01$$

For the observer O' the sum of times of the track of both rays of light along the x' axis is:

$$\sum t'_{x'}=t'_c+t'_u=\frac{L}{c}+\frac{L}{c}=\frac{2L}{c} \qquad 18.02$$

In the traditional analysis for the observer O' the sum of times of the track of both rays of light along the y' axis is:

$$\sum t'_{y'}=t'_+ +t'_-=\frac{L}{c}+\frac{L}{c}=\frac{2L}{c} \qquad 18.03$$

As we have $\sum t'_{x'}=\sum t'_{y'}=\frac{2L}{c}$ there is no interference fringe and it is applied the null result of the Michelson & Morley experience.

In this traditional analysis the identical track of the clockwise and counter clockwise rays of light in the equation 18.01 that originates the null result of the Michelson & Morley experience contradicts the Sagnac effect that is exactly the time difference existing between the track of the clockwise and counter clockwise rays of light.

Based on the Undulating Relativity we make a deeper analysis of the Michelson & Morley experience obtaining a result that complies completely with the Sagnac effect.

Observing that the equation 18.01 corresponds to the hypothesis C of the paragraph §9.

Applying 18.01 in 9.19 we have:

$$\begin{cases} D' = ct'_c + ct'_u = ct_c K_c + ct_u K_u \Rightarrow D' = L + L = ct_c K_c + ct_u K_u & A \\ D = ct_c + ct_u = ct'_c K'_c + ct'_u K'_u \Rightarrow D = ct_c + ct_u = LK'_c + LK'_u = L(K'_c + K'_u) & B \end{cases} \qquad 18.04$$

From 18.04 A we have:

$$D' = 2L = ct_c\left(1 - \frac{v_c}{c}\right) + ct_u\left(1 + \frac{v_u}{c}\right) \Rightarrow D' = 2L = ct_c - v_c t_c + ct_u + v_u t_u \qquad 18.05$$

Where applying 9.26 we have:

$$D' = 2L = ct_c + ct_u \Rightarrow \sum t_x = t_c + t_u = \frac{2L}{c} \qquad 18.06$$

In 18.04 B we have:

$$D = ct_c + ct_u = L\left[\left(1 + \frac{v'_c}{c}\right) + \left(1 - \frac{v'_u}{c}\right)\right] \qquad 18.07$$

Where applying 9.25 B we have:

$$D = ct_c + ct_u = 2L \Rightarrow \sum t_x = t_c + t_u = \frac{2L}{c} \qquad 18.08$$

The equations 18.06 and 18.08 demonstrate that the Doppler effect in the clockwise and counter clockwise rays of light compensate itself in the referential of the observer O resulting in:

$$\sum t'_{y'} = \sum t'_{x'} = \sum t_x = \frac{2L}{c} \qquad 18.09$$

Because of this, according to the Undulating Relativity in the Michelson & Morley experience we can predict that the clockwise ray of light has a different track from the counter clockwise ray of light according to the formula 18.08 obtaining also the null result for the experience and matching then with the Sagnac effect. This supposition cannot be made based on the Einstein's Special Relativity because according to 17.26 we have:

$$\sum t'_{x'} \neq \sum t_x \qquad 18.10$$

§19 Regression of the perihelion of Mercury of 7,13”

Let us imagine the Sun located in the focus of an ellipse that coincides with the origin of a system of coordinates (x,y,z) with no movement in relation to denominated fixed stars and that the planet Mercury is in a movement governed by the force of gravitational attraction with the Sun describing an elliptic orbit in the plan (x,y) according to the laws of Kepler and the formula of the Newton's gravitational attraction law:

$$\vec{F}=\frac{-GM_o m_o}{r^2}\hat{r}=\frac{-(6{,}67.10^{-11})(1{,}98.10^{30})(3{,}28.10^{23})}{r^2}\hat{r}=\frac{-k}{r^2}\hat{r} \qquad 19.01$$

The sub index "o" indicating mass in relative rest to the observer.

To describe the movement we will use the known formulas:

$$\vec{r}=r\hat{r} \qquad 19.02$$

$$\vec{u}=\frac{d\vec{r}}{dt}=\frac{d(r\hat{r})}{dt}=\frac{dr}{dt}\hat{r}+r\frac{d\phi}{dt}\hat{\phi} \qquad 19.03$$

$$u^2=\vec{u}.\vec{u}=\left(\frac{dr}{dt}\right)^2+\left(r\frac{d\phi}{dt}\right)^2 \qquad 19.04$$

$$\vec{a}=\frac{d\vec{u}}{dt}=\frac{d^2\vec{r}}{dt^2}=\frac{d^2(r\hat{r})}{dt^2}=\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}+\left[2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right]\hat{\phi} \qquad 19.05$$

The formula of the relativity force is given by:

$$\vec{F}=\frac{d}{dt}\left(\frac{m_o\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}\right)=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\vec{a}+\frac{m_o}{\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}}\frac{u}{c^2}\frac{du}{dt}\vec{u}=\frac{m_o}{\left(1-\frac{u^2}{c^2}\right)^{3/2}}\left[\left(1-\frac{u^2}{c^2}\right)\vec{a}+\left(\vec{u}\frac{d\vec{u}}{dt}\right)\frac{\vec{u}}{c^2}\right] \qquad 19.06$$

In this the first term corresponds to the variation of the mass with the speed and the second as we will see later in 19.22 corresponds to the variation of the energy with the time.

With this and the previous formulas we obtain:

$$\vec{F}=\frac{m_o}{\left(1-\frac{u^2}{c^2}\right)^{3/2}}\left\{\begin{array}{l}\left(1-\frac{u^2}{c^2}\right)\left\{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}+\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\hat{\phi}\right\}+\\+\left\{\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\right\}\frac{1}{c^2}\left(\frac{dr}{dt}\hat{r}+r\frac{d\phi}{dt}\hat{\phi}\right)\end{array}\right\} \qquad 19.07$$

$$\vec{F}=\frac{m_o}{(1-u^2/c^2)^{3/2}}\left\{\begin{array}{l}\left\{\left(1-\frac{u^2}{c^2}\right)\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+\left\{\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\right\}\frac{1}{c^2}\frac{dr}{dt}\right\}\hat{r}+\\+\left\{\left(1-\frac{u^2}{c^2}\right)\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)+\left\{\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\right\}\frac{r}{c^2}\frac{d\phi}{dt}\right\}\hat{\phi}\end{array}\right\} \qquad 19.08$$

In this we have the transverse and radial component given by:

$$\vec{F}_{\hat{r}}=\frac{m_o}{\left(1-u^2/c^2\right)^{3/2}}\left\{\left(1-\frac{u^2}{c^2}\right)\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+\left\{\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\right\}\frac{1}{c^2}\frac{dr}{dt}\right\}\hat{r} \tag{19.09}$$

$$\vec{F}_{\hat{\phi}}=\frac{m_o}{\left(1-u^2/c^2\right)^{3/2}}\left\{\left(1-\frac{u^2}{c^2}\right)\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)+\left\{\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\right\}\frac{r}{c^2}\frac{d\phi}{dt}\right\}\hat{\phi} \tag{19.10}$$

As the gravitational force is central we should have to null the traverse component $\vec{F}_{\hat{\phi}}=zero$ so we have:

$$\vec{F}_{\hat{\phi}}=\frac{m_o}{\left(1-u^2/c^2\right)^{3/2}}\left\{\left(1-\frac{u^2}{c^2}\right)\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)+\left\{\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]+r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\right\}\frac{r}{c^2}\frac{d\phi}{dt}\right\}\hat{\phi}=zero \tag{19.11}$$

From where we have:

$$\frac{\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)}{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]}=\frac{\frac{-r}{c^2}\frac{dr}{dt}\frac{d\phi}{dt}}{\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]} \qquad \frac{\left(2r\frac{dr}{dt}\frac{d\phi}{dt}+r^2\frac{d^2\phi}{dt^2}\right)}{r^2\frac{d\phi}{dt}}=\frac{\frac{-1}{c^2}\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]}{\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]} \tag{19.12}$$

From the radial component $\vec{F}_{\hat{r}}$ we have:

$$\vec{F}_{\hat{r}}=\frac{m_o}{\left(1-u^2/c^2\right)^{3/2}}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\left\{\left(1-\frac{u^2}{c^2}\right)+\left\{\frac{dr}{dt}+\frac{r\frac{d\phi}{dt}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)}{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]}\right\}\frac{1}{c^2}\frac{dr}{dt}\right\}\hat{r} \tag{19.13}$$

That applying 19.12 we have:

$$\vec{F}_{\hat{r}}=\frac{m_o}{\left(1-u^2/c^2\right)^{3/2}}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\left\{\left(1-\frac{u^2}{c^2}\right)+\left\{\frac{dr}{dt}-\frac{r\frac{d\phi}{dt}\left(\frac{r}{c^2}\frac{dr}{dt}\frac{d\phi}{dt}\right)}{\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]}\right\}\frac{1}{c^2}\frac{dr}{dt}\right\}\hat{r} \tag{19.14}$$

That simplifying results in:

$$\vec{F}_{\hat{r}}=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\frac{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]}{\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]}\hat{r} \tag{19.15}$$

This equaled to Newton's gravitational force results in the relativistic gravitational force:

$$\vec{F}_{\hat{r}}=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\frac{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]}{\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]}\hat{r}=\frac{-GM_om_o}{r^2}\hat{r}=\frac{-k}{r^2}\hat{r} \tag{19.16}$$

As the gravitational force is central it should assist the theory of conservation of the energy (E) that is written as:

$E = E_k + E_p$ = constant. 19.17

Where the kinetic energy (E_k) is given by:

$$E_k = mc^2 - m_o c^2 = m_o c^2 \left(\frac{1}{\sqrt{1 - \frac{u^2}{c^2}}} - 1 \right) \quad 19.18$$

And the potential energy (E_p) gravitational by:

$$E_p = \frac{-GM_o m_o}{r} = \frac{-k}{r} \quad 19.19$$

Resulting in:

$$E = m_o c^2 \left[\frac{1}{\sqrt{1 - \frac{u^2}{c^2}}} - 1 \right] - \frac{k}{r} = \text{Constant} . \quad 19.20$$

As the total energy (E) it is constant we should have:

$$\frac{dE}{dt} = \frac{dE_k}{dt} + \frac{dE_p}{dt} = zero . \quad 19.21$$

Then we have:

$$\frac{dE_k}{dt} = \frac{m_o u}{\left(1 - \frac{u^2}{c^2} \right)^{\frac{3}{2}}} \frac{du}{dt} \quad 19.22$$

$$\frac{dE_p}{dt} = \frac{k}{r^2} \frac{dr}{dt} \quad 19.23$$

Resulting in:

$$\frac{dE}{dt} = \frac{dE_k}{dt} + \frac{dE_p}{dt} = zero \Rightarrow \frac{m_o u}{\left(1 - \frac{u^2}{c^2} \right)^{\frac{3}{2}}} \frac{du}{dt} + \frac{k}{r^2} \frac{dr}{dt} = zero \Rightarrow \frac{m_o u}{\left(1 - \frac{u^2}{c^2} \right)^{\frac{3}{2}}} \frac{du}{dt} = \frac{-k}{r^2} \frac{dr}{dt} \quad 19.24$$

This applied in the relativistic force 19.06 and equaled to the gravitational force 19.01 results in:

$$\vec{F} = \frac{m_o}{\sqrt{1 - \frac{u^2}{c^2}}} \vec{a} - \frac{1}{c^2} \frac{k}{r^2} \frac{dr}{dt} \vec{u} = \frac{-k}{r^2} \hat{r} \quad 19.25$$

In this substituting the previous variables we get:

$$\vec{F}=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\left\{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}+\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\hat{\phi}\right\}-\frac{1}{c^2}\frac{k}{r^2}\frac{dr}{dt}\left(\frac{dr}{dt}\hat{r}+r\frac{d\phi}{dt}\hat{\phi}\right)=\frac{-k}{r^2}\hat{r} \qquad 19.26$$

From this we obtain the radial component $\vec{F}_{\hat{r}}$ equals to:

$$\vec{F}_{\hat{r}}=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]-\frac{1}{c^2}\frac{k}{r^2}\left(\frac{dr}{dt}\right)^2=\frac{-k}{r^2} \qquad 19.27$$

That easily becomes the relativistic gravitational force 19.16.

From 19.26 we obtain the traverse component $\vec{F}_{\hat{\phi}}$ equals to:

$$\vec{F}_{\hat{\phi}}=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)-\frac{1}{c^2}\frac{k}{r}\frac{dr}{dt}\frac{d\phi}{dt}=zero \qquad 19.28$$

From this last one we have:

$$\frac{2r\frac{dr}{dt}\frac{d\phi}{dt}+r^2\frac{d^2\phi}{dt^2}}{r^2\frac{d\phi}{dt}}=\frac{1}{m_o}\frac{k}{c^2r^2}\frac{dr}{dt}\sqrt{1-\frac{u^2}{c^2}} \qquad 19.29$$

As the gravitational force is central it should also assist the theory of conservation of the angular moment that is written as:

$\vec{L}=\vec{r}\times\vec{p}=$constant. 19.30

$$\vec{L}=\vec{r}\times\vec{p}=\vec{r}\times\frac{m_o\vec{u}}{\sqrt{1-\frac{u^2}{c^2}}}=r\hat{r}\times\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\left(\frac{dr}{dt}\hat{r}+r\frac{d\phi}{dt}\hat{\phi}\right)=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}r^2\frac{d\phi}{dt}\left(\hat{r}\times\hat{\phi}\right)=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}r^2\frac{d\phi}{dt}\hat{k} \qquad 19.31$$

$\vec{L}=\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}r^2\frac{d\phi}{dt}\vec{k}=L\hat{k}$ = constant. 19.32

$$\frac{d\vec{L}}{dt}=\frac{d\left(L\hat{k}\right)}{dt}=\frac{d(L)\hat{k}}{dt}+\frac{Ld\left(\hat{k}\right)}{dt}=\frac{d(L)\hat{k}}{dt}=zero\Rightarrow\frac{d(L)}{dt}=zero \qquad 19.33$$

Resulting in L that is constant.

In 19.33 we had $\frac{d\hat{k}}{dt}=zero$ because the movement is in the plane (x,y).

Deriving L we find:

$$\frac{dL}{dt}=\frac{d}{dt}\left(\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}r^2\frac{d\phi}{dt}\right)=\frac{1}{c^2}\frac{m_o u}{\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}}\frac{du}{dt}r^2\frac{d\phi}{dt}+\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}\left(2r\frac{dr}{dt}\frac{d\phi}{dt}+r^2\frac{d^2\phi}{dt^2}\right)=zero \qquad 19.34$$

From that we have:

$$\frac{\left(2r\frac{dr}{dt}\frac{d\phi}{dt}+r^2\frac{d^2\phi}{dt^2}\right)}{r^2\frac{d\phi}{dt}}=\frac{-u}{\left(1-\frac{u^2}{c^2}\right)}\frac{du}{dt}\frac{1}{c^2} \qquad 19.35$$

Equaling 19.12 originating from the theory of the central force with 19.29 originating from the theory of conservation of the energy and 19.35 originating from the theory of conservation of the angular moment we have:

$$\frac{\left(2r\frac{dr}{dt}\frac{d\phi}{dt}+r^2\frac{d^2\phi}{dt^2}\right)}{r^2\frac{d\phi}{dt}}=\frac{\frac{-1}{c^2}\frac{dr}{dt}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]}{\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]}=\frac{k}{m_o c^2 r^2}\frac{dr}{dt}\sqrt{1-\frac{u^2}{c^2}}=\frac{-u}{\left(1-\frac{u^2}{c^2}\right)}\frac{du}{dt}\frac{1}{c^2} \qquad 19.36$$

From the last two equality we obtain 19.24 and from the two of the middle we obtain 19.16.

For solution of the differential equations we will use the same method used in the Newton's theory.

Let us assume $w=\frac{1}{r}$ 19.37

The differential total of this is $dw=\frac{\partial w}{\partial r}dr\Rightarrow dw=\frac{-1}{r^2}dr$ 19.38

From where we have $\frac{dw}{d\phi}=\frac{-1}{r^2}\frac{dr}{d\phi}$ e $\frac{dw}{dt}=\frac{-1}{r^2}\frac{dr}{dt}$ 19.39

From the module of the angular moment we have $\frac{d\phi}{dt}=\frac{L}{m_o r^2}\sqrt{1-\frac{u^2}{c^2}}$ 19.40

From where we have $\frac{dr}{dt}=\frac{L}{m_o r^2}\frac{dr}{d\phi}\sqrt{1-\frac{u^2}{c^2}}$ 19.41

Where applying 19.39 we have $\frac{dr}{dt}=\frac{-L}{m_o}\frac{dw}{d\phi}\sqrt{1-\frac{u^2}{c^2}}$ 19.42

That derived supplies $\frac{d^2r}{dt^2}=\frac{d\phi}{dt}\frac{dt}{d\phi}\frac{d}{dt}\left(\frac{-L}{m_o}\frac{dw}{d\phi}\sqrt{1-\frac{u^2}{c^2}}\right)$ 19.43

Where applying 19.40 and deriving we have:

$$\frac{d^2r}{dt^2}=\frac{L}{m_o r^2}\sqrt{1-\frac{u^2}{c^2}}\frac{d}{d\phi}\left(\frac{-L}{m_o}\frac{dw}{d\phi}\sqrt{1-\frac{u^2}{c^2}}\right)=\frac{-L^2}{m_o^2 r^2}\sqrt{1-\frac{u^2}{c^2}}\left[\frac{d^2w}{d\phi^2}\sqrt{1-\frac{u^2}{c^2}}+\frac{dw}{d\phi}\frac{d}{d\phi}\left(\sqrt{1-\frac{u^2}{c^2}}\right)\right] \qquad 19.44$$

In this with 19.36 the radical derived is obtained this way:

$$\frac{d}{dt}\left(\sqrt{1-\frac{u^2}{c^2}}\right)=\frac{-1}{\sqrt{1-u^2/c^2}}\frac{u}{c^2}\frac{du}{dt}=\frac{k}{m_o c^2 r^2}\frac{dr}{dt}\left(1-\frac{u^2}{c^2}\right)=\frac{-k}{m_o c^2}\frac{dw}{dt}\left(1-\frac{u^2}{c^2}\right) \qquad 19.45$$

$$\frac{d}{d\phi}\left(\sqrt{1-\frac{u^2}{c^2}}\right)=\frac{-1}{\sqrt{1-u^2/c^2}}\frac{u}{c^2}\frac{du}{d\phi}=\frac{k}{m_o c^2 r^2}\frac{dr}{d\phi}\left(1-\frac{u^2}{c^2}\right)=\frac{-k}{m_o c^2}\frac{dw}{d\phi}\left(1-\frac{u^2}{c^2}\right) \qquad 19.46$$

That applied in 19.44 supplies:

$$\frac{d^2r}{dt^2}=\frac{-L^2}{m_o^2 r^2}\sqrt{1-\frac{u^2}{c^2}}\left[\frac{d^2w}{d\phi^2}\sqrt{1-\frac{u^2}{c^2}}-\frac{k}{m_o c^2}\left(\frac{dw}{d\phi}\right)^2\left(1-\frac{u^2}{c^2}\right)\right] \qquad 19.47$$

Simplified results:

$$\frac{d^2r}{dt^2}=\frac{L^2k}{m_o^3 c^2 r^2}\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}\left(\frac{dw}{d\phi}\right)^2-\frac{L^2}{m_o^2 r^2}\left(1-\frac{u^2}{c^2}\right)\frac{d^2w}{d\phi^2} \qquad 19.48$$

Let us find the second derived of the angle deriving 19.40:

$$\frac{d^2\phi}{dt^2}=\frac{d}{dt}\left(\frac{L}{m_o r^2}\sqrt{1-\frac{u^2}{c^2}}\right)=\frac{-2L}{m_o r^3}\frac{dr}{dt}\sqrt{1-\frac{u^2}{c^2}}+\frac{L}{m_o r^2}\frac{d}{dt}\left(\sqrt{1-\frac{u^2}{c^2}}\right) \qquad 19.49$$

In this applying 19.42 and 19.45 and simplifying we have:

$$\frac{d^2\phi}{dt^2}=\frac{2L^2}{m_o^2 r^3}\frac{dw}{d\phi}\left(1-\frac{u^2}{c^2}\right)-\frac{L^2k}{m_o^3 c^2 r^4}\frac{dw}{d\phi}\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}} \qquad 19.50$$

Applying in 19.04 the equations 19.40 and 19.42 and simplifying we have:

$$u^2=\frac{L^2}{m_o^2}\left(1-\frac{u^2}{c^2}\right)\left[\left(\frac{dw}{d\phi}\right)^2+\frac{1}{r^2}\right] \qquad 19.51$$

The equation of the relativistic gravitational force 19.16 remodeled is:

$$\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2=\sqrt{1-\frac{u^2}{c^2}}\left[1-\frac{1}{c^2}\left(\frac{dr}{dt}\right)^2\right]\frac{-k}{m_o r^2} \qquad 19.52$$

In this applying the formulas above we have:

$$\frac{L^2k}{m_o^3 c^2 r^2}\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}\left(\frac{dw}{d\phi}\right)^2-\frac{L^2}{m_o^2 r^2}\left(1-\frac{u^2}{c^2}\right)\frac{d^2w}{d\phi^2}-r\left(\frac{L}{m_o r^2}\sqrt{1-\frac{u^2}{c^2}}\right)^2=\sqrt{1-\frac{u^2}{c^2}}\left[1-\frac{1}{c^2}\left(\frac{-L}{m_o}\frac{dw}{d\phi}\sqrt{1-\frac{u^2}{c^2}}\right)^2\right]\frac{-k}{m_o r^2}$$

$$\frac{L^2k}{m_o^3c^2r^2}\left(1-\frac{u^2}{c^2}\right)\left(\frac{dw}{d\phi}\right)^2-\frac{L^2}{m_o^2r^2}\sqrt{1-\frac{u^2}{c^2}}\frac{d^2w}{d\phi^2}-\frac{L^2}{m_o^2r^3}\sqrt{1-\frac{u^2}{c^2}}=\left[1-\frac{1}{c^2}\left(\frac{-L}{m_o}\frac{dw}{d\phi}\sqrt{1-\frac{u^2}{c^2}}\right)^2\right]\frac{-k}{m_or^2}$$

$$\frac{L^2k}{m_o^3c^2r^2}\left(1-\frac{u^2}{c^2}\right)\left(\frac{dw}{d\phi}\right)^2-\frac{L^2}{m_o^2r^2}\sqrt{1-\frac{u^2}{c^2}}\frac{d^2w}{d\phi^2}-\frac{L^2}{m_o^2r^3}\sqrt{1-\frac{u^2}{c^2}}=\frac{-k}{m_or^2}+\frac{L^2k}{m_o^3r^2}\frac{1}{c^2}\left(1-\frac{u^2}{c^2}\right)\left(\frac{dw}{d\phi}\right)^2$$

$$-\frac{L^2}{m_o^2r^2}\sqrt{1-\frac{u^2}{c^2}}\frac{d^2w}{d\phi^2}-\frac{L^2}{m_o^2r^3}\sqrt{1-\frac{u^2}{c^2}}=\frac{-k}{m_or^2}$$

$$\frac{d^2w}{d\phi^2}+\frac{1}{r}=\frac{m_ok}{L^2\sqrt{1-\frac{u^2}{c^2}}}$$

$$\frac{d^2w}{d\phi^2}+\frac{1}{r}=\frac{m_ok}{\left(\frac{m_o}{\sqrt{1-\frac{u^2}{c^2}}}r^2\frac{d\phi}{dt}\right)^2\sqrt{1-\frac{u^2}{c^2}}}$$

$$\frac{d^2w}{d\phi^2}+\frac{1}{r}=\frac{m_ok\sqrt{1-\frac{u^2}{c^2}}}{m_o^2r^4\left(\frac{d\phi}{dt}\right)^2}$$

$$\left(\frac{d^2w}{d\phi^2}+\frac{1}{r}\right)^2=\left[\frac{k\sqrt{1-\frac{u^2}{c^2}}}{m_or^4\left(\frac{d\phi}{dt}\right)^2}\right]^2$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2\left(1-\frac{u^2}{c^2}\right)}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}u^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}\left[\left(\frac{dr}{dt}\right)^2+\left(r\frac{d\phi}{dt}\right)^2\right]}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}\left(\frac{dr}{dt}\right)^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}\left(r\frac{d\phi}{dt}\right)^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}\left(\frac{dr}{d\phi}\right)^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^2}-\frac{k^2}{m_o^2c^2r^6\left(\frac{d\phi}{dt}\right)^2}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}\left(-r^2\frac{dw}{d\phi}\right)^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^2}-\frac{k^2}{m_o^2c^2r^6\left(\frac{d\phi}{dt}\right)^2}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2r^8\left(\frac{d\phi}{dt}\right)^4}-\frac{\frac{k^2}{c^2}\left(\frac{dw}{d\phi}\right)^2}{m_o^2r^4\left(\frac{d\phi}{dt}\right)^2}-\frac{k^2}{m_o^2c^2r^6\left(\frac{d\phi}{dt}\right)^2}$$

In this we will consider constant the Newton's angular moment in the form:

$$L=r^2\frac{d\phi}{dt} \qquad 19.53$$

That it is really the known theoretical angular moment.

$$\left(\frac{d^2w}{d\phi^2}\right)^2+\frac{2}{r}\frac{d^2w}{d\phi^2}+\frac{1}{r^2}=\frac{k^2}{m_o^2L^4}-\frac{k^2}{m_o^2c^2L^2}\left(\frac{dw}{d\phi}\right)^2-\frac{k^2}{m_o^2c^2r^2L^2}$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+2\frac{d^2w}{d\phi^2}2w+w^2=\frac{k^2}{m_o^2L^4}-\frac{k^2}{m_o^2c^2L^2}\left(\frac{dw}{d\phi}\right)^2-\frac{k^2}{m_o^2c^2L^2}w^2$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+2\frac{d^2w}{d\phi^2}w+w^2=B-A\left(\frac{dw}{d\phi}\right)^2-Aw^2$$

$$\left(\frac{d^2w}{d\phi^2}\right)^2+2\frac{d^2w}{d\phi^2}w+A\left(\frac{dw}{d\phi}\right)^2+(A+1)w^2-B=zero \qquad 19.54$$

Where we have:

$$A=\frac{k^2}{m_o^2c^2L^2} \qquad 19.55$$

$$B=\frac{k^2}{m_o^2L^4} \qquad 19.56$$

The equation 19.54 has as solution:

$$w=\frac{1}{\varepsilon D}\left[1-\varepsilon cos\left(\phi\sqrt{1+A}+\phi_o\right)\right]\Rightarrow w=\frac{1}{\varepsilon D}\left[1-\varepsilon cos(\phi Q)\right] \qquad 19.57$$

Where we consider $\phi_o=zero$.

It is denominated in 19.57 $Q^2=1+A$. 19.58

The equation 19.58 is function only of A demonstrating the intrinsic union between the variation of the mass with the variation of the energy in the time, because both as already described, participate in the relativistic force 19.06 in this relies the essential difference between the mass and the electric charge that is invariable and indivisible in the electromagnetic theory.

From 19.57 we obtain the ray of a conical:

$$r=\frac{1}{w}=\frac{\varepsilon D}{1-\varepsilon cos\left(\phi\sqrt{1+A}\right)}\Rightarrow r=\frac{\varepsilon D}{1-\varepsilon cos(\phi Q)} \qquad 19.59$$

Where ε is the eccentricity and D the directory distance of the focus.

Deriving 19.57 we have $\frac{dw}{d\phi}=\frac{Q\,sen(\phi Q)}{D}$ 19.60

That derived results in $\frac{d^2w}{d\phi^2}=\frac{Q^2 cos(\phi Q)}{D}$ 19.61

Applying in 19.54 the variables we have:

$$\left(\frac{d^2w}{d\phi^2}\right)^2+2\frac{d^2w}{d\phi^2}w+A\left(\frac{dw}{d\phi}\right)^2+(A+1)w^2-B=zero.$$

$$\frac{Q^4cos^2(\phi Q)}{D^2}+2\frac{Q^2cos(\phi Q)}{D}\left[\frac{1-\varepsilon cos(\phi Q)}{\varepsilon D}\right]+A\frac{Q^2sen^2(\phi Q)}{D^2}+(A+1)\left[\frac{1-\varepsilon cos(\phi Q)}{\varepsilon D}\right]^2-B=zero \qquad 19.62$$

$$\frac{Q^4cos^2(\phi Q)}{D^2}+2\frac{Q^2cos(\phi Q)}{\varepsilon D^2}-2\frac{Q^2cos^2(\phi Q)}{D^2}+A\frac{Q^2}{D^2}-A\frac{Q^2cos^2(\phi Q)}{D^2}+(A+1)\left[\frac{1-\varepsilon cos(\phi Q)}{\varepsilon D}\right]^2-B=zero$$

$$\frac{Q^4cos^2(\phi Q)}{D^2}+2\frac{Q^2cos(\phi Q)}{\varepsilon D^2}-2\frac{Q^2cos^2(\phi Q)}{D^2}+A\frac{Q^2}{D^2}-A\frac{Q^2cos^2(\phi Q)}{D^2}+\frac{(A+1)}{\varepsilon^2D^2}-2\frac{(A+1)cos(\phi Q)}{\varepsilon D^2}+\frac{(A+1)cos^2(\phi Q)}{D^2}-B=zero$$

$$\left(Q^4-2Q^2-AQ^2+A+1\right)\frac{cos^2(\phi Q)}{D^2}+\left(\frac{2Q^2}{\varepsilon D}-\frac{2A}{\varepsilon D}-\frac{2}{\varepsilon D}\right)\frac{cos(\phi Q)}{D}+\frac{AQ^2}{D^2}+\frac{(A+1)}{\varepsilon^2D^2}-B=zero \qquad 19.63$$

In this applying in the first parenthesis $Q^2=1+A$ we have:

$$\left(Q^4-2Q^2-AQ^2+A+1\right)=\left[(1+A)^2-2(1+A)-A(1+A)+A+1\right]=\left(1+2A+A^2-2-2A-A-A^2+A+1\right)=zero$$

In 19.63 applying in the second parenthesis $Q^2=1+A$ we have:

$$\left(\frac{2Q^2}{\varepsilon D}-\frac{2A}{\varepsilon D}-\frac{2}{\varepsilon D}\right)=\left[\frac{2(1+A)}{\varepsilon D}-\frac{2A}{\varepsilon D}-\frac{2}{\varepsilon D}\right]=zero$$

The rest of the equation 19.63 is therefore:

$$\frac{AQ^2}{D^2}+\frac{(A+1)}{\varepsilon^2 D^2}-B=zero \qquad 19.64$$

The data of the elliptic orbit of the planet Mercury is [1]:

Eccentricity of the orbit $\varepsilon=0{,}206$.

Larger semi-axis = a = $5{,}79.10^{10}$m.

Smaller semi-axis $b=a\sqrt{1-\varepsilon^2}=5{,}79.10^{10}\sqrt{1-0{,}206^2}=56.658.160.305{,}80m$.

$\varepsilon D=a\left(1-\varepsilon^2\right)=5{,}79.10^{10}\left(1-0{,}206^2\right)=55.442.955.600{,}00m$.

$$D=\frac{a\left(1-\varepsilon^2\right)}{\varepsilon}=\frac{5{,}79.10^{10}\left(1-0{,}206^2\right)}{0{,}206}=269.140.561.165{,}00m.$$

The orbital period of the Earth (PT) and Mercury (PM) around the Sun in seconds are:

$PT=3{,}16\,.\,10^7\,s.$

$PM=7{,}60\,.\,10^6\,s.$

The number of turns that Mercury (m_o) makes around the Sun (M_o) in one century is, therefore:

$$N=100\frac{3{,}16\,.\,10^7}{7{,}60\,.\,10^6}=415{,}79\,. \qquad 19.65$$

Theoretical angular moment of Mercury:

$$L^2=\left(r^2\frac{d\phi}{dt}\right)^2=GM_o a\left(1-\varepsilon^2\right)=6{,}67.10^{-11}1{,}98.10^{30}5{,}79.10^{10}\left(1-0{,}206^2\right)=7{,}32212937427.10^{30} \qquad 19.66$$

$$A=\frac{(GM_0 m_o)^2}{m_o^2c^2L^2}=\frac{(GM_0)^2}{c^2L^2}=\frac{\left(6{,}67.10^{-11}\right)^2\left(1{,}98.10^{30}\right)^2}{\left(3{,}0.10^8\right)^2\left(7{,}32.10^{30}\right)}=2{,}65.10^{-8}. \qquad 19.67$$

$$B=\frac{(GM_0 m_o)^2}{m_o^2L^4}=\frac{(GM_0)^2}{L^4}=\frac{\left(6{,}67.10^{-11}\right)^2\left(1{,}98.10^{30}\right)^2}{\left(7{,}32.10^{30}\right)^2}=3{,}25.10^{-22} \qquad 19.68$$

$$Q=\sqrt{1+A}=\sqrt{1+2{,}63.10^{-8}}=1{,}000.000.013.23 \qquad 19.69$$

Applying the numeric data with several decimal numbers to the rest of the equation 19.63 we have:

$$\frac{AQ^2}{D^2}+\frac{(A+1)}{\varepsilon^2D^2}-B=\frac{2{,}65.10^{-8}(1{,}000.000.013.23)^2}{(269.140.561.165{,}00)^2}+\frac{2{,}65.10^{-8}+1}{(55.442.955.600{,}00)^2}-3{,}25.10^{-22}=8{,}976.10^{-30} \qquad 19.70$$

Result that we can consider null.

We will obtain the relativistic angular moment of the rest of the equation 19.63 in this applying the variables we have:

$$\frac{AQ^2}{D^2}+\frac{(A+1)}{\varepsilon^2 D^2}-B=\frac{(GM_0)^2}{c^2L^2D^2}\left[1+\frac{(GM_0)^2}{c^2L^2}\right]+\frac{1}{\varepsilon^2D^2}\left[1+\frac{(GM_0)^2}{c^2L^2}\right]-\frac{(GM_0)^2}{L^4}=zero \qquad 19.71$$

$$\varepsilon^2L^2(GM_0)^2\left[1+\frac{(GM_0)^2}{c^2L^2}\right]+L^4c^2\left[1+\frac{(GM_0)^2}{c^2L^2}\right]-c^2\varepsilon^2D^2(GM_0)^2=zero$$

$$\varepsilon^2L^2(GM_0)^2+\varepsilon^2L^2(GM_0)^2\frac{(GM_0)^2}{c^2L^2}+L^4c^2+L^4c^2\frac{(GM_0)^2}{c^2L^2}-c^2\varepsilon^2D^2(GM_0)^2=zero$$

$$\varepsilon^2L^2(GM_0)^2+\varepsilon^2\frac{(GM_0)^4}{c^2}+L^4c^2+L^2(GM_0)^2-c^2\varepsilon^2D^2(GM_0)^2=zero$$

$$c^2L^4+(1+\varepsilon^2)(GM_0)^2L^2+\varepsilon^2\frac{(GM_0)^4}{c^2}-c^2\varepsilon^2D^2(GM_0)^2=zero \qquad 19.72$$

$$L^2=\frac{-(1+\varepsilon^2)(GM_0)^2\pm\sqrt{\left[(1+\varepsilon^2)(GM_0)^2\right]^2-4c^2\left[\varepsilon^2\frac{(GM_0)^4}{c^2}-c^2\varepsilon^2D^2(GM_0)^2\right]}}{2c^2}$$

$$L^2=\frac{-(1+\varepsilon^2)(GM_0)^2\pm\sqrt{(1+\varepsilon^2)^2(GM_0)^4-4\varepsilon^2(GM_0)^4+4c^4\varepsilon^2D^2(GM_0)^2}}{2c^2}$$

$$L^2=\frac{-(1+\varepsilon^2)(GM_0)^2\pm\sqrt{(1+2\varepsilon^2+\varepsilon^4)(GM_0)^4-4\varepsilon^2(GM_0)^4+4c^4\varepsilon^2D^2(GM_0)^2}}{2c^2}$$

$$L^2=\frac{-(1+\varepsilon^2)(GM_0)^2\pm\sqrt{(GM_0)^4+2\varepsilon^2(GM_0)^4+\varepsilon^4(GM_0)^4-4\varepsilon^2(GM_0)^4+4c^4\varepsilon^2D^2(GM_0)^2}}{2c^2}$$

$$L^2=\frac{-(1+\varepsilon^2)(GM_0)^2\pm\sqrt{(GM_0)^4+\varepsilon^4(GM_0)^4-2\varepsilon^2(GM_0)^4+4c^4\varepsilon^2D^2(GM_0)^2}}{2c^2}$$

$$L^2=\frac{-(1+\varepsilon^2)(GM_0)^2+\sqrt{(1-\varepsilon^2)^2(GM_0)^4+4c^4\varepsilon^2D^2(GM_0)^2}}{2c^2}=7{,}32212927328.10^{30} \qquad 19.73$$

This last equation has the exclusive property of relating the speed c to the denominated relativistic angular moment that is smaller than the theoretical angular moment 19.66.

The variation of the relativistic angular moment in relation to the theoretical angular moment is very small and given by:

$$\Delta L=\frac{7{,}32212927328.10^{30}-7{,}32212937427.10^{30}}{7{,}32212937427.10^{30}}=-1{,}38.10^{-8}=\frac{-1}{72.503.509{,}00}. \qquad 19.74$$

That demonstrates the accuracy of the principle of constancy of the speed of the light.

In reality, the equation 19.06 provides a secular retrocession perihelion of Mercury, which is given by in

$$\Delta\phi = 2\pi 415{,}79\left(\frac{1}{Q}-1\right) = 2\pi 415{,}79(-0{,}000.000.013.23) = -3{,}46.10^{-5}\,rad. \qquad 19.75$$

Converting for the second we have:

$$\Delta\phi = \frac{-3{,}46.10^{-5}.180{,}00.3.600{,}00}{\pi} = -7{,}13'' . \qquad 19.76$$

This retrocession, is not expected in Newtonian theory is due to relativistic variation of mass and energy and is shrouded in total observed precession of 5599. "

§§19 Advance of Mercury's perihelion of 42.79"

If we write the equation for the gravitational relativity energy E_R covering the terms for the kinetic energy, the potential energy E_p and the resting energy:

$$E_R = m_o c^2 \left(\frac{1}{\sqrt{1-\frac{u^2}{c^2}}} - 1 \right) + E_p + m_o c^2 = \frac{m_o c^2}{\sqrt{1-\frac{u^2}{c^2}}} + E_p . \qquad 19.77$$

Being the conservative the gravitational force its energy is constant. Assuming then that in 19.77 when the radius tends to infinite, the speed and potential energy tends to zero, resulting then:

$$E_R = \frac{m_o c^2}{\sqrt{1-\frac{u^2}{c^2}}} + E_p = m_o c^2 \qquad 19.78$$

Writing the equation to the Newton's gravitation energy E_N having the correspondent Newton's terms to the 19.77:

$$E_N = \frac{m_o u^2}{2} - \frac{k}{r} + m_o c^2 = m_o c^2 \qquad 19.79$$

Where $\frac{m_o u^2}{2}$ is the kinetic energy, $\frac{-k}{r}$ the potential energy and $m_o c^2$ the resting energy or better saying the inertial energy.

From this 19.79 we have:

$$\frac{m_o u^2}{2} - \frac{k}{r} + m_o c^2 = m_o c^2 \Rightarrow \frac{m_o u^2}{2} = \frac{k}{r} \Rightarrow u^2 = \frac{2k}{m_o r} = \frac{2GM_o m_o}{m_o r} \Rightarrow u^2 = \frac{2GM_o}{r} \qquad 19.80$$

Deriving 19.79 we have:

$$\frac{dE_N}{dt} = \frac{d}{dt}\left(\frac{m_o u^2}{2} - \frac{k}{r} + m_o c^2 \right) = \text{zero}$$

$$\frac{m_o 2u}{2}\frac{du}{dt} + \frac{k}{r^2}\frac{dr}{dt} = \text{zero}$$

$$u\frac{du}{dt} = \frac{-k}{m_o r^2}\frac{dr}{dt} = \frac{-GM_o}{r^2}\frac{dr}{dt}$$

$$u\frac{du}{dt} = \frac{-GM_o}{r^2}\frac{dr}{dt}$$

$$u\frac{du}{dr} = \frac{-GM_o}{r^2} \qquad 19.81$$

Making the relativity energy 19.78 equal to the Newton's energy 19.79 we have:

$$E_R = E_N \Rightarrow \frac{m_o c^2}{\sqrt{1-\frac{u^2}{c^2}}} + E_p = \frac{m_o u^2}{2} - \frac{k}{r} + m_o c^2 \qquad 19.82$$

$$\frac{m_o c^2}{m_o\sqrt{1-\frac{u^2}{c^2}}} + \frac{E_p}{m_o} = \frac{m_o u^2}{m_o 2} - \frac{GM_o m_o}{m_o r} + \frac{m_o c^2}{m_o} \qquad 19.83$$

In that denominating the relativity potential (φ) as:

$$\varphi = \frac{E_p}{m_o} \qquad 19.84$$

We have:

$$\frac{c^2}{\sqrt{1-\frac{u^2}{c^2}}} + \varphi = \frac{u^2}{2} - \frac{GM_o}{r} + c^2$$

$$\varphi = \frac{u^2}{2} - \frac{GM_o}{r} + c^2 - \frac{c^2}{\sqrt{1-\frac{u^2}{c^2}}} \qquad 19.85$$

In this one replacing the approximation:

$$\frac{1}{\sqrt{1-\frac{u^2}{c^2}}} \approx 1 + \frac{u^2}{2c^2} \qquad 19.86$$

We have:

$$\varphi = \frac{u^2}{2} - \frac{GM_o}{r} + c^2 - c^2\left(1 + \frac{u^2}{2c^2}\right)$$

That simplified results in the Newton's potential:

$$\varphi = \frac{u^2}{2} - \frac{GM_o}{r} + c^2 - c^2 - \frac{u^2}{2} = \frac{-GM_o}{r} \qquad 19.87$$

Replacing 19.84 and the relativity potential 19.85 in the relativity energy 19.78:

$$E_R = \frac{m_o c^2}{\sqrt{1-\frac{u^2}{c^2}}} + m_o\left(\frac{u^2}{2} - \frac{GM_o}{r} + c^2 - \frac{c^2}{\sqrt{1-\frac{u^2}{c^2}}}\right) \qquad 19.88$$

We have the Newton's energy 19.79:

$$E_N = \frac{m_o u^2}{2} - \frac{GM_o m_o}{r} + m_o c^2$$

Deriving the relativity potential 19.85 we have the relativity gravitational acceleration modulus exactly as in the Newton's theory:

$$a=\frac{-d\varphi}{dr}$$

$$a=\frac{-d\varphi}{dr}=\frac{-d}{dr}\left(\frac{u^2}{2}-\frac{GM_o}{r}+c^2-\frac{c^2}{\sqrt{1-\frac{u^2}{c^2}}}\right)$$

$$a=\frac{-d}{dr}\left(\frac{u^2}{2}-\frac{GM_o}{r}+c^2\right)-\frac{d}{dr}\left(-\frac{c^2}{\sqrt{1-\frac{u^2}{c^2}}}\right)$$

Where we have:

$\frac{-d}{dr}\left(\frac{u^2}{2}-\frac{GM_o}{r}+c^2\right)=\frac{-d}{dr}\left(\frac{E_N}{m_o}\right)=\text{zero}$. Because the term to be derived is the Newton's energy divided by m_o that is $\frac{E_N}{m_o}=\frac{u^2}{2}-\frac{GM_o}{r}+c^2$ that is constant, resulting then in:

$$a=-\frac{d}{dr}\left(-\frac{c^2}{\sqrt{1-\frac{u^2}{c^2}}}\right)$$

$$a=-\left[-\frac{u}{\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}}\frac{du}{dr}\right]$$

In this one applying 19.81 we have:

$$a=\frac{-1}{\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}}\frac{GM_o}{r^2} \qquad 19.89$$

The vector acceleration is given by 19.05:

$$\vec{a}=\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}+\left[2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right]\hat{\phi}$$

The relativity gravitational acceleration modulus 19.89 is equal to the component of the vector radius ($\hat{r}$) thus we have:

$$a=\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]=\frac{-1}{\left(1-\frac{u^2}{c^2}\right)^{\frac{3}{2}}}\frac{GM_o}{r^2} \qquad 19.90$$

Being null the transversal acceleration we have:

$$\left[2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right]\hat{\phi}=zero \qquad 19.91$$

$$2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}=zero$$

That is equal to the derivative of the constant angular momentum $L=r^2\frac{d\phi}{dt}$ 19.92

$$\frac{dL}{dt}=\frac{d}{dt}\left(r^2\frac{d\phi}{dt}\right)=2r\frac{dr}{dt}\frac{d\phi}{dt}+r^2\frac{d^2\phi}{dt^2}=zero \qquad 19.93$$

Rewriting some equations already described we have:

$$w=\frac{1}{r}$$

$$dw=\frac{\partial w}{\partial r}dr\Rightarrow dw=\frac{-1}{r^2}dr$$

$\frac{dw}{d\phi}=\frac{-1}{r^2}\frac{dr}{d\phi}$ or $\frac{dr}{d\phi}=-r^2\frac{dw}{d\phi}$ and $\frac{dw}{dt}=\frac{-1}{r^2}\frac{dr}{dt}$

$$\frac{dr}{dt}=\frac{d\phi}{dt}\frac{dt}{d\phi}\frac{dr}{dt}=\frac{L}{r^2}\frac{dr}{d\phi}=\frac{-L}{r^2}r^2\frac{dw}{d\phi}\Rightarrow\frac{dr}{dt}=-L\frac{dw}{d\phi}$$

$$\frac{d^2r}{dt^2}=\frac{d}{dt}\left(\frac{dr}{dt}\right)=\frac{d\phi}{dt}\frac{dt}{d\phi}\frac{d}{dt}\left(-L\frac{dw}{d\phi}\right)=\frac{L}{r^2}\frac{d}{d\phi}\left(-L\frac{dw}{d\phi}\right)=\frac{-L^2}{r^2}\frac{d^2w}{d\phi^2} \qquad 19.94$$

From 19.90 we have:

$$\left(1-\frac{3u^2}{2c^2}\right)\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]=\frac{-GM_o}{r^2}$$

In this one we 19.94 the speed of 19.80 and the angular momentum we have:

$$\left[1-\frac{3}{2c^2}\left(\frac{2GM_o}{r}\right)\right]\left[\frac{-L^2}{r^2}\frac{d^2w}{d\phi^2}-r\left(\frac{L}{r^2}\right)^2\right]=-\frac{GM_o}{r^2}$$

$$\left(1-\frac{3GM_o}{c^2}\frac{1}{r}\right)\left(\frac{d^2w}{d\phi^2}+\frac{1}{r}\right)=\frac{GM_o}{L^2}$$

$$\left(1-\frac{3GM_o}{c^2}\frac{1}{r}\right)\frac{d^2w}{d\phi^2}+\left(1-\frac{3Gm_o}{c^2}\frac{1}{r}\right)\frac{1}{r}=\frac{GM_o}{L^2}$$

$$\frac{d^2w}{d\phi^2}-\frac{3GM_o}{c^2}\frac{d^2w}{d\phi^2}\frac{1}{r}+\frac{1}{r}-\frac{3GM_o}{c^2}\frac{1}{r^2}-\frac{GM_o}{L^2}=zero$$

$$\frac{d^2w}{d\phi^2}-A\frac{d^2w}{d\phi^2}\frac{1}{r}+\frac{1}{r}-A\frac{1}{r^2}-B=zero$$

$$\frac{d^2w}{d\phi^2}-A\frac{d^2w}{d\phi^2}w+w-Aw^2-B=zero$$

$$\frac{d^2w}{d\phi^2}-A\frac{d^2w}{d\phi^2}w-Aw^2+w-B=zero \quad 19.95$$

Where we have:

$$A=\frac{3GM_o}{c^2} \quad B=\frac{GM_o}{L^2} \quad 19.96$$

The solution to the differential equation 19.95 is:

$$w=\frac{1}{\varepsilon D}\left[1-\varepsilon\cos(\phi Q+\phi_o)\right]\Rightarrow w=\frac{1}{\varepsilon D}\left[1-\varepsilon\cos(\phi Q)\right]. \quad 19.97$$

Where we consider $\phi_o=zero$

Then the radius is given by:

$$r=\frac{1}{w}=\frac{\varepsilon D}{1-\varepsilon\cos(\phi Q)}\Rightarrow r=\frac{\varepsilon D}{1-\varepsilon\cos(\phi Q)} \quad 19.98$$

Where ε is the eccentricity and D the focus distance to the directory.

Deriving 19.97 we have $\frac{dw}{d\phi}=\frac{Q\,sen(\phi Q)}{D}$ and $\frac{d^2w}{d\phi^2}=\frac{Q^2\cos(\phi Q)}{D}$ 19.99

Applying the derivatives in 19.95 we have:

$$\frac{d^2w}{d\phi^2}-A\frac{d^2w}{d\phi^2}w-Aw^2+w-B=zero$$

$$\frac{Q^2\cos(\phi Q)}{D}-\frac{AQ^2\cos(\phi Q)}{D}\frac{1}{\varepsilon D}\left[1-\varepsilon\cos(\phi Q)\right]-\frac{A}{\varepsilon^2D^2}\left[1-\varepsilon\cos(\phi Q)\right]^2+\frac{1}{\varepsilon D}\left[1-\varepsilon\cos(\phi Q)\right]-B=zero$$

$$\frac{Q^2\cos(\phi Q)}{D}-\frac{AQ^2\cos(\phi Q)}{\varepsilon D^2}\left[1-\varepsilon\cos(\phi Q)\right]-\frac{A}{\varepsilon^2D^2}\left[1-2\varepsilon\cos(\phi Q)+\varepsilon^2\cos^2(\phi Q)\right]+\left[\frac{1}{\varepsilon D}-\frac{1}{\varepsilon D}\varepsilon\cos(\phi Q)\right]-B=zero$$

$$\frac{Q^2\cos(\phi Q)}{D}-\frac{AQ^2\cos(\phi Q)}{\varepsilon D^2}+\frac{AQ^2\cos(\phi Q)}{\varepsilon D^2}\varepsilon\cos(\phi Q)-$$

$$-\frac{A}{\varepsilon^2D^2}+\frac{A}{\varepsilon^2D^2}2\varepsilon\cos(\phi Q)-\frac{A}{\varepsilon^2D^2}\varepsilon^2\cos^2(\phi Q)+\frac{1}{\varepsilon D}-\frac{1}{\varepsilon D}\varepsilon\cos(\phi Q)-B=zero$$

$$\frac{\cos(\phi Q)}{D}\left(Q^2-\frac{AQ^2}{\varepsilon D}+\frac{2A}{\varepsilon D}-1\right)+\frac{AQ^2\cos^2(\phi Q)}{D^2}-\frac{A\cos^2(\phi Q)}{D^2}-\frac{A}{\varepsilon^2 D^2}+\frac{1}{\varepsilon D}-B=\text{zero}$$

$$\frac{\cos(\phi Q)}{AD}\left(Q^2-\frac{AQ^2}{\varepsilon D}+\frac{2A}{\varepsilon D}-1\right)+\frac{AQ^2\cos^2(\phi Q)}{AD^2}-\frac{A\cos^2(\phi Q)}{AD^2}-\frac{A}{A\varepsilon^2 D^2}+\frac{1}{A\varepsilon D}-\frac{B}{A}=\text{zero}$$

$$\frac{\cos(\phi Q)}{D}\left(\frac{Q^2}{A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}-\frac{1}{A}\right)+\frac{Q^2\cos^2(\phi Q)}{D^2}-\frac{\cos^2(\phi Q)}{D^2}-\frac{1}{\varepsilon^2 D^2}+\frac{1}{A\varepsilon D}-\frac{B}{A}=\text{zero}$$

$$\frac{\cos^2(\phi Q)}{D^2}\left(Q^2-1\right)+\frac{\cos(\phi Q)}{D}\left(\frac{Q^2}{A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}-\frac{1}{A}\right)-\frac{1}{\varepsilon^2 D^2}+\frac{1}{A\varepsilon D}-\frac{B}{A}=\text{zero} \qquad 19.100$$

The coefficient of the squared co-cosine can be considered null because $Q\approx 1$ and D^2 is a very large number:

$$\frac{\cos^2(\phi Q)}{D^2}\left(Q^2-1\right)=\text{zero} \qquad 19.101$$

Resulting from the equation 19.100:

$$\frac{\cos(\phi Q)}{D}\left(\frac{Q^2}{A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}-\frac{1}{A}\right)-\frac{1}{\varepsilon^2 D^2}+\frac{1}{A\varepsilon D}-\frac{B}{A}=\text{zero} \qquad 19.102$$

Due to the unicity of the equation 19.102 we must have the only solution that makes it null simultaneously the parenthesis and the rest of the equation, that is, we must have a unique solution for both the following equations:

$$\frac{Q^2}{A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}-\frac{1}{A}=\text{zero} \quad \text{and} \quad -\frac{1}{\varepsilon^2 D^2}+\frac{1}{A\varepsilon D}-\frac{B}{A}=\text{zero} \qquad 19.103$$

These equations can be written as:

$$[a=b]\Rightarrow\frac{1}{A}-\frac{1}{\varepsilon D}=\frac{1}{Q^2}\left(\frac{1}{A}-\frac{2}{\varepsilon D}\right) \qquad 19.104$$

$$[a=c]\Rightarrow\frac{1}{A}-\frac{1}{\varepsilon D}=\frac{\varepsilon DB}{A} \qquad 19.105$$

In these ones the common term $a=\frac{1}{A}-\frac{1}{\varepsilon D}$ must have a single solution then we have:

$$[b=c]\Rightarrow\frac{1}{Q^2}\left(\frac{1}{A}-\frac{2}{\varepsilon D}\right)=\frac{\varepsilon DB}{A} \qquad 19.106$$

With 19.96 and the theoretical momentum we have:

$$A=\frac{3GM_o}{c^2} \quad B=\frac{GM_o}{L^2} \quad L^2=\varepsilon DGM_o \qquad \varepsilon DB=\frac{\varepsilon DGM_o}{L^2}=1 \qquad 19.107$$

It is applied in 19.105 and 19.106 resulting in:

$$[a=c]\Rightarrow \frac{1}{A}-\frac{1}{\varepsilon D}=\frac{1}{A} \qquad 19.108$$

$$[b=c]\Rightarrow \frac{1}{Q^2}\left(\frac{1}{A}-\frac{2}{\varepsilon D}\right)=\frac{1}{A} \qquad 19.109$$

From 19.108 we have the mistake made in 19.105:

$$\frac{1}{A}-\frac{1}{\varepsilon D}=\frac{1}{A}\Rightarrow -\frac{1}{\varepsilon D}\approx \text{zero} \qquad 19.110$$

$$-\frac{1}{\varepsilon D}=\frac{-1}{55.442.955.600{,}00}=-1{,}80.10^{-11}\approx \text{zero} \qquad 19.111$$

From 19.109 we have Q:

$$\frac{1}{Q^2}\left(\frac{1}{A}-\frac{2}{\varepsilon D}\right)=\frac{1}{A}\Rightarrow Q^2=1-\frac{2A}{\varepsilon D}\Rightarrow Q^2=1-\frac{2}{\varepsilon D}\frac{3GM_\circ}{c^2} \qquad 19.112$$

It is applied in 19.104 resulting in 19.110:

$$\frac{1}{A}-\frac{1}{\varepsilon D}=\frac{1}{Q^2}\left(\frac{1}{A}-\frac{2}{\varepsilon D}\right)\Rightarrow \frac{1}{A}-\frac{1}{\varepsilon D}=\frac{1}{\left(1-\frac{2A}{\varepsilon D}\right)}\left(\frac{1}{A}-\frac{2}{\varepsilon D}\right)\Rightarrow \frac{1}{A}-\frac{1}{\varepsilon D}=\frac{1}{A}\Rightarrow -\frac{1}{\varepsilon D}\approx \text{zero}$$

From 19.112 we have:

$$Q=\sqrt{1-\frac{6GM_\circ}{\varepsilon Dc^2}}=\sqrt{1-\frac{6\left(6{,}67.10^{-11}\right)\left(1{,}98.10^{30}\right)}{\left(55.442.955.600{,}00\right)\left(3.10^{8}\right)^2}}=0{,}999.999.920.599 \qquad 19.113$$

That corresponds to the advance of Mercury’s perihelion in one century of:

$$\sum\Delta\phi=\Delta\phi.\,415{,}79=\left(\frac{1}{Q}-1\right).1.296.000{,}00.415{,}79=42{,}79'' \qquad 19.114$$

Calculated in this way:

In one trigonometric turn we have $360\times60\times60=1.296.000{,}00''$ seconds.

The angle ϕ in seconds ran by the planet in one trigonometric turn is given by:

$$\phi Q=1.296.000{,}00\Rightarrow \phi=\frac{1.296.000{,}00}{Q}.$$

If $Q>1{,}00$ we have a regression. $\phi<1.296.000{,}00$.

If $Q<1{,}00$ we have an advance. $\phi>1.296.000{,}00$.

The angular variation in seconds in one turn is given by:

$$\Delta\phi = \frac{1.296.000{,}00}{Q} - 1.296.000{,}00 = \left(\frac{1}{Q} - 1\right) 1.296.000{,}00.$$

If $\Delta\phi < zero$ we have a regression.

If $\Delta\phi > zero$ we have an advance.

In one century we have 415,79 turns that supply a total angular variation of:

$$\sum \Delta\phi = \Delta\phi .\, 415{,}79 = \left(\frac{1}{Q} - 1\right) .\, 1.296.000{,}00 .\, 415{,}79 = 42{,}79''$$

If $\sum \Delta\phi < zero$ we have a regression.

If $\sum \Delta\phi > zero$ we have an advance.

§20 Inertia

Imagine in an infinite universe totally empty, a point O' which is the beginning of the coordinates of the observer O'. In the cases of the observer O' being at rest or in uniform motion the law of inertia requires that the spherical electromagnetic waves with speed c issued by a source located at point O' is always observed by O', regardless of time, with spherical speed c and therefore the uniform motion and rest are indistinguishable from each other remain valid in both cases the law of inertia. To the observer O' the equations of electromagnetic theory describe the spread just like a spherical wave. The image of an object located in O' will always be centered on the object itself and a beam of light emitted from O' will always remain straight and perpendicular to the spherical waves.

Imagine another point O what will be the beginning of the coordinates of the observer which has the same properties as described for the inertial observer O'.

Obviously two imaginary points without any form of interaction between them remain individually and together perfectly meeting the law of inertia even though there is a uniform motion between them only detectable due to the presence of two observers who will be considered individually in rest, setting in motion the other referential.

The intrinsic properties of these two observers are described by the equations of relativistic transformations.

Note: the infinite universe is one in which any point can be considered the central point of this universe.

(§ 20 electronic translation)

§21 Advance of Mercury's perihelion of 42.79" calculated with the Undulating Relativity

Assuming $ux=v$

(2.3) $u'x'=\dfrac{ux-v}{\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vux}{c^2}}}=\dfrac{v-v}{\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vv}{c^2}}}\Rightarrow u'x'=zero$

$ux=v$ $\qquad\qquad u'x'=zero$ 21.01

(1.17) $dt'=dt\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vux}{c^2}}=dt\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vv}{c^2}}\Rightarrow dt'=dt\sqrt{1-\dfrac{v^2}{c^2}}$

(1.22) $dt=dt'\sqrt{1+\dfrac{v'^2}{c^2}+\dfrac{2v'u'x'}{c^2}}=dt'\sqrt{1+\dfrac{v'^2}{c^2}+\dfrac{2v'(0)}{c^2}}\Rightarrow dt=dt'\sqrt{1+\dfrac{v'^2}{c^2}}$

$dt'=dt\sqrt{1-\dfrac{v^2}{c^2}}$ $\qquad\qquad dt=dt'\sqrt{1+\dfrac{v'^2}{c^2}}$ 21.02

$\sqrt{1-\dfrac{v^2}{c^2}}\sqrt{1+\dfrac{v'^2}{c^2}}=1$ 21.03

$v=\dfrac{v'}{\sqrt{1+\dfrac{v'^2}{c^2}}}$ $\qquad\qquad v'=\dfrac{v}{\sqrt{1-\dfrac{v^2}{c^2}}}$ 21.04

$dt>dt'$ $\qquad v<v'$ $\qquad vdt=v'dt'$ 21.05

(1.33) $\vec{v}=\dfrac{-\vec{v}'}{\sqrt{1+\dfrac{v'^2}{c^2}+\dfrac{2v'u'x'}{c^2}}}=\dfrac{-\vec{v}'}{\sqrt{1+\dfrac{v'^2}{c^2}+\dfrac{2v'(0)}{c^2}}}\Rightarrow\vec{v}=\dfrac{-\vec{v}'}{\sqrt{1+\dfrac{v'^2}{c^2}}}$

(1.34) $\vec{v}'=\dfrac{-\vec{v}}{\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vux}{c^2}}}=\dfrac{-\vec{v}}{\sqrt{1+\dfrac{v^2}{c^2}-\dfrac{2vv}{c^2}}}\Rightarrow\vec{v}'=\dfrac{-\vec{v}}{\sqrt{1-\dfrac{v^2}{c^2}}}$

$\vec{v}=\dfrac{-\vec{v}'}{\sqrt{1+\dfrac{v'^2}{c^2}}}$ $\qquad\qquad -\vec{v}'=\dfrac{\vec{v}}{\sqrt{1-\dfrac{v^2}{c^2}}}$ 21.06

$\vec{r}=r\hat{r}=-\vec{r}'$ $\qquad \vec{r}'=-r\hat{r}=-\vec{r}$ $\qquad |\vec{r}|=|\vec{r}'|=r$ 21.07

$d\vec{r}=dr\hat{r}+rd\hat{r}=-d\vec{r}'$ $\qquad d\vec{r}'=-dr\hat{r}-rd\hat{r}=-d\vec{r}$ 21.08

$\hat{r}d\vec{r}=dr\hat{r}\hat{r}+r\hat{r}d\hat{r}=dr$ $\qquad \hat{r}d\vec{r}'=-dr\hat{r}\hat{r}-r\hat{r}d\hat{r}=-dr$ 21.09

$\vec{v}=\dfrac{d\vec{r}}{dt}=\dfrac{d(r\hat{r})}{dt}=\dfrac{dr}{dt}\hat{r}+r\dfrac{d\phi}{dt}\hat{\phi}$ $\qquad v^2=\vec{v}\vec{v}=\left(\dfrac{dr}{dt}\right)^2+\left(r\dfrac{d\phi}{dt}\right)^2$ 21.10

$\vec{v}'=\dfrac{d\vec{r}'}{dt'}=\dfrac{d(-r\hat{r})}{dt'}=-\left(\dfrac{dr}{dt'}\hat{r}+r\dfrac{d\phi}{dt'}\hat{\phi}\right)$ $\qquad v'^2=\vec{v}'\vec{v}'=\left(\dfrac{dr}{dt'}\right)^2+\left(r\dfrac{d\phi}{dt'}\right)^2$ 21.11

$$\vec{a}=\frac{d\vec{v}}{dt}=\frac{d^2\vec{r}}{dt^2}=\frac{d^2(r\hat{r})}{dt^2}=\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}+\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\hat{\phi} \qquad 21.12$$

$$\vec{a}'=\frac{d\vec{v}'}{dt'}=\frac{d^2\vec{r}'}{dt'^2}=\frac{d^2(-r\hat{r})}{dt'^2}=-\left[\frac{d^2r}{dt'^2}-r\left(\frac{d\phi}{dt'}\right)^2\right]\hat{r}-\left(2\frac{dr}{dt'}\frac{d\phi}{dt'}+r\frac{d^2\phi}{dt'^2}\right)\hat{\phi} \qquad 21.13$$

$$-\vec{v}'=\frac{\vec{v}}{\sqrt{1-\frac{v^2}{c^2}}} \qquad 21.06$$

$$-\vec{a}'=\frac{d(-\vec{v}')}{dt'}=\frac{d}{dt'}\left(\frac{\vec{v}}{\sqrt{1-\frac{v^2}{c^2}}}\right)=\frac{dt}{dt'}\frac{d}{dt}\left(\frac{\vec{v}}{\sqrt{1-\frac{v^2}{c^2}}}\right)=\sqrt{1+\frac{v'^2}{c^2}}\frac{d}{dt}\left(\frac{\vec{v}}{\sqrt{1-\frac{v^2}{c^2}}}\right) \qquad 21.14$$

$$-\vec{a}'=-\frac{d\vec{v}'}{dt'}=\sqrt{1+\frac{v'^2}{c^2}}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}\left[\sqrt{1-\frac{v^2}{c^2}}\frac{d\vec{v}}{dt}-\vec{v}\frac{d}{dt}\left(\sqrt{1-\frac{v^2}{c^2}}\right)\right]$$

$$-\vec{a}'=-\frac{d\vec{v}'}{dt'}=\sqrt{1+\frac{v'^2}{c^2}}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}\left[\sqrt{1-\frac{v^2}{c^2}}\frac{d\vec{v}}{dt}-\vec{v}\frac{1}{2}\left(1-\frac{v^2}{c^2}\right)^{\frac{1}{2}-\frac{2}{2}=\frac{-1}{2}}\left(\frac{-2v}{c^2}\frac{dv}{dt}\right)\right]$$

$$-\vec{a}'=-\frac{d\vec{v}'}{dt'}=\sqrt{1+\frac{v'^2}{c^2}}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}\left(\sqrt{1-\frac{v^2}{c^2}}\frac{d\vec{v}}{dt}+\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right)$$

$$-\vec{a}'=-\frac{d\vec{v}'}{dt'}=\sqrt{1+\frac{v'^2}{c^2}}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}\left(\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}\sqrt{1-\frac{v^2}{c^2}}\sqrt{1-\frac{v^2}{c^2}}\frac{d\vec{v}}{dt}+\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right)$$

$$-\vec{a}'=-\frac{d\vec{v}'}{dt'}=\sqrt{1+\frac{v'^2}{c^2}}\frac{1}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right]$$

$$-m'\vec{a}'=\frac{-m_o\vec{a}'}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{-m_o}{\sqrt{1+\frac{v'^2}{c^2}}}\frac{d\vec{v}'}{dt'}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right]$$

$$\vec{F}'=-m'\vec{a}'=\frac{-m_o\vec{a}'}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{-m_o}{\sqrt{1+\frac{v'^2}{c^2}}}\frac{d\vec{v}'}{dt'} \qquad 21.15$$

$$\vec{F}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right] \qquad (=19.06) \tag{21.16}$$

$$\vec{F}'=-m'\vec{a}'=\frac{-m_o\vec{a}'}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{-m_o}{\sqrt{1+\frac{v'^2}{c^2}}}\frac{d\vec{v}'}{dt'}=\vec{F}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right] \tag{21.17}$$

$$E_k=\int\vec{F}'(-d\vec{r}')=\int\vec{F}.d\vec{r}=\int\frac{-k}{r^2}\hat{r}d\vec{r} \tag{21.18}$$

$$E_k=\int\vec{F}'\,.(-d\vec{r}')=\int\vec{F}.d\vec{r}=\int\frac{-m_o}{\sqrt{1+\frac{v'^2}{c^2}}}\frac{d\vec{v}'}{dt'}(-d\vec{r}')=\int\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right]d\vec{r}=\int\frac{-k}{r^2}\hat{r}.d\vec{r} \tag{21.19}$$

$$E_k=\int\frac{m_o}{\sqrt{1+\frac{v'^2}{c^2}}}d\vec{v}'\frac{d\vec{r}'}{dt'}=\int\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)d\vec{v}\frac{d\vec{r}}{dt}+vdv\frac{d\vec{r}}{dt}\frac{\vec{v}}{c^2}\right]=\int\frac{-k}{r^2}\hat{r}.d\vec{r}$$

$$E_k=\int\frac{m_o d\vec{v}'\vec{v}'}{\sqrt{1+\frac{v'^2}{c^2}}}=\int\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)d\vec{v}\vec{v}+vdv\frac{\vec{v}\vec{v}}{c^2}\right]=\int\frac{-k}{r^2}\hat{r}.d\vec{r}$$

$$E_k=\int\frac{m_o v'dv'}{\sqrt{1+\frac{v'^2}{c^2}}}=\int\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)vdv+vdv\frac{v^2}{c^2}\right]=\int\frac{-k}{r^2}dr$$

$$E_k=\int\frac{m_o v'dv'}{\sqrt{1+\frac{v'^2}{c^2}}}=\int\frac{m_o vdv}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left(1-\frac{v^2}{c^2}+\frac{v^2}{c^2}\right)=\int\frac{-k}{r^2}dr$$

$$E_k=\int\frac{m_o v'dv'}{\sqrt{1+\frac{v'^2}{c^2}}}=\int\frac{m_o vdv}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}=\int\frac{-k}{r^2}dr \qquad dE_k=\frac{m_o v'dv'}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{m_o vdv}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}=\frac{-k}{r^2}dr \tag{21.20}$$

$$E_k=m_o c^2\sqrt{1+\frac{v'^2}{c^2}}=\frac{m_o c^2}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{k}{r}+constant \tag{21.21}$$

$$E_R=m_o c^2\sqrt{1+\frac{v'^2}{c^2}}-\frac{k}{r}=constant \qquad E_R=\frac{m_o c^2}{\sqrt{1-\frac{v^2}{c^2}}}-\frac{k}{r}=constant \tag{21.22}$$

$$E_R=\frac{m_o c^2}{\sqrt{1-\frac{v^2}{c^2}}}-\frac{k}{r}=m_o c^2+\frac{m_o v^2}{2}-\frac{k}{r} \qquad E_R=\frac{m_o c^2}{\sqrt{1-\frac{(0)^2}{c^2}}}-\frac{k}{\infty}=m_o c^2 \tag{21.23}$$

$$\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{E_R}{m_o c^2}+\frac{k}{m_o c^2}\frac{1}{r} \qquad H=\frac{E_R}{m_o c^2} \qquad A=\frac{k}{m_o c^2}=\frac{GM_o m_o}{m_o c^2}=\frac{GM_o}{c^2} \tag{21.24}$$

$$\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}=H+A\frac{1}{r} \qquad \frac{1}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}=\left(H+A\frac{1}{r}\right)^3 \tag{21.25}$$

$$\vec{L}=\vec{r}\times\vec{v}=r\hat{r}\times\left(\frac{dr}{dt}\hat{r}+r\frac{d\phi}{dt}\hat{\phi}\right)=r^2\frac{d\phi}{dt}\left(\hat{r}\times\hat{\phi}\right)=r^2\frac{d\phi}{dt}\hat{k} \tag{21.26}$$

$$\vec{L}=\vec{r}\times\vec{v}=\vec{r}\times\frac{-\vec{v}'}{\sqrt{1+\frac{v'^2}{c^2}}}=r\hat{r}\times\frac{-1}{\sqrt{1+\frac{v'^2}{c^2}}}\left[-\left(\frac{dr}{dt'}\hat{r}+r\frac{d\phi}{dt'}\hat{\phi}\right)\right]=\frac{1}{\sqrt{1+\frac{v'^2}{c^2}}}r^2\frac{d\phi}{dt'}\left(\hat{r}\times\hat{\phi}\right)=r^2\frac{d\phi}{dt}\hat{k} \tag{21.26}$$

$$\vec{L}=r^2\frac{d\phi}{dt}\vec{k}=L\hat{k}=\text{constant} \qquad L=r^2\frac{d\phi}{dt} \tag{21.27}$$

$$dE_k=\frac{m_o v'dv'}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{m_o v dv}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}=\frac{-k}{r^2}dr=\frac{-k}{r^2}\hat{r}.d\vec{r} \tag{21.20}$$

$$\frac{dE_k}{dt}=\vec{F}\vec{v}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\vec{v}\frac{d\vec{v}}{dt}=\frac{-k}{r^2}\hat{r}.\frac{d\vec{r}}{dt}=\frac{-k}{r^2}\hat{r}.\vec{v}$$

$$\vec{F}=\frac{m_o\vec{a}}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}=\frac{-k}{r^2}\hat{r} \tag{21.28}$$

$$\vec{F}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left\{\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}+\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\hat{\phi}\right\}=\frac{-k}{r^2}\hat{r} \tag{21.29}$$

$$\vec{F}_{\hat{\phi}}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left(2\frac{dr}{dt}\frac{d\phi}{dt}+r\frac{d^2\phi}{dt^2}\right)\hat{\phi}=zero \tag{21.30}$$

$$\vec{F}_{\hat{r}}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\frac{d^2r}{dt^2}-r\left(\frac{d\phi}{dt}\right)^2\right]\hat{r}=\frac{-k}{r^2}\hat{r} \tag{21.31}$$

$$\frac{d\phi}{dt}=\frac{L}{r^2} \qquad \frac{dr}{dt}=-L\frac{dw}{d\phi} \qquad \frac{d^2r}{dt^2}=\frac{-L^2}{r^2}\frac{d^2w}{d\phi^2} \qquad \frac{d^2\phi}{dt^2}=\frac{2L^2}{r^3}\frac{dw}{d\phi} \tag{21.32}$$

$$\vec{F}_{\hat{r}} = \frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\frac{-L^2}{r^2}\frac{d^2w}{d\phi^2} - r\left(\frac{L}{r^2}\right)^2\right]\hat{r} = \frac{-k}{r^2}\hat{r} \qquad 21.33$$

$$\frac{1}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left(\frac{-L^2}{r^2}\frac{d^2w}{d\phi^2} - \frac{L^2}{r^3}\right) = \frac{-GM_o}{r^2}$$

$$\frac{1}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left(\frac{d^2w}{d\phi^2} + \frac{1}{r}\right)\left(\frac{-L^2}{r^2}\right) = \frac{-GM_o}{r^2}$$

$$\frac{1}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left(\frac{d^2w}{d\phi^2} + \frac{1}{r}\right) = \frac{GM_o}{L^2} \qquad 21.34$$

$$\left(H + A\frac{1}{r}\right)^3\left(\frac{d^2w}{d\phi^2} + \frac{1}{r}\right) = \frac{GM_o}{L^2} \qquad 21.35$$

$$\left(H + 3A\frac{1}{r}\right)\left(\frac{d^2w}{d\phi^2} + \frac{1}{r}\right) = \frac{GM_o}{L^2}$$

$$H\frac{d^2w}{d\phi^2} + H\frac{1}{r} + 3A\frac{d^2w}{d\phi^2}\frac{1}{r} + 3A\frac{1}{r^2} = \frac{GM_o}{L^2}$$

$$H\frac{d^2w}{d\phi^2} + Hw + 3A\frac{d^2w}{d\phi^2}w + 3Aw^2 - \frac{GM_o}{L^2} = zero$$

$$H = \frac{E_R}{m_o c^2} \qquad A = \frac{k}{m_o c^2} = \frac{GM_o m_o}{m_o c^2} = \frac{GM_o}{c^2} \qquad B = \frac{GM_o}{L^2} \qquad 21.36$$

$$H\frac{d^2w}{d\phi^2} + Hw + 3A\frac{d^2w}{d\phi^2}w + 3Aw^2 - B = zero \qquad 21.37$$

$$w = \frac{1}{r} = \frac{1}{\varepsilon D}\left[1 + \varepsilon\cos(\phi Q)\right] \qquad \frac{dw}{d\phi} = \frac{-Q\,sen(\phi Q)}{D} \qquad \frac{d^2w}{d\phi^2} = \frac{-Q^2\cos(\phi Q)}{D} \qquad 21.38$$

$$H\frac{-Q^2\cos(\phi Q)}{D} + H\frac{1}{\varepsilon D}\left[1 + \varepsilon\cos(\phi Q)\right] + 3A\frac{-Q^2\cos(\phi Q)}{D}\frac{1}{\varepsilon D}\left[1 + \varepsilon\cos(\phi Q)\right] + 3A\left\{\frac{1}{\varepsilon D}\left[1 + \varepsilon\cos(\phi Q)\right]\right\}^2 - B = zero \qquad 21.39$$

$$-Q^2H\frac{\cos(\phi Q)}{D} + H\frac{1}{\varepsilon D} + H\frac{1}{\varepsilon D}\varepsilon\cos(\phi Q) - \frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}\left[1 + \varepsilon\cos(\phi Q)\right] + \frac{3A}{\varepsilon^2 D^2}\left[1 + 2\varepsilon\cos(\phi Q) + \varepsilon^2\cos^2(\phi Q)\right] - B = zero$$

$$-Q^2H\frac{\cos(\phi Q)}{D} + H\frac{1}{\varepsilon D} + H\frac{\cos(\phi Q)}{D} - \frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D} - \frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}\varepsilon\cos(\phi Q) +$$
$$+ \frac{3A}{\varepsilon^2 D^2} + \frac{3A}{\varepsilon^2 D^2}2\varepsilon\cos(\phi Q) + \frac{3A}{\varepsilon^2 D^2}\varepsilon^2\cos^2(\phi Q) - B = zero$$

$$-Q^2H\frac{\cos(\phi Q)}{D}+H\frac{1}{\varepsilon D}+H\frac{\cos(\phi Q)}{D}-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}-3Q^2A\frac{\cos^2(\phi Q)}{D^2}+$$
$$+\frac{3A}{\varepsilon^2D^2}+\frac{6A}{\varepsilon D}\frac{\cos(\phi Q)}{D}+3A\frac{\cos^2(\phi Q)}{D^2}-B=zero$$

$$-Q^2H\frac{\cos(\phi Q)}{D}+H\frac{\cos(\phi Q)}{D}-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}+\frac{6A}{\varepsilon D}\frac{\cos(\phi Q)}{D}-$$
$$-3Q^2A\frac{\cos^2(\phi Q)}{D^2}+3A\frac{\cos^2(\phi Q)}{D^2}+H\frac{1}{\varepsilon D}+\frac{3A}{\varepsilon^2D^2}-B=zero$$

$$\left(-Q^2H+H-\frac{3Q^2A}{\varepsilon D}+\frac{6A}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\left(-3Q^2A+3A\right)\frac{\cos^2(\phi Q)}{D^2}+H\frac{1}{\varepsilon D}+\frac{3A}{\varepsilon^2D^2}-B=zero$$

$$\left(-3Q^2A+3A\right)\frac{\cos^2(\phi Q)}{3AD^2}+\left(-Q^2H+H-\frac{3Q^2A}{\varepsilon D}+\frac{6A}{\varepsilon D}\right)\frac{\cos(\phi Q)}{3AD}+H\frac{1}{3A\varepsilon D}+\frac{3A}{3A\varepsilon^2D^2}-\frac{B}{3A}=zero$$

$$\left(1-Q^2\right)\frac{\cos^2(\phi Q)}{D^2}+\left(\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}-\frac{B}{3A}=zero \qquad 21.40$$

$$Q^2\approx 1 \qquad \left(1-Q^2\right)\frac{\cos^2(\phi Q)}{D^2}=zero \qquad 21.41$$

$$\left(\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}-\frac{B}{3A}=zero \qquad 21.42$$

$$\frac{\cos(\phi Q)}{D}=zero\Rightarrow\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}-\frac{B}{3A}=zero$$

$$\frac{\cos(\phi Q)}{D}\neq zero\Rightarrow\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}=zero$$

$$\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}=zero \qquad \frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}-\frac{B}{3A}=zero \qquad 21.43$$

$$[a=b]\Rightarrow\frac{H}{3A}+\frac{1}{\varepsilon D}=\frac{1}{Q^2}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right) \qquad [a=c]\Rightarrow\frac{H}{3A}+\frac{1}{\varepsilon D}=\frac{\varepsilon DB}{3A} \qquad 21.44$$

$$Q^2=1 \qquad H=\frac{E_R}{m_oc^2}=\frac{m_oc^2}{m_oc^2}=1 \qquad \varepsilon DB=\frac{\varepsilon DGM_o}{L^2}=\frac{\varepsilon DGM_o}{\varepsilon DGM_o}=1$$

$$[a=b]\Rightarrow\frac{H}{3A}+\frac{1}{\varepsilon D}=\frac{1}{1}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right)\Rightarrow\frac{1}{\varepsilon D}=zero \qquad [a=c]\Rightarrow\frac{1}{3A}+\frac{1}{\varepsilon D}=\frac{1}{3A}\Rightarrow\frac{1}{\varepsilon D}=zero$$

$$[b=c]\Rightarrow\frac{1}{Q^2}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right)=\frac{\varepsilon DB}{3A} \qquad 21.45$$

$$\varepsilon DB=\frac{\varepsilon DGM_o}{L^2}=\frac{\varepsilon DGM_o}{\varepsilon DGM_o}=1 \qquad 21.46$$

$$[b=c] \Rightarrow \frac{1}{Q^2}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right)=\frac{1}{3A} \qquad Q^2=H+\frac{6A}{\varepsilon D} \qquad 21.47$$

$Q=Q(H)$ The regression is a function of positive energy that governs the movement.

$$H=\frac{E_R}{m_o c^2}=\frac{m_o c^2}{m_o c^2}=1 \qquad Q^2=1+\frac{6A}{\varepsilon D}\ \text{Regression} \qquad 21.48$$

$$[a=b] \Rightarrow \frac{1}{3A}+\frac{1}{\varepsilon D}=\frac{1}{\left(1+\frac{6A}{\varepsilon D}\right)}\left(\frac{1}{3A}+\frac{2}{\varepsilon D}\right) \Rightarrow \frac{1}{\varepsilon D}=zero \qquad 21.49$$

$$3A\varepsilon D\left(\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)=zero \qquad 3A\varepsilon^2D^2\left(\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}-\frac{B}{3A}\right)=zero \qquad 21.43$$

$$H=\frac{E_R}{m_o c^2} \qquad A=\frac{GM_o}{c^2} \qquad B=\frac{GM_o}{L^2}$$

$$-Q^2H\varepsilon D+H\varepsilon D-Q^2 3A+6A=zero \qquad H\varepsilon D+3A-\varepsilon D(\varepsilon DB)=zero$$

$$-Q^2(-3A+\varepsilon D)-3A+\varepsilon D-Q^2 3A+6A=zero \qquad H\varepsilon D=-3A+\varepsilon D$$

$$Q^2 3A-Q^2\varepsilon D+\varepsilon D-Q^2 3A+3A=zero$$

$$-Q^2\varepsilon D+\varepsilon D+3A=zero \qquad Q^2=1+\frac{3A}{\varepsilon D}$$

This regression is not governed by the positive energy

$$\vec{v}=\frac{-\vec{v}'}{\sqrt{1+\frac{v'^2}{c^2}}} \quad 21.06$$

$$\vec{a}=\frac{d\vec{v}}{dt}=\frac{d}{dt}\left(\frac{-\vec{v}'}{\sqrt{1+\frac{v'^2}{c^2}}}\right)=\frac{dt'}{dt}\frac{d}{dt'}\left(\frac{-\vec{v}'}{\sqrt{1+\frac{v'^2}{c^2}}}\right)=\sqrt{1-\frac{v^2}{c^2}}\frac{d}{dt'}\left(\frac{-\vec{v}'}{\sqrt{1+\frac{v'^2}{c^2}}}\right) \quad 21.50$$

$$\vec{a}=\frac{d\vec{v}}{dt}=\sqrt{1-\frac{v^2}{c^2}}\frac{-1}{\left(1+\frac{v'^2}{c^2}\right)}\left[\sqrt{1+\frac{v'^2}{c^2}}\frac{d\vec{v}'}{dt'}-\vec{v}'\frac{d}{dt'}\left(\sqrt{1+\frac{v'^2}{c^2}}\right)\right]$$

$$\vec{a}=\frac{d\vec{v}}{dt}=\sqrt{1-\frac{v^2}{c^2}}\frac{-1}{\left(1+\frac{v'^2}{c^2}\right)}\left[\sqrt{1+\frac{v'^2}{c^2}}\frac{d\vec{v}'}{dt'}-\vec{v}'\frac{1}{2}\left(1+\frac{v'^2}{c^2}\right)^{\frac{1}{2}-\frac{2}{2}=\frac{-1}{2}}\left(\frac{2v'}{c^2}\frac{dv'}{dt'}\right)\right]$$

$$\vec{a}=\frac{d\vec{v}}{dt}=\sqrt{1-\frac{v^2}{c^2}}\frac{-1}{\left(1+\frac{v'^2}{c^2}\right)}\left(\sqrt{1+\frac{v'^2}{c^2}}\frac{d\vec{v}'}{dt'}-\frac{1}{\sqrt{1+\frac{v'^2}{c^2}}}v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right)$$

$$\vec{a}=\frac{d\vec{v}}{dt}=\sqrt{1-\frac{v^2}{c^2}}\frac{-1}{\left(1+\frac{v'^2}{c^2}\right)}\left(\frac{1}{\sqrt{1+\frac{v'^2}{c^2}}}\sqrt{1+\frac{v'^2}{c^2}}\frac{d\vec{v}'}{dt'}\sqrt{1+\frac{v'^2}{c^2}}-\frac{1}{\sqrt{1+\frac{v'^2}{c^2}}}v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right)$$

$$\vec{a}=\frac{d\vec{v}}{dt}=\sqrt{1-\frac{v^2}{c^2}}\frac{-1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)\frac{d\vec{v}'}{dt'}-v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right]$$

$$m\vec{a}=\frac{m_o\vec{a}}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}\frac{d\vec{v}}{dt}=\frac{-m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)\frac{d\vec{v}'}{dt'}-v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right]$$

$$\vec{F}=m\vec{a}=\frac{m_o\vec{a}}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}\frac{d\vec{v}}{dt} \quad 21.51$$

$$\vec{F}'=\frac{-m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)\frac{d\vec{v}'}{dt'}-v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right] \quad 21.52$$

$$\vec{F}=m\vec{a}=\frac{m_o\vec{a}}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}\frac{d\vec{v}}{dt}=\vec{F}'=\frac{-m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)\frac{d\vec{v}'}{dt'}-v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right] \quad 21.53$$

$$E_k=\int\vec{F}.d\vec{r}=\int\vec{F}'\ .(-d\vec{r}')=\int\frac{-k}{r^2}\hat{r}(-d\vec{r}')\tag{21.54}$$

$$E_k=\int\vec{F}.d\vec{r}=\int\vec{F}'\ .(-d\vec{r}')=\int\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}\frac{d\vec{v}}{dt}d\vec{r}=\int\frac{-m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)\frac{d\vec{v}'}{dt'}-v'\frac{dv'}{dt'}\frac{\vec{v}'}{c^2}\right](-d\vec{r}')=\int\frac{-k}{r^2}\hat{r}(-d\vec{r}')\tag{21.55}$$

$$E_k=\int\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}d\vec{v}\frac{d\vec{r}}{dt}=\int\frac{m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)d\vec{v}'\frac{d\vec{r}'}{dt'}-v'dv'\frac{d\vec{r}'}{dt'}\frac{\vec{v}'}{c^2}\right]=\int\frac{k}{r^2}\hat{r}d\vec{r}'$$

$$E_k=\int\frac{m_o}{\sqrt{1-\frac{v^2}{c^2}}}d\vec{v}\vec{v}=\int\frac{m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)d\vec{v}'\vec{v}'-v'dv'\frac{\vec{v}'\vec{v}'}{c^2}\right]=\int\frac{-k}{r^2}dr$$

$$E_k=\int\frac{m_ovdv}{\sqrt{1-\frac{v^2}{c^2}}}=\int\frac{m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1+\frac{v'^2}{c^2}\right)d\vec{v}'\vec{v}'-v'dv'\frac{v'^2}{c^2}\right]=\int\frac{-k}{r^2}dr$$

$$E_k=\int\frac{m_ovdv}{\sqrt{1-\frac{v^2}{c^2}}}=\int\frac{m_ov'dv'}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left(1+\frac{v'^2}{c^2}-\frac{v'^2}{c^2}\right)=\int\frac{-k}{r^2}dr$$

$$E_k=\int\frac{m_ovdv}{\sqrt{1-\frac{v^2}{c^2}}}=\int\frac{m_ov'dv'}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=\int\frac{-k}{r^2}dr\qquad dE_k=\frac{m_ovdv}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_ov'dv'}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=\frac{-k}{r^2}dr\tag{21.56}$$

$$E_k=-m_oc^2\sqrt{1-\frac{v^2}{c^2}}=\frac{-m_oc^2}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{k}{r}+constant\tag{21.57}$$

$$E_R=-m_oc^2\sqrt{1-\frac{v^2}{c^2}}-\frac{k}{r}=constant\qquad E_R=\frac{-m_oc^2}{\sqrt{1+\frac{v'^2}{c^2}}}-\frac{k}{r}=constant\tag{21.58}$$

$$E_R=\frac{-m_oc^2}{\sqrt{1+\frac{v'^2}{c^2}}}-\frac{k}{r}=-m_oc^2+\frac{m_ov'^2}{2}-\frac{k}{r}\qquad E_R=\frac{-m_oc^2}{\sqrt{1+\frac{(0)^2}{c^2}}}-\frac{k}{\infty}=-m_oc^2\tag{21.59}$$

$$\frac{-1}{\sqrt{1+\frac{v'^2}{c^2}}}=\frac{E_R}{m_oc^2}+\frac{k}{m_oc^2}\frac{1}{r}\tag{21.60}$$

$$H=\frac{E_R}{m_oc^2}\qquad A=\frac{k}{m_oc^2}=\frac{GM_om_o}{m_oc^2}=\frac{GM_o}{c^2}\tag{21.61}$$

$$\frac{-1}{\sqrt{1+\frac{v'^2}{c^2}}}=H+A\frac{1}{r} \qquad \frac{1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=-\left(H+A\frac{1}{r}\right)^3 \tag{21.62}$$

$$\vec{L}'=\vec{r}'\times\vec{v}'=-r\hat{r}\times\left[-\left(\frac{dr}{dt'}\hat{r}+r\frac{d\phi}{dt'}\hat{\phi}\right)\right]=r^2\frac{d\phi}{dt'}\left(\hat{r}\times\hat{\phi}\right)=r^2\frac{d\phi}{dt'}\hat{k} \tag{21.63}$$

$$\vec{L}'=\vec{r}'x\vec{v}'=-r\hat{r}\times\frac{-\vec{v}}{\sqrt{1-\frac{v^2}{c^2}}}=r\hat{r}\times\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}\left[\left(\frac{dr}{dt}\hat{r}+r\frac{d\phi}{dt}\hat{\phi}\right)\right]=\frac{1}{\sqrt{1-\frac{v^2}{c^2}}}r^2\frac{d\phi}{dt}\left(\hat{r}\times\hat{\phi}\right)=r^2\frac{d\phi}{dt'}\hat{k} \tag{21.63}$$

$$\vec{L}'=r^2\frac{d\phi}{dt'}\vec{k}=L'\hat{k} \qquad L'=r^2\frac{d\phi}{dt'} \tag{21.64}$$

$$dE_k=\frac{m_o v dv}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_o v' dv'}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=\frac{-k}{r^2}dr=\frac{k}{r^2}\hat{r}d\vec{r}' \tag{21.56}$$

$$\frac{dE_k}{dt'}=\vec{F}'\vec{v}'=\frac{m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}v'\frac{dv'}{dt'}=\frac{k}{r^2}\hat{r}\frac{d\vec{r}'}{dt'}=\frac{k}{r^2}\hat{r}\vec{v}'$$

$$\vec{F}'=\frac{m_o\vec{a}'}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=\frac{k}{r^2}\hat{r} \tag{21.65}$$

$$\vec{F}'=\frac{m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left\{-\left[\frac{d^2r}{dt'^2}-r\left(\frac{d\phi}{dt'}\right)^2\right]\hat{r}-\left(2\frac{dr}{dt'}\frac{d\phi}{dt'}+r\frac{d^2\phi}{dt'^2}\right)\hat{\phi}\right\}=\frac{k}{r^2}\hat{r} \tag{21.66}$$

$$\vec{F}'_{\hat{\phi}}=\frac{-m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left(2\frac{dr}{dt'}\frac{d\phi}{dt'}+r\frac{d^2\phi}{dt'^2}\right)\hat{\phi}=zero \tag{21.67}$$

$$\vec{F}'_{\hat{r}}=\frac{-m_o}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\frac{d^2r}{dt'^2}-r\left(\frac{d\phi}{dt'}\right)^2\right]\hat{r}=\frac{k}{r^2}\hat{r} \tag{21.68}$$

$$\frac{1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\frac{d^2r}{dt'^2}-r\left(\frac{d\phi}{dt'}\right)^2\right]\hat{r}=\frac{-GM_o}{r^2}\hat{r}$$

$$\frac{d\phi}{dt'}=\frac{L'}{r^2} \qquad \frac{dr}{dt'}=-L'\frac{dw}{d\phi} \qquad \frac{d^2r}{dt'^2}=\frac{-L'^2}{r^2}\frac{d^2w}{d\phi^2} \qquad \frac{d^2\phi}{dt'^2}=\frac{2L'^2}{r^3}\frac{dw}{d\phi} \tag{21.69}$$

$$\frac{1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left[\frac{-L'^2}{r^2}\frac{d^2w}{d\phi^2}-r\left(\frac{L'}{r^2}\right)^2\right]=\frac{-GM_o}{r^2}$$

$$\frac{1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left(\frac{-L'^2}{r^2}\frac{d^2w}{d\phi^2}-\frac{L'^2}{r^3}\right)=\frac{-GM_o}{r^2}$$

$$\frac{1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left(\frac{d^2w}{d\phi^2}+\frac{1}{r}\right)\left(\frac{-L'^2}{r^2}\right)=\frac{-GM_o}{r^2}$$

$$\frac{1}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}\left(\frac{d^2w}{d\phi^2}+\frac{1}{r}\right)=\frac{GM_o}{L'^2} \qquad 21.70$$

$$-\left(H+A\frac{1}{r}\right)^3\left(\frac{d^2w}{d\phi^2}+\frac{1}{r}\right)=\frac{GM_o}{L'^2} \qquad 21.71$$

$$\left(H+3A\frac{1}{r}\right)\left(\frac{d^2w}{d\phi^2}+\frac{1}{r}\right)=\frac{-GM_o}{L'^2}$$

$$H\frac{d^2w}{d\phi^2}+H\frac{1}{r}+3A\frac{d^2w}{d\phi^2}\frac{1}{r}+3A\frac{1}{r^2}=\frac{-GM_o}{L'^2}$$

$$H\frac{d^2w}{d\phi^2}+Hw+3A\frac{d^2w}{d\phi^2}w+3Aw^2+\frac{GM_o}{L'^2}=zero$$

$$H=\frac{E_R}{m_o c^2} \qquad A=\frac{GM_o}{c^2} \qquad B=\frac{GM_o}{L'^2} \qquad 21.72$$

$$H\frac{d^2w}{d\phi^2}+Hw+3A\frac{d^2w}{d\phi^2}w+3Aw^2+B=zero \qquad 21.73$$

$$w=\frac{1}{r}=\frac{1}{\varepsilon D}[1+\varepsilon\cos(\phi Q)] \qquad \frac{dw}{d\phi}=\frac{-Qsen(\phi Q)}{D} \qquad \frac{d^2w}{d\phi^2}=\frac{-Q^2\cos(\phi Q)}{D} \qquad 21.38$$

$$H\frac{-Q^2\cos(\phi Q)}{D}+H\frac{1}{\varepsilon D}[1+\varepsilon\cos(\phi Q)]+3A\frac{-Q^2\cos(\phi Q)}{D}\frac{1}{\varepsilon D}[1+\varepsilon\cos(\phi Q)]+3A\left\{\frac{1}{\varepsilon D}[1+\varepsilon\cos(\phi Q)]\right\}^2+B=zero \qquad 21.74$$

$$-Q^2H\frac{\cos(\phi Q)}{D}+H\frac{1}{\varepsilon D}+H\frac{1}{\varepsilon D}\varepsilon\cos(\phi Q)-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}[1+\varepsilon\cos(\phi Q)]+\frac{3A}{\varepsilon^2D^2}\left[1+2\varepsilon\cos(\phi Q)+\varepsilon^2\cos^2(\phi Q)\right]+B=zero$$

$$-Q^2H\frac{\cos(\phi Q)}{D}+H\frac{1}{\varepsilon D}+H\frac{\cos(\phi Q)}{D}-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}\varepsilon\cos(\phi Q)+$$
$$+\frac{3A}{\varepsilon^2D^2}+\frac{3A}{\varepsilon^2D^2}2\varepsilon\cos(\phi Q)+\frac{3A}{\varepsilon^2D^2}\varepsilon^2\cos^2(\phi Q)+B=zero$$

$$-Q^2H\frac{\cos(\phi Q)}{D}+H\frac{1}{\varepsilon D}+H\frac{\cos(\phi Q)}{D}-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}-3Q^2A\frac{\cos^2(\phi Q)}{D^2}+$$
$$+\frac{3A}{\varepsilon^2D^2}+\frac{6A}{\varepsilon D}\frac{\cos(\phi Q)}{D}+3A\frac{\cos^2(\phi Q)}{D^2}+B=zero$$

$$-Q^2H\frac{\cos(\phi Q)}{D}+H\frac{\cos(\phi Q)}{D}-\frac{3Q^2A}{\varepsilon D}\frac{\cos(\phi Q)}{D}+\frac{6A}{\varepsilon D}\frac{\cos(\phi Q)}{D}-$$
$$-3Q^2A\frac{\cos^2(\phi Q)}{D^2}+3A\frac{\cos^2(\phi Q)}{D^2}+H\frac{1}{\varepsilon D}+\frac{3A}{\varepsilon^2D^2}+B=zero$$

$$\left(-Q^2H+H-\frac{3Q^2A}{\varepsilon D}+\frac{6A}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\left(-3Q^2A+3A\right)\frac{\cos^2(\phi Q)}{D^2}+H\frac{1}{\varepsilon D}+\frac{3A}{\varepsilon^2D^2}+B=zero$$

$$\left(-3Q^2A+3A\right)\frac{\cos^2(\phi Q)}{3AD^2}+\left(-Q^2H+H-\frac{3Q^2A}{\varepsilon D}+\frac{6A}{\varepsilon D}\right)\frac{\cos(\phi Q)}{3AD}+H\frac{1}{3A\varepsilon D}+\frac{3A}{3A\varepsilon^2D^2}+\frac{B}{3A}=zero$$

$$\left(1-Q^2\right)\frac{\cos^2(\phi Q)}{D^2}+\left(\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}+\frac{B}{3A}=zero \qquad 21.75$$

$$Q^2\approx 1 \qquad \left(1-Q^2\right)\frac{\cos^2(\phi Q)}{D^2}=zero \qquad 21.76$$

$$\left(\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}+\frac{B}{3A}=zero \qquad 21.77$$

$$\frac{\cos(\phi Q)}{D}=zero\Rightarrow\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}+\frac{B}{3A}=zero$$

$$\frac{\cos(\phi Q)}{D}\neq zero\Rightarrow\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}=zero$$

$$\frac{-Q^2H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}=zero \qquad \frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2D^2}+\frac{B}{3A}=zero \qquad 21.78$$

$$[a=b]\Rightarrow\frac{H}{3A}+\frac{1}{\varepsilon D}=\frac{1}{Q^2}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right) \qquad [a=c]\Rightarrow\frac{H}{3A}+\frac{1}{\varepsilon D}=-\frac{\varepsilon DB}{3A} \qquad 21.79$$

$$Q^2=1 \qquad H=\frac{E_R}{m_oc^2}=\frac{-m_oc^2}{m_oc^2}=-1 \qquad \varepsilon DB=\frac{\varepsilon DGM_o}{L'^2}=\frac{\varepsilon DGM_o}{\varepsilon DGM_o}=1$$

$$[a=b]\Rightarrow\frac{H}{3A}+\frac{1}{\varepsilon D}=\frac{1}{1}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right)\Rightarrow\frac{1}{\varepsilon D}=zero \qquad [a=c]\Rightarrow\frac{-1}{3A}+\frac{1}{\varepsilon D}=-\frac{1}{3A}\Rightarrow\frac{1}{\varepsilon D}=zero$$

$$[b=c]\Rightarrow\frac{1}{Q^2}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right)=-\frac{\varepsilon DB}{3A} \qquad 21.80$$

$$\varepsilon DB=\frac{\varepsilon DGM_o}{L'^2}=\frac{\varepsilon DGM_o}{\varepsilon DGM_o}=1 \qquad 21.81$$

$$[b=c] \Rightarrow \frac{1}{Q^2}\left(\frac{H}{3A}+\frac{2}{\varepsilon D}\right)=-\frac{1}{3A} \qquad Q^2=-H-\frac{6A}{\varepsilon D} \tag{21.82}$$

$Q=Q(H)$ The advance is a function of negative energy that governs the movement

$$H=\frac{E_R}{m_o c^2}=\frac{-m_o c^2}{m_o c^2}=-1 \qquad Q^2=-(-1)-\frac{6A}{\varepsilon D} \Rightarrow Q^2=1-\frac{6A}{\varepsilon D} \text{ Advance} \tag{21.83}$$

$$[a=b] \Rightarrow \frac{-1}{3A}+\frac{1}{\varepsilon D}=\frac{1}{\left(1-\frac{6A}{\varepsilon D}\right)}\left(\frac{-1}{3A}+\frac{2}{\varepsilon D}\right) \Rightarrow \frac{1}{\varepsilon D}=zero \tag{21.84}$$

$$H=\frac{E_R}{m_o c^2} \qquad A=\frac{GM_o}{c^2} \qquad B=\frac{GM_o}{L'^2}$$

$$\frac{-Q^2 H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}=zero \qquad \frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2 D^2}+\frac{B}{3A}=zero \tag{21.78}$$

$$3A\varepsilon D\left(\frac{-Q^2 H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)=zero \qquad 3A\varepsilon^2 D^2\left(\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2 D^2}+\frac{B}{3A}\right)=zero$$

$$-Q^2 H\varepsilon D+H\varepsilon D-Q^2 3A+6A=zero \qquad H\varepsilon D+3A+\varepsilon D(\varepsilon DB)=zero \tag{21.85}$$

$$\varepsilon DB=\frac{\varepsilon DGM_o}{L'^2}=\frac{\varepsilon DGM_o}{\varepsilon DGM_o}=1 \qquad H\varepsilon D=-3A-\varepsilon D \tag{21.86}$$

$$-Q^2(-3A-\varepsilon D)-3A-\varepsilon D-Q^2 3A+6A=zero \tag{21.87}$$

$$Q^2 3A+Q^2\varepsilon D-\varepsilon D-Q^2 3A+3A=zero$$

$$Q^2\varepsilon D-\varepsilon D+3A=zero \qquad Q^2=1-\frac{3A}{\varepsilon D} \tag{21.88}$$

This advance is not governed by negative energy

$$-Q^2 H\varepsilon D+H\varepsilon D-Q^2 3A+6A=zero \tag{21.85}$$

$$-Q^2(-3A-\varepsilon D)+H\varepsilon D-Q^2 3A+6A=zero \tag{21.89}$$

$$Q^2 3A+Q^2\varepsilon D+H\varepsilon D-Q^2 3A+6A=zero$$

$$Q^2\varepsilon D+H\varepsilon D+6A=zero \qquad Q^2=-H-\frac{6A}{\varepsilon D} \tag{21.90}$$

$$\left(\frac{-Q^2 H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2 D^2}+\frac{B}{3A}=zero \tag{21.77}$$

$$3A\varepsilon^2 D^2\left[\left(\frac{-Q^2 H}{3A}+\frac{H}{3A}-\frac{Q^2}{\varepsilon D}+\frac{2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H}{3A\varepsilon D}+\frac{1}{\varepsilon^2 D^2}+\frac{B}{3A}\right]=zero$$

$$\varepsilon D\left(\frac{-Q^2 H3A\varepsilon D}{3A}+\frac{H3A\varepsilon D}{3A}-\frac{Q^2 3A\varepsilon D}{\varepsilon D}+\frac{2\,.\,3A\varepsilon D}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{H3A\varepsilon^2 D^2}{3A\varepsilon D}+\frac{3A\varepsilon^2 D^2}{\varepsilon^2 D^2}+\frac{B3A\varepsilon^2 D^2}{3A}=zero$$

$$\varepsilon D\left(-Q^2H\varepsilon D+H\varepsilon D-Q^2 3A+6A\right)\frac{\cos(\phi Q)}{D}+H\varepsilon D+3A+\varepsilon D(\varepsilon DB)=zero$$

$$\varepsilon DB=\frac{\varepsilon DGM_o}{L'^2}=\frac{\varepsilon DGM_o}{\varepsilon DGM_o}=1 \qquad H=\frac{E_R}{m_o c^2}=\frac{-m_o c^2}{m_o c^2}=-1$$

$$\varepsilon D\left(-Q^2H\varepsilon D+H\varepsilon D-Q^2 3A+6A\right)\frac{\cos(\phi Q)}{D}-\varepsilon D+3A+\varepsilon D=zero$$

$$\left(-Q^2H\varepsilon D+H\varepsilon D-Q^2 3A+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero \qquad 21.91$$

$$Q^2=1-\frac{3A}{\varepsilon D}$$

$$\left[-\left(1-\frac{3A}{\varepsilon D}\right)H\varepsilon D+H\varepsilon D-\left(1-\frac{3A}{\varepsilon D}\right)3A+6A\right]\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\left(-H\varepsilon D+H\varepsilon D\frac{3A}{\varepsilon D}+H\varepsilon D-3A+3A\frac{3A}{\varepsilon D}+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\left(-H\varepsilon D+H3A+H\varepsilon D-3A+\frac{9A^2}{\varepsilon D}+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\left(H3A+\frac{9A^2}{\varepsilon D}+3A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$H=\frac{E_R}{m_o c^2}=\frac{-m_o c^2}{m_o c^2}=-1$$

$$\left(-3A+\frac{9A^2}{\varepsilon D}+3A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\frac{9A^2}{\varepsilon D}\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero \qquad \frac{\cos(\phi Q)}{D}+\frac{1}{3A}=zero \qquad 21.92$$

$$\left(-Q^2H\varepsilon D+H\varepsilon D-Q^2 3A+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero \qquad 21.91$$

$$Q^2=1-\frac{6A}{\varepsilon D}$$

$$\left[-\left(1-\frac{6A}{\varepsilon D}\right)H\varepsilon D+H\varepsilon D-\left(1-\frac{6A}{\varepsilon D}\right)3A+6A\right]\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\left(-H\varepsilon D+H\varepsilon D\frac{6A}{\varepsilon D}+H\varepsilon D-3A+3A\frac{6A}{\varepsilon D}+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\left(-H\varepsilon D+H6A+H\varepsilon D-3A+\frac{18A^2}{\varepsilon D}+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\left(H6A+\frac{18A^2}{\varepsilon D}+3A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$H=\frac{E_R}{m_o c^2}=\frac{-m_o c^2}{m_o c^2}=-1$$

$$\left(-6A+\frac{18A^2}{\varepsilon D}+3A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$\frac{1}{3A}\left[\left(-3A+\frac{18A^2}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}\right]=zero$$

$$\left(-1+\frac{6A}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{1}{\varepsilon D}=zero$$

$$-\left(1-\frac{6A}{\varepsilon D}\right)\frac{\cos(\phi Q)}{D}+\frac{1}{\varepsilon D}=zero \qquad -Q^2\frac{\cos(\phi Q)}{D}+\frac{1}{\varepsilon D}=zero \qquad 21.93$$

$$\left(-Q^2H\varepsilon D+H\varepsilon D-Q^2 3A+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero \qquad 21.91$$

$$Q^2=1 \qquad H=\frac{E_R}{m_o c^2}=\frac{-m_o c^2}{m_o c^2}=-1$$

$$\left(\varepsilon D-\varepsilon D-3A+6A\right)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero$$

$$(3A)\frac{\cos(\phi Q)}{D}+\frac{3A}{\varepsilon D}=zero \qquad \frac{\cos(\phi Q)}{D}+\frac{1}{\varepsilon D}=zero \qquad 21.94$$

$$Q^2=1-\frac{6A}{\varepsilon D} \qquad Q^2=1 \qquad Q^2=1-\frac{3A}{\varepsilon D}$$

$$\left|-Q^2\frac{\cos(\phi Q)}{D}+\frac{1}{\varepsilon D}\right| << \left|\frac{\cos(\phi Q)}{D}+\frac{1}{\varepsilon D}\right| <<<<<< \left|\frac{\cos(\phi Q)}{D}+\frac{1}{3A}\right| \qquad 21.95$$

Energy Newtonian (E_N)

$$E_N = \frac{m_o u^2}{2} - \frac{k}{r}$$

$$u^2 = \left(\frac{dr}{dt}\right)^2 + \left(r\frac{d\phi}{dt}\right)^2 = \left(\frac{dr}{dt}\right)^2 + \frac{L^2}{r^2}$$

$$E_N = \frac{m_o}{2}\left[\left(\frac{dr}{dt}\right)^2 + \frac{L^2}{r^2}\right] - \frac{k}{r}$$

$$\frac{2E_N}{m_o} = \left(\frac{dr}{dt}\right)^2 + \frac{L^2}{r^2} - \frac{2k}{m_o}\frac{1}{r}$$

$$\left(\frac{dr}{dt}\right)^2 + \frac{L^2}{r^2} - \frac{2k}{m_o}\frac{1}{r} - \frac{2E_N}{m_o} = zero$$

$$\frac{d\phi}{dt} = \frac{L}{r^2} \qquad \frac{dr}{dt} = -L\frac{dw}{d\phi} \qquad \frac{d^2r}{dt^2} = \frac{-L^2}{r^2}\frac{d^2w}{d\phi^2} \qquad \frac{d^2\phi}{dt^2} = \frac{2L^2}{r^3}\frac{dw}{d\phi}$$

$$\left(-L\frac{dw}{d\phi}\right)^2 + \frac{L^2}{r^2} - \frac{2k}{m_o}\frac{1}{r} - \frac{2E_N}{m_o} = zero$$

$$\left(\frac{dw}{d\phi}\right)^2 + \frac{1}{r^2} - \frac{2k}{m_o L^2}\frac{1}{r} - \frac{2E_N}{m_o L^2} = zero$$

$$\left(\frac{dw}{d\phi}\right)^2 + \frac{1}{r^2} - \frac{2k}{m_o L^2}\frac{1}{r} - \frac{2E_N}{m_o L^2} = zero$$

$$\left(\frac{dw}{d\phi}\right)^2 + w^2 - \frac{2k}{m_o L^2}w - \frac{2E_N}{m_o L^2} = zero$$

$$x = \frac{2k}{m_o L^2} \qquad y = \frac{2E_N}{m_o L^2}$$

$$\left(\frac{dw}{d\phi}\right)^2 + w^2 - xw - y = zero$$

$$w = \frac{1}{r} = \frac{1}{\varepsilon D}\left[1 + \varepsilon\cos(\phi Q)\right] \qquad \frac{dw}{d\phi} = \frac{-Q\,sen(\phi Q)}{D} \qquad \frac{d^2w}{d\phi^2} = \frac{-Q^2\cos(\phi Q)}{D}$$

$$\left[\frac{-Q\,sen(\phi Q)}{D}\right]^2 + \left\{\frac{1}{\varepsilon D}\left[1+\varepsilon\cos(\phi Q)\right]\right\}^2 - x\frac{1}{\varepsilon D}\left[1+\varepsilon\cos(\phi Q)\right] - y = zero$$

$$\frac{Q^2}{D^2}\left[1-\cos^2(\phi Q)\right] + \frac{1}{\varepsilon^2 D^2}\left[1 + 2\varepsilon\cos(\phi Q) + \varepsilon^2\cos^2(\phi Q)\right] - x\frac{1}{\varepsilon D} - x\frac{1}{\varepsilon D}\varepsilon\cos(\phi Q) - y = zero$$

$$\frac{Q^2}{D^2}-\frac{Q^2}{D^2}\cos^2(\phi Q)+\frac{1}{\varepsilon^2 D^2}+\frac{1}{\varepsilon^2 D^2}2\varepsilon\cos(\phi Q)+\frac{1}{\varepsilon^2 D^2}\varepsilon^2\cos^2(\phi Q)-\frac{x}{\varepsilon D}-x\frac{\cos(\phi Q)}{D}-y=zero$$

$$\frac{Q^2}{D^2}-Q^2\frac{\cos^2(\phi Q)}{D^2}+\frac{1}{\varepsilon^2 D^2}+\frac{2}{\varepsilon D}\frac{\cos(\phi Q)}{D}+\frac{\cos^2(\phi Q)}{D^2}-\frac{x}{\varepsilon D}-x\frac{\cos(\phi Q)}{D}-y=zero$$

$$\frac{\cos^2(\phi Q)}{D^2}-Q^2\frac{\cos^2(\phi Q)}{D^2}+\frac{2}{\varepsilon D}\frac{\cos(\phi Q)}{D}-x\frac{\cos(\phi Q)}{D}+\frac{Q^2}{D^2}+\frac{1}{\varepsilon^2 D^2}-\frac{x}{\varepsilon D}-y=zero$$

$$(1-Q^2)\frac{\cos^2(\phi Q)}{D^2}+\left(\frac{2}{\varepsilon D}-x\right)\frac{\cos(\phi Q)}{D}+\frac{Q^2}{D^2}+\frac{1}{\varepsilon^2 D^2}-\frac{x}{\varepsilon D}-y=zero$$

$$Q^2\approx 1 \qquad\qquad (1-Q^2)\frac{\cos^2(\phi Q)}{D^2}=zero$$

$$\left(\frac{2}{\varepsilon D}-x\right)\frac{\cos(\phi Q)}{D}+\frac{1}{D^2}+\frac{1}{\varepsilon^2 D^2}-\frac{x}{\varepsilon D}-y=zero$$

$$\left(\frac{2}{\varepsilon D}-x\right)=zero \qquad\qquad \frac{1}{D^2}+\frac{1}{\varepsilon^2 D^2}-\frac{x}{\varepsilon D}-y=zero$$

$$x=\frac{2k}{m_o L^2} \qquad\qquad y=\frac{2E_N}{m_o L^2}$$

$$\frac{2}{\varepsilon D}-x=zero\Rightarrow x=\frac{2}{\varepsilon D}=\frac{2k}{m_o L^2}\Rightarrow\frac{1}{\varepsilon D}=\frac{GM_o m_o}{m_o L^2}\Rightarrow L^2=\varepsilon DGM_o$$

$$\frac{\varepsilon^2 D^2}{D^2}+\frac{\varepsilon^2 D^2}{\varepsilon^2 D^2}-\frac{\varepsilon^2 D^2 x}{\varepsilon D}-\varepsilon^2 D^2 y=zero$$

$$\varepsilon^2+1-\varepsilon Dx-\varepsilon^2 D^2 y=zero$$

$$\varepsilon Dx=\varepsilon D\frac{2}{\varepsilon D}\Rightarrow\varepsilon Dx=2 \qquad\qquad \varepsilon^2 D^2 y=\varepsilon^2 D^2\frac{2E_N}{m_o L^2}=\varepsilon^2 D^2\frac{2E_N}{m_o\varepsilon DGM_o}=\frac{2\varepsilon DE_N}{k}$$

$$\varepsilon^2+1-2-\frac{2\varepsilon DE_N}{k}=zero \qquad\qquad E_N=\frac{k}{2\varepsilon D}(\varepsilon^2-1)$$

$$\frac{1}{a}=\frac{1}{\varepsilon D}(1-\varepsilon^2) \qquad\qquad E_N=\frac{-k}{2a}$$

§22 Spatial deformation

$$t=\frac{t'}{\sqrt{1-\frac{v^2}{c^2}}} \qquad t>t'$$

$$t=t_1+t_2=\frac{L}{c-v}+\frac{L}{c+v}=\frac{2L}{c}\frac{1}{\left(1-\frac{v^2}{c^2}\right)} \qquad t'=\frac{2L'}{c}$$

$$t=\frac{2L}{c}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}=\frac{\frac{2L'}{c}}{\sqrt{1-\frac{v^2}{c^2}}}\Rightarrow L=L'\sqrt{1-\frac{v^2}{c^2}} \qquad L'>L$$

This is the spatial deformation.

The length L' at rest in the reference frame of the observer O' is greater than the length L that is moving with velocity relative v on reference frame the observer O.

Now compute to the observer O' the distance $d'=vt'$ between $O\leftrightarrow O'$:

$$d'=vt'=v\frac{2L'}{c}$$

Thus we obtain the velocity v: $d'=v\frac{2L'}{c}\Rightarrow v=\frac{cd'}{2L'}$.

Now compute to the observer O the distance $d=vt$ between $O\leftrightarrow O'$:

$$d=vt=v(t_1+t_2)=v\frac{2L}{c}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}$$

Thus we obtain the velocity v: $d=v\frac{2L}{c}\frac{1}{\left(1-\frac{v^2}{c^2}\right)}\Rightarrow v=\frac{cd}{2L}\left(1-\frac{v^2}{c^2}\right)$.

The speed v is the same to both observers so we have:

$$v=\frac{cd'}{2L'}=\frac{cd}{2L}\left(1-\frac{v^2}{c^2}\right)$$

Where applying the relation $L=L'\sqrt{1-\frac{v^2}{c^2}}$ we obtain:

$$\frac{cd'}{2L'}=\frac{cd}{2L'\sqrt{1-\frac{v^2}{c^2}}}\left(1-\frac{v^2}{c^2}\right)\Rightarrow d'=d\sqrt{1-\frac{v^2}{c^2}} \qquad d>d'.$$

Where the distance d and d’ varies inversely with the distances L and L’.

In general, we obtain (14.2, 14.4):

$$d' = \frac{d\left(1 - \frac{vux}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad \text{or} \qquad d = \frac{d'\left(1 + \frac{vu'x'}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$u'x' = zero \qquad d = \frac{d'\left[1 + \frac{v(0)}{c^2}\right]}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad d = \frac{d'}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$u'x' = c \qquad d = \frac{d'\left(1 + \frac{vc}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad d = d'\sqrt{\frac{1 + \frac{v}{c}}{1 - \frac{v}{c}}}$$

$$u'x' = -v \qquad d = \frac{d'\left[1 + \frac{v(-v)}{c^2}\right]}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad d = d'\sqrt{1 - \frac{v^2}{c^2}}$$

$$ux = v \qquad d' = \frac{d\left[1 - \frac{v(v)}{c^2}\right]}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad d' = d\sqrt{1 - \frac{v^2}{c^2}}$$

$$ux = c \qquad d' = \frac{d\left(1 - \frac{vc}{c^2}\right)}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad d' = d\sqrt{\frac{1 - \frac{v}{c}}{1 + \frac{v}{c}}}$$

$$ux = zero \qquad d' = \frac{d\left[1 - \frac{v(0)}{c^2}\right]}{\sqrt{1 - \frac{v^2}{c^2}}} \qquad d' = \frac{d}{\sqrt{1 - \frac{v^2}{c^2}}}$$

§23 Space and Time Bend

Variables with line $t',v',x',y',\vec{r}'$ etc ...They are used in §21.

Geometry of space and time in the plan $xy \to y \perp x$.

$$y=f(x)$$

$$x=ct' \qquad y=\int ds'=\int\sqrt{d\vec{r}'.d\vec{r}'}$$

$$\int ds'=f(ct')$$

$$dx=cdt' \qquad dy=ds'=\sqrt{d\vec{r}'.d\vec{r}'}$$

$$\vec{r}=x\hat{\imath}+y\hat{\jmath}=ct'\hat{\imath}+\int ds'\hat{\jmath} \qquad \vec{r}'=x'\hat{\imath}+y'\hat{\jmath}$$

$$d\vec{r}=dx\hat{\imath}+dy\hat{\jmath}=cdt'\hat{\imath}+ds'\hat{\jmath} \qquad d\vec{r}'=dx'\hat{\imath}+dy'\hat{\jmath}$$

$$dr=\frac{\vec{r}.d\vec{r}}{r}=\frac{x}{r}dx+\frac{y}{r}dy$$

$$\vec{v}=\frac{d\vec{r}}{dt'}=\frac{dx}{dt'}\hat{\imath}+\frac{dy}{dt'}\hat{\jmath}=\frac{cdt'}{dt'}\hat{\imath}+\frac{ds'}{dt'}\hat{\jmath}=c\hat{\imath}+v'\hat{\jmath} \qquad \vec{v}'=\frac{d\vec{r}'}{dt'}=\frac{dx'}{dt'}\hat{\imath}+\frac{dy'}{dt'}\hat{\jmath}$$

$$\frac{dx}{dt'}=c \qquad \frac{dy}{dt'}=\frac{ds'}{dt'}=v' \qquad c=v\cos\varphi \qquad v'=v\,sen\varphi$$

$$tg\varphi=\frac{dy}{dx}=\frac{\frac{dy}{dt'}}{\frac{dx}{dt'}}=\frac{\frac{ds'}{dt'}}{c}=\frac{1}{c}\frac{ds'}{dt'} \qquad \frac{d^2y}{dx^2}=\frac{d}{dx}\left(\frac{dy}{dx}\right)=\frac{1}{c}\frac{d}{dt'}\left(\frac{1}{c}\frac{ds'}{dt'}\right)=\frac{1}{c^2}\frac{d^2s'}{dt'^2}$$

$$\vec{v}=\vec{c}+\vec{v}' \qquad \vec{c}=c\hat{\imath} \qquad \vec{v}'=v'\hat{\jmath}$$

$$\vec{a}=\frac{d\vec{v}}{dt'}=\frac{d\vec{c}}{dt'}+\frac{d\vec{v}'}{dt'} \qquad \frac{d\vec{c}}{dt'}=zero \qquad \frac{d\vec{v}}{dt'}=\frac{d\vec{v}'}{dt'}\to\vec{a}=\vec{a}'$$

$$ds^2=d\vec{r}.d\vec{r}=(dx\hat{\imath}+dy\hat{\jmath})(dx\hat{\imath}+dy\hat{\jmath})=(cdt'\hat{\imath}+ds'\hat{\jmath})(cdt'\hat{\imath}+ds'\hat{\jmath})=dx^2+dy^2=c^2dt'^2+ds'^2$$

$$ds=\sqrt{c^2dt'^2+ds'^2} \qquad ds'=\sqrt{ds^2-c^2dt'^2}$$

$$v=\frac{ds}{dt'}=\sqrt{c^2+\left(\frac{ds'}{dt'}\right)^2}=\sqrt{c^2+v'^2}>c \qquad v'=\frac{ds'}{dt'}=\sqrt{\left(\frac{ds}{dt'}\right)^2-c^2}=\sqrt{v^2-c^2}$$

$$\mathrm{K}=\left|\frac{d\varphi}{ds}\right|\to\varphi\equiv\angle$$ theoretical curve

$$tg\varphi=\frac{dy}{dx} \qquad \varphi=arctg\frac{dy}{dx} \qquad \frac{d\varphi}{dx}=\frac{\frac{d^2y}{dx^2}}{1+\left(\frac{dy}{dx}\right)^2}=\frac{\frac{1}{c^2}\frac{d^2s'}{dt'^2}}{1+\frac{1}{c^2}\left(\frac{ds'}{dt'}\right)^2}$$

$$\frac{ds}{dx}=\sqrt{1+\left(\frac{dy}{dx}\right)^2}=\sqrt{1+\frac{1}{c^2}\left(\frac{ds'}{dt'}\right)^2}$$

$$\mathrm{K}=\frac{d\varphi}{ds}=\frac{\frac{d\varphi}{dx}}{\frac{ds}{dx}}=\frac{\frac{\frac{1}{c^2}\frac{d^2s'}{dt'^2}}{1+\frac{1}{c^2}\left(\frac{ds'}{dt'}\right)^2}}{\sqrt{1+\frac{1}{c^2}\left(\frac{ds'}{dt'}\right)^2}}=\frac{\frac{1}{c^2}\frac{d^2s'}{dt'^2}}{\left[1+\frac{1}{c^2}\left(\frac{ds'}{dt'}\right)^2\right]^{\frac{3}{2}}}$$

$$\frac{ds'}{dt'}\mathrm{K}=\frac{ds'}{dt'}\frac{d\varphi}{ds}=v'\mathrm{K}=v'\frac{d\varphi}{ds}=\frac{\frac{1}{c^2}\frac{ds'}{dt'}\frac{d^2s'}{dt'^2}}{\left[1+\frac{1}{c^2}\left(\frac{ds'}{dt'}\right)^2\right]^{\frac{3}{2}}}=\frac{\frac{1}{c^2}\vec{v}'\frac{d\vec{v}'}{dt'}}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}$$

$$\vec{v}'\vec{\mathrm{K}}=\vec{v}'\frac{d\vec{\varphi}}{ds}=\frac{\frac{1}{c^2}\vec{v}'\frac{d\vec{v}'}{dt'}}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}} \qquad \vec{\mathrm{K}}=\frac{d\vec{\varphi}}{ds}=\frac{\frac{1}{c^2}\frac{d\vec{v}'}{dt'}}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}$$

$$dE_k=\frac{m_o v dv}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{m_o v' dv'}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=-\frac{k}{r^2}dr=\frac{k}{r^2}\hat{r}d\vec{r}' \qquad 21.56$$

$$\frac{dE_k}{dt'}=\vec{F}'\,.\vec{v}'=\frac{m_o\frac{c^2}{c^2}\vec{v}'\frac{d\vec{v}'}{dt'}}{\left(1+\frac{v'^2}{c^2}\right)^{\frac{3}{2}}}=\frac{k}{r^2}\hat{r}\frac{d\vec{r}'}{dt'}=\frac{k}{r^2}\hat{r}\vec{v}'$$

$$\frac{dE_k}{dt'}=\vec{F}'\,.\vec{v}'=m_oc^2\vec{v}'\frac{d\vec{\varphi}}{ds}=\frac{k}{r^2}\hat{r}\vec{v}'$$

$$\vec{F}'=m_oc^2\frac{d\vec{\varphi}}{ds}=\frac{k}{r^2}\hat{r} \qquad \vec{\mathrm{K}}=\frac{d\vec{\varphi}}{ds}=\frac{k}{m_oc^2}\frac{1}{r^2}\hat{r}$$

§ 23 electronic translation

§24 Variational Principle

$$E_k=\frac{m_o c^2}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{k}{r}+\text{constant} \qquad 21.21$$

$$E_k=\frac{m_o v^2}{\sqrt{1-\frac{v^2}{c^2}}}+m_o c^2\sqrt{1-\frac{v^2}{c^2}}=\frac{m_o c^2}{\sqrt{1-\frac{v^2}{c^2}}}=\frac{k}{r}+\text{constant}$$

$$\frac{m_o v^2}{\sqrt{1-\frac{v^2}{c^2}}}-\left(-m_o c^2\sqrt{1-\frac{v^2}{c^2}}+\frac{k}{r}\right)=m_o c^2 \qquad p=\frac{d}{dv}\left(-m_o c^2\sqrt{1-\frac{v^2}{c^2}}\right)=\frac{m_o v}{\sqrt{1-\frac{v^2}{c^2}}}$$

$L=-m_o c^2\sqrt{1-\frac{v^2}{c^2}}+\frac{k}{r}$ Lagrangeana.

$\frac{m_o v^2}{\sqrt{1-\frac{v^2}{c^2}}}-L=m_o c^2$ What is the initial energy of the particle of mass mo.

$$pv-L=m_o c^2 \qquad L=pv-m_o c^2=-m_o c^2\sqrt{1-\frac{v^2}{c^2}}+\frac{k}{r}$$

Variational Principle

$\text{Ação}=S=\int_{t_1}^{t_2} L[(x(t),\dot{x}(t),t)]dt$ $\qquad \dot{x}=\frac{dx}{dt}=ux$ This is the velocity component in x axis.

$\delta S=\delta\int_{t_1}^{t_2} L(x,\dot{x},t)dt=\text{zero}$ Variation of the action along the X axis.

Building the variable $x'=x+\varepsilon\eta$ in the range $t_1\le t\le t_2$ we have seen this when $\varepsilon\to\text{zero}\Rightarrow x'=x$ and where $\varepsilon\ne\text{zero}$ we will have the conditions:

$$\frac{d\varepsilon}{dt}=\text{zero} \qquad \eta=\eta(t) \qquad \eta(t_1)=\text{zero} \qquad \eta(t_2)=\text{zero} \qquad \frac{d\eta}{d\varepsilon}=\text{zero} \qquad \dot{\eta}=\frac{d\eta}{dt}$$

$$x'=x+\varepsilon\eta \qquad \dot{x}'=\dot{x}+\varepsilon\dot{\eta} \qquad \frac{dx'}{d\varepsilon}=\eta \qquad \frac{d\dot{x}'}{d\varepsilon}=\dot{\eta} \qquad \frac{dx}{d\varepsilon}=\text{zero} \qquad \frac{d\dot{x}}{d\varepsilon}=\text{zero}$$

Then we have a new function $I(\varepsilon)=\int_{t_1}^{t_2} G(x+\varepsilon\eta,\dot{x}+\varepsilon\dot{\eta},t)dt=\int_{t_1}^{t_2} F(x',\dot{x}',t)dt$ and where:

$$\varepsilon=\text{zero}\to x'=x\to\dot{x}'=\dot{x}\to F=L\Rightarrow\int_{t_1}^{t_2} F(x',\dot{x}',t)dt=\int_{t_1}^{t_2} L(x,\dot{x},t)dt$$

$$\varepsilon\ne\text{zero}\to x'\ne x\to\dot{x}'\ne\dot{x}\to F\ne L\Rightarrow\int_{t_1}^{t_2} F(x',\dot{x}',t)dt\ne\int_{t_1}^{t_2} L(x,\dot{x},t)dt$$

So we have $I(\varepsilon)=\int_{t_1}^{t_2} F[x'(\varepsilon),\dot{x}'(\varepsilon),t]dt$ that provides derived:

$$\frac{\delta I(\varepsilon)}{d\varepsilon}=\int_{t_1}^{t_2}\frac{\partial F(x',\dot{x}',t)}{\partial x'}\frac{dx'}{d\varepsilon}dt+\int_{t_1}^{t_2}\frac{\partial F(x',\dot{x}',t)}{\partial \dot{x}'}\frac{d\dot{x}'}{d\varepsilon}dt=\int_{t_1}^{t_2}\frac{\partial F}{\partial x'}\eta dt+\int_{t_1}^{t_2}\frac{\partial F}{\partial \dot{x}'}\dot{\eta} dt=zero$$

$$\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\eta\right)=\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\eta+\frac{\partial F}{\partial \dot{x}'}\frac{d\eta}{dt}\Rightarrow\frac{\partial F}{\partial \dot{x}'}\dot{\eta}=\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\eta\right)-\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\eta$$

$$\frac{\delta I(\varepsilon)}{d\varepsilon}=\int_{t_1}^{t_2}\frac{\partial F}{\partial x'}\eta dt+\int_{t_1}^{t_2}\frac{\partial F}{\partial \dot{x}'}\dot{\eta} dt=\int_{t_1}^{t_2}\frac{\partial F}{\partial x'}\eta dt+\int_{t_1}^{t_2}\left[\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\eta\right)-\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\eta\right]dt=zero$$

$$\frac{\delta I(\varepsilon)}{d\varepsilon}=\int_{t_1}^{t_2}\frac{\partial F}{\partial x'}\eta dt+\int_{t_1}^{t_2}d\left(\frac{\partial F}{\partial \dot{x}'}\eta\right)-\int_{t_1}^{t_2}\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\eta dt=zero$$

$$\int_{t_1}^{t_2}d\left(\frac{\partial F}{\partial \dot{x}'}\eta\right)=\frac{\partial F}{\partial \dot{x}'}\eta\Big|_{t_1}^{t_2}=\frac{\partial F}{\partial \dot{x}'}\eta(t_2)-\frac{\partial F}{\partial \dot{x}'}\eta(t_1)=zero$$

$$\frac{\delta I(\varepsilon)}{d\varepsilon}=\int_{t_1}^{t_2}\frac{\partial F}{\partial x'}\eta dt-\int_{t_1}^{t_2}\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\eta dt=\int_{t_1}^{t_2}\left[\frac{\partial F}{\partial x'}-\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\right]\eta dt=zero$$

$$\frac{\delta I(\varepsilon)}{d\varepsilon}=\int_{t_1}^{t_2}\left[\frac{\partial F}{\partial x'}-\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)\right]\eta dt=zero\Rightarrow\eta\neq zero\rightarrow\frac{\partial F}{\partial x'}-\frac{d}{dt}\left(\frac{\partial F}{\partial \dot{x}'}\right)=zero$$

$$\varepsilon=zero\rightarrow x'=x\rightarrow\dot{x}'=\dot{x}\rightarrow F=L\Rightarrow\frac{\partial L}{\partial x}-\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{x}}\right)=zero$$

$\frac{\partial L}{\partial x}=\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{x}}\right)$ This is the X axis component $\qquad L=-m_oc^2\sqrt{1-\frac{v^2}{c^2}}+\frac{k}{r}$

$$\frac{\partial}{\partial x}\left(-m_oc^2\sqrt{1-\frac{v^2}{c^2}}+\frac{k}{r}\right)=\frac{d}{dt}\left[\frac{\partial}{\partial \dot{x}}\left(-m_oc^2\sqrt{1-\frac{v^2}{c^2}}+\frac{k}{r}\right)\right]$$

$$\frac{\partial}{\partial x}\left(-m_oc^2\sqrt{1-\frac{v^2}{c^2}}\right)=zero \qquad \frac{\partial}{\partial \dot{x}}\left(\frac{k}{r}\right)=zero \qquad v=\sqrt{\frac{dx}{dt}^2+\frac{dy}{dt}^2+\frac{dz}{dt}^2}=\sqrt{\dot{x}^2+\dot{y}^2+\dot{z}^2}$$

$\frac{\partial}{\partial x}\left(\frac{k}{r}\right)=\frac{d}{dt}\left[\frac{\partial}{\partial \dot{x}}\left(-m_oc^2\sqrt{1-\frac{v^2}{c^2}}\right)\right]$ This is the X axis component

$$\frac{\partial}{\partial x}\left(\frac{k}{r}\right)=k\frac{\partial}{\partial x}\left(r^{-1}\right)=k(-1)r^{-1-1}\frac{\partial r}{\partial x}=-k\frac{1}{r^2}\frac{x}{r}=-k\frac{x}{r^3} \qquad r^2=x^2+y^2+z^2$$

$$\frac{\partial}{\partial\dot{x}}\left(-m_o c^2\sqrt{1-\frac{v^2}{c^2}}\right)=-m_o c^2\frac{1}{2}\left(1-\frac{v^2}{c^2}\right)^{\frac{1}{2}-1}\left(-\frac{2v}{c^2}\frac{dv}{d\dot{x}}\right)=\frac{m_o v}{\sqrt{1-\frac{v^2}{c^2}}}\frac{d}{d\dot{x}}\left(\sqrt{\dot{x}^2+\dot{y}^2+\dot{z}^2}\right)$$

$$\frac{\partial}{\partial\dot{x}}\left(-m_o c^2\sqrt{1-\frac{v^2}{c^2}}\right)=\frac{m_o v}{\sqrt{1-\frac{v^2}{c^2}}}\left[\frac{1}{2}\left(\dot{x}^2+\dot{y}^2+\dot{z}^2\right)^{\frac{1}{2}-1}2\dot{x}\right]=\frac{m_o v}{\sqrt{1-\frac{v^2}{c^2}}}\left(\frac{\dot{x}}{\sqrt{\dot{x}^2+\dot{y}^2+\dot{z}^2}}\right)=\frac{m_o\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}$$

$$\frac{d}{dt}\left(\frac{m_o\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}\right)=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)}\left[\frac{d\dot{x}}{dt}\sqrt{1-\frac{v^2}{c^2}}-\dot{x}\frac{d}{dt}\left(\sqrt{1-\frac{v^2}{c^2}}\right)\right]=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)}\left[\frac{d\dot{x}}{dt}\sqrt{1-\frac{v^2}{c^2}}-\dot{x}\frac{1}{2}\left(1-\frac{v^2}{c^2}\right)^{\frac{1}{2}-1}\left(-\frac{2v}{c^2}\frac{dv}{dt}\right)\right]$$

$$\frac{d}{dt}\left(\frac{m_o\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}\right)=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)}\left[\frac{d\dot{x}}{dt}\sqrt{1-\frac{v^2}{c^2}}+\frac{\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}\left(\frac{v}{c^2}\frac{dv}{dt}\right)\right]=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)}\left[\frac{\frac{d\dot{x}}{dt}\sqrt{1-\frac{v^2}{c^2}}\sqrt{1-\frac{v^2}{c^2}}}{\sqrt{1-\frac{v^2}{c^2}}}+\frac{\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}\left(\frac{v}{c^2}\frac{dv}{dt}\right)\right]$$

$$\frac{d}{dt}\left(\frac{m_o\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}\right)=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)}\left[\frac{\frac{d\dot{x}}{dt}\sqrt{1-\frac{v^2}{c^2}}\sqrt{1-\frac{v^2}{c^2}}}{\sqrt{1-\frac{v^2}{c^2}}}+\frac{\dot{x}}{\sqrt{1-\frac{v^2}{c^2}}}\left(\frac{v}{c^2}\frac{dv}{dt}\right)\right]=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\dot{x}}{dt}+v\frac{dv}{dt}\frac{\dot{x}}{c^2}\right]$$

$$-k\frac{x}{r^3}\hat{i}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{x}+v\frac{dv}{dt}\frac{\dot{x}}{c^2}\right]\hat{i}\ \text{ X axis}$$

$$-k\frac{y}{r^3}\hat{j}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{y}+v\frac{dv}{dt}\frac{\dot{y}}{c^2}\right]\hat{j}\ \text{ Y axis}$$

$$-k\frac{z}{r^3}\hat{k}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{z}+v\frac{dv}{dt}\frac{\dot{z}}{c^2}\right]\hat{k}\ \text{ Z axis}$$

$$-k\frac{x}{r^3}\hat{i}-k\frac{y}{r^3}\hat{j}-k\frac{z}{r^3}\hat{k}=\frac{-k}{r^3}\left(x\hat{i}+y\hat{j}+z\hat{k}\right)=\frac{-k}{r^3}\vec{r}=\frac{-k}{r^2}\hat{r}$$

$$\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{x}+v\frac{dv}{dt}\frac{\dot{x}}{c^2}\right]\hat{i}+\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{y}+v\frac{dv}{dt}\frac{\dot{y}}{c^2}\right]\hat{j}+\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{z}+v\frac{dv}{dt}\frac{\dot{z}}{c^2}\right]\hat{k}=\frac{-k}{r^2}\hat{r}$$

$$\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left\{\left[\left(1-\frac{v^2}{c^2}\right)\ddot{x}+v\frac{dv}{dt}\frac{\dot{x}}{c^2}\right]\hat{i}+\left[\left(1-\frac{v^2}{c^2}\right)\ddot{y}+v\frac{dv}{dt}\frac{\dot{y}}{c^2}\right]\hat{j}+\left[\left(1-\frac{v^2}{c^2}\right)\ddot{z}+v\frac{dv}{dt}\frac{\dot{z}}{c^2}\right]\hat{k}\right\}=\frac{-k}{r^2}\hat{r}$$

$$\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\ddot{x}\hat{i}+v\frac{dv}{dt}\frac{\dot{x}}{c^2}\hat{i}+\left(1-\frac{v^2}{c^2}\right)\ddot{y}\hat{j}+v\frac{dv}{dt}\frac{\dot{y}}{c^2}\hat{j}+\left(1-\frac{v^2}{c^2}\right)\ddot{z}\hat{k}+v\frac{dv}{dt}\frac{\dot{x}}{c^2}\hat{k}\right]=\frac{-k}{r^2}\hat{r}$$

$$\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\left(\ddot{x}\hat{i}+\ddot{y}\hat{j}+\ddot{z}\hat{k}\right)+\frac{v}{c^2}\frac{dv}{dt}\left(\dot{x}\hat{i}++\dot{y}\hat{j}+\dot{z}\hat{k}\right)\right]=\frac{-k}{r^2}\hat{r}$$

$$\vec{a}=\ddot{x}\hat{i}+\ddot{y}\hat{j}+\ddot{z}\hat{k}=\frac{d}{dt}\left(\dot{x}\hat{i}+\dot{y}\hat{j}+\dot{z}\hat{k}\right)=\frac{d\vec{v}}{dt} \qquad \vec{v}=\dot{x}\hat{i}++\dot{y}\hat{j}+\dot{z}\hat{k}$$

$$\vec{F}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right]=\frac{-k}{r^2}\hat{r} \qquad = 21.16$$

$$\vec{F}=\frac{m_o}{\left(1-\frac{v^2}{c^2}\right)^{\frac{3}{2}}}\left[\left(1-\frac{v^2}{c^2}\right)\frac{d\vec{v}}{dt}+v\frac{dv}{dt}\frac{\vec{v}}{c^2}\right]=\frac{-k}{r^2}\hat{r} \qquad = 21.19$$

§ 24 electronic translation

"Although nobody can return behind and perform a new beginning,
any one can begin now and create a new end"
(Chico Xavier)

Author: Alfredo Dimas Moreira Garcia.

e-mail: avaliac@sjc.sp.gov.br

Translator: Rodolfo Marcos Venâncio

e-mail: sooter@terra.com.br

Bibliography

[1] Alonso & Finn, Vol. I e II, Ed. Edgard Blücher LTDA, 1972.
[2] Robert Resnik, Introduçao à Relatividade Especial, Ed. Univ. de S. Paulo e Ed. Polígono S.A., 1971.
[3] E. Terradas Y R. Ortiz, Relatividad, Cia. Ed. Espasa-Calpe Argentina S. A., 1952.
[4] Abraham Pais, “Sutil é o Senhor...” A ciência e a vida de Albert Einstein, Ed. Nova Fronteira, 1995.
[5] Marcel Rouault, Física Atômica, Ao Livro Técnico LTDA, 1959.
[6] H. A. Lorentz, A. Einstein e H. Minkowisk, O Princípio da Relatividade, edição da Fundação Calouste Gulbenkian, 1989.
[7] L. Landau e E. Lifchitz, Teoria do Campo, Hemus – Livraria Editora LTDA.
[8] Robert Martin Eisberg, Fundamentos da Física Moderna, Ed. Guanabara Dois S.A., 1979.
[9] Max Born, Física Atômica, edição da Fundação Calouste Gulbenkian, 1986.

http://www.wbabin.net/physics/faraj7.htm