Electro-Magnetic Field Equation and Lorentz gauge in Rindler spacetime

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ABSTRACT

In the general relativity theory, we find the electro-magnetic field transformation and the electro-magnetic field equation (Maxwell equation) in Rindler spacetime. We treat Lorentz gauge transformation, Lorentz gauge, Lorentz gauge fixing condition in Rindler spacetime. We prove the electro-magnetic wave function cannot exist in Rindler spacetime in Appendix A. Specially, this article say the uniqueness of the accelerated frame because the accelerated frame can treat electro-magnetic field equation.

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Key words:General relativity theory;

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Lorentz gauge fixing condition

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1. Introduction

In 2007 year, G.F.Torres del Castillio and C.I.Perez Sanchez already discovered Maxwell equations in uniformly accelerated frame in vacuum (see Ref [13]). In 2011 year, J.W.Maluf and F.F.Faria discovered electro-magnetic field transformation in Rindler space-time on ArXiv preprint(see also Ref[11]). But they did mistake they used Maxwell equations of gravity field. Maxwell equations of uniformly accelerated frame have to treat in flat Minkowski space-time not in gravity space-time.

Our theory's aim is that we find electro-magnetic field equation in Rindler space-time in vacuum also not in vacuum in the general relativity theory. In Section 2, we prepare for finding electro-magnetic field equation in Rindler space-time. In this section, we discover Lorentz gauge transformation and Lorentz gauge, Lorentz gauge fixing condition, transformation of the electro-magnetic 4-vector potential in Rindler space-time. In Section 3, we define the electro-magnetic field in Rindler space-time and we find the transformation of the electro-magnetic field. In Section 4, we obtain the electro-magnetic field equation in Rindler space-time and we apply the gauge theory to Maxwell equations (discovered by us) in Rindler space-time for viewing invariant about the gauge transformation.

We think seriously electro-magnetic wave function (radiation) in Rindler space-time but we know it doesn't satisfy electro-magnetic wave equation in mathematically (See APPENDIX A). Hence, in 2007 year and in 2011 year, all researchers mistake calculation of electro-magnetic wave function (see Ref [11],[13]) .

We understand electro-magnetic wave function can exist in inertial frame by J.C.Maxwell or A. Einstein.

2. Transformation of the electro-magnetic 4-vector potential, Lorentz gauge transformation and Lorentz gauge, Lorentz gauge fixing condition

The Rindler coordinate transformation is

$$ct = \left(\frac{c^2}{a_0} + \xi^1\right)\sinh\left(\frac{a_0\xi^0}{c}\right)$$

$$x = \left(\frac{c^2}{a_0} + \xi^1\right)\cosh\left(\frac{a_0\xi^0}{c}\right) - \frac{c^2}{a_0} \quad , y = \xi^2, z = \xi^3 \tag{1}$$

In this time, the tetrad $e^a{}_\mu$ is (see Ref [12], [14])

$$d\tau^2 = dt^2 - \frac{1}{c^2}[dx^2 + dy^2 + dz^2]$$

$$= -\frac{1}{c^2}\eta_{ab}\frac{\partial x^a}{\partial \xi^\mu}\frac{\partial x^b}{\partial \xi^\upsilon}d\xi^\mu d\xi^\upsilon$$

$$= -\frac{1}{c^2}\eta_{ab}e^a{}_\mu e^b{}_\upsilon d\xi^\mu d\xi^\upsilon = -\frac{1}{c^2} g_{\mu\upsilon} d\xi^\mu d\xi^\upsilon \ , \ e^a{}_\mu = \frac{\partial x^a}{\partial \xi^\mu} \quad (2)$$

$$e^\alpha{}_0(\xi^0) = \frac{\partial x^\alpha}{c\partial \xi^0} = ((1+\frac{a_0}{c^2}\xi^1)\cosh(\frac{a_0\xi^0}{c}),(1+\frac{a_0}{c^2}\xi^1)\sinh(\frac{a_0\xi^0}{c}),0,0) \quad (3)$$

About y -axis's and z -axis's orientation

$$e^\alpha{}_2(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^2} = (0,0,1,0) \quad , \quad e^\alpha{}_3(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^3} = (0,0,0,1) \quad (4)$$

The other unit vector $e^\alpha{}_1(\xi^0)$ is

$$e^\alpha{}_1(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^1} = (\sinh(\frac{a_0\xi^0}{c}),\cosh(\frac{a_0\xi^0}{c}),0,0) \quad (5)$$

Therefore,

$$cdt = c\cosh(\frac{a_0\xi^0}{c})d\xi^0(1+\frac{a_0}{c^2}\xi^1) + \sinh(\frac{a_0\xi^0}{c})d\xi^1$$

$$dx = c\sinh(\frac{a_0\xi^0}{c})d\xi^0(1+\frac{a_0}{c^2}\xi^1) + \cosh(\frac{a_0\xi^0}{c})d\xi^1, dy = d\xi^2, \ dz = d\xi^3 \quad (6)$$

The vector transformation is

$$V'^\mu = \frac{\partial x'^\mu}{\partial x^\alpha}V^\alpha, \qquad U'_\mu = \frac{\partial x^\alpha}{\partial x'^\mu}U_\alpha \quad (7)$$

Therefore, the transformation of the electro-magnetic 4-vector potential $(\phi,\vec{A}) = A^\alpha$ is

$$A^\alpha = \frac{\partial x^\alpha}{\partial x'^\mu}A'^\mu = \frac{\partial x^\alpha}{\partial \xi^\mu}A_\xi{}^\mu = e^\alpha{}_\mu A_\xi{}^\mu \ , \ e^\alpha{}_\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} \quad (8)$$

The transformation of differential coordinates is

$$dx^\alpha = \frac{\partial x^\alpha}{\partial x'^\mu}dx'^\mu = \frac{\partial x^\alpha}{\partial \xi^\mu}d\xi^\mu = e^\alpha{}_\mu d\xi^\mu \ , e^\alpha{}_\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} \quad (8\text{-i})$$

The equation of the electro-magnetic 4-vector potential $(\phi,\vec{A}) = A^\alpha$ is

$$(\frac{1}{c^2}\frac{\partial^2}{\partial t^2} - \nabla^2)\phi = 4\pi\rho$$

$$(\frac{1}{c^2}\frac{\partial^2}{\partial t^2} - \nabla^2)\vec{A} = \frac{4\pi}{c}\vec{j}$$

4-vector $(c\rho, \vec{j}) = \rho_0 \frac{dx^\alpha}{d\tau}$ (9)

Lorentz gauge transformation is in Rindler space-time,

$$A^\mu \to A^\mu + \partial^\mu \Lambda = A^\mu + g^{\mu\upsilon}\partial_\upsilon \Lambda \quad , \ \Lambda \text{ is a scalar function.}$$

$$g^{00} = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}, g^{11} = g^{22} = g^{33} = 1$$

$$\phi_\xi \to \phi_\xi - \frac{1}{c}\frac{\partial \Lambda}{\partial \xi^0}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}, \ \vec{A}_\xi \to \vec{A}_\xi + \vec{\nabla}_\xi \Lambda, \ \Lambda \text{ is a scalar function.} \tag{10}$$

Lorentz gauge is in Rindler space-time,

$$A^\mu{}_{;\mu} = \frac{\partial A^\mu}{\partial \xi^\mu} + \Gamma^\mu{}_{\mu\rho} A^\rho,$$

$$\Gamma^\mu{}_{\mu\rho} = \Gamma^0{}_{01} = \frac{1}{2} g^{00}(\frac{\partial g_{00}}{\partial \xi^1}) = \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}$$

$$g^{00} = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}, g^{11} = g^{22} = g^{33} = 1$$

$$0 = \frac{\partial \phi}{c\partial t} + \vec{\nabla}\cdot\vec{A} = \frac{\partial \phi_\xi}{c\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi + \frac{A_{\xi^1} a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \tag{11}$$

Hence, Lorentz gauge transformation and Lorentz gauge are in Rindler space-time,

$$A^\mu{}_{;\mu} = \frac{\partial A^\mu}{\partial \xi^\mu} + \Gamma^\mu{}_{\mu\rho}A^\rho \to \partial_\mu (A^\mu + g^{\mu\upsilon}\partial_\upsilon \Lambda) + \Gamma^0{}_{01}(A^1 + \frac{\partial \Lambda}{\partial \xi^1})$$

$$= \partial_\mu A^\mu + (g^{\mu\upsilon}\partial_\mu \partial_\upsilon)\Lambda + \Gamma^0{}_{01}(A^1 + \frac{\partial \Lambda}{\partial \xi^1})$$

$$0 = \frac{\partial \phi_\xi}{c\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi + \frac{A_{\xi^1} a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}$$

$$\to \frac{1}{c}\frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi - [\frac{1}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}(\frac{\partial}{\partial \xi^0})^2 - \nabla_\xi{}^2]\Lambda$$

$$+\frac{A_{\xi^1}a_0}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}+\frac{\partial\Lambda}{\partial\xi^1}\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}=0 \tag{11-i}$$

Hence, Lorentz gauge fixing condition is in Rindler space-time,

$$[\frac{1}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}(\frac{\partial}{\partial\xi^0})^2-\nabla_\xi{}^2]\Lambda-\frac{\partial\Lambda}{\partial\xi^1}\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}=0 \tag{12}$$

Hence, the transformation of the electro-magnetic 4-vector potential $(\phi,\vec{A})$ in inertial frame and the electro-magnetic 4-vector potential $(\phi_\xi,\vec{A}_\xi)$ in uniformly accelerated frame is

$$\phi=\cosh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\phi_\xi+\sinh(\frac{a_0\xi^0}{c})A_{\xi^1}$$
$$A_x=\sinh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\phi_\xi+\cosh(\frac{a_0\xi^0}{c})A_{\xi^1}$$
$$A_y=A_{\xi^2},A_z=A_{\xi^3} \tag{13}$$

If we make the matrix of the transformation of the differential coordinate,

$$\begin{pmatrix}cdt\\dx\\dy\\dz\end{pmatrix}=\begin{pmatrix}\cosh(\frac{a_0\xi^0}{c})(1+\frac{a_0\xi^1}{c^2}) & \sinh(\frac{a_0\xi^0}{c}) & 0 & 0\\ \sinh(\frac{a_0\xi^0}{c})(1+\frac{a_0\xi^1}{c^2}) & \cosh(\frac{a_0\xi^0}{c}) & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & 1\end{pmatrix}\begin{pmatrix}cd\xi^0\\d\xi^1\\d\xi^2\\d\xi^3\end{pmatrix}$$

$$=A\begin{pmatrix}cd\xi^0\\d\xi^1\\d\xi^2\\d\xi^3\end{pmatrix} \tag{14}$$

If we make the inverse matrix of Eq(14),

$$e_{\mu}{}^{\alpha} = \frac{\partial \xi^{\alpha}}{\partial x^{\mu}} = A^{-1} = \begin{pmatrix} \frac{c\partial \xi^0}{c\partial t} & \frac{c\partial \xi^0}{\partial x} & \frac{c\partial \xi^0}{\partial y} & \frac{c\partial \xi^0}{\partial z} \\ \frac{\partial \xi^1}{c\partial t} & \frac{\partial \xi^1}{\partial x} & \frac{\partial \xi^1}{\partial y} & \frac{\partial \xi^1}{\partial z} \\ \frac{\partial \xi^2}{c\partial t} & \frac{\partial \xi^2}{\partial x} & \frac{\partial \xi^2}{\partial y} & \frac{\partial \xi^2}{\partial z} \\ \frac{\partial \xi^3}{c\partial t} & \frac{\partial \xi^3}{\partial x} & \frac{\partial \xi^3}{\partial y} & \frac{\partial \xi^3}{\partial z} \end{pmatrix}$$

$$= \begin{pmatrix} \frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})} & -\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})} & 0 & 0 \\ -\sinh(\frac{a_0\xi^0}{c}) & \cosh(\frac{a_0\xi^0}{c}) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \tag{15}$$

Hence, we can obtain the matrix of transformation of differential operation.

$$\begin{pmatrix} \frac{1}{c}\frac{\partial}{\partial t} \\ \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{pmatrix} = (A^{-1})^T \begin{pmatrix} \frac{1}{c}\frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{pmatrix} = (A^T)^{-1} \begin{pmatrix} \frac{1}{c}\frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{pmatrix}$$

$$= \begin{pmatrix} \frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})} & -\sinh(\frac{a_0\xi^0}{c}) & 0 & 0 \\ -\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})} & \cosh(\frac{a_0\xi^0}{c}) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \frac{1}{c}\frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{pmatrix} \tag{16}$$

Hence, the transformation of differential operation is

$$\frac{1}{c}\frac{\partial}{\partial t}=\frac{c\partial\xi^0}{c\partial t}\frac{1}{c}\frac{\partial}{\partial\xi^0}+\frac{\partial\xi^1}{c\partial t}\frac{\partial}{\partial\xi^1}$$

$$=\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}-\sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}$$

$$\frac{\partial}{\partial x}=\frac{c\partial\xi^0}{\partial x}\frac{1}{c}\frac{\partial}{\partial\xi^0}+\frac{\partial\xi^1}{\partial x}\frac{\partial}{\partial\xi^1}$$

$$=-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}$$

$$\frac{\partial}{\partial y}=\frac{\partial}{\partial\xi^2},\frac{\partial}{\partial z}=\frac{\partial}{\partial\xi^3} \qquad (17)$$

The differential operation satisfy the following equation.

$$\frac{1}{c^2}\frac{\partial^2}{\partial t^2}-\nabla^2=\frac{1}{c^2(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2-\nabla_\xi{}^2$$

$$\vec{\nabla}=(\frac{\partial}{\partial x},\frac{\partial}{\partial y},\frac{\partial}{\partial z}),\qquad \vec{\nabla}_\xi=(\frac{\partial}{\partial\xi^1},\frac{\partial}{\partial\xi^2},\frac{\partial}{\partial\xi^3}) \qquad (18)$$

3. Electro-magnetic Field in the Rindler space-time

The electro-magnetic field $(\vec{E},\vec{B})$ is in the inertial frame,

$$\vec{E}=-\vec{\nabla}\phi-\frac{\partial\vec{A}}{c\partial t},\vec{B}=\vec{\nabla}\times\vec{A} \qquad (19)$$

We have to calculate for define electro-magnetic field in Rindler space-time. We have to calculate for the electro-magnetic field transformation in Rindler space-time. We will straightforward calculate it by the electro-magnetic 4-vector potential transformation, Eq(13) and the transformation of differential operation, Eq(17).

The x-component E_x of electric field $\vec{E}$ is in the inertial frame,

$$E_x=-\frac{\partial\phi}{\partial x}-\frac{\partial A_x}{c\partial t}$$

$$= -[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0} + \cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[\cosh(\frac{a_0\xi^0}{c})(1+\frac{a_0\xi^1}{c^2})\phi_\xi + \sinh(\frac{a_0\xi^0}{c})A_{\xi^1}]$$

$$-[\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0} - \sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[\sinh(\frac{a_0\xi^0}{c})(1+\frac{a_0\xi^1}{c^2})\phi_\xi + \cosh(\frac{a_0\xi^0}{c})A_{\xi^1}]$$

$$= -\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial A_{\xi^1}}{c\partial\xi^0} - (1+\frac{a_0\xi^1}{c^2})\frac{\partial\phi_\xi}{\partial\xi^1} - 2\phi_\xi\frac{a_0}{c^2}$$

$$= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{\partial\xi^1}[(1+\frac{a_0}{c^2}\xi^1)^2\phi_\xi] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial A_{\xi^1}}{c\partial\xi^0} \qquad (20)$$

The y-component E_y of electric field $\vec{E}$ is in the inertial frame,

$$E_y = -\frac{\partial\phi}{\partial y} - \frac{\partial A_y}{c\partial t} = -\frac{\partial}{\partial\xi^2}[\cosh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\phi_\xi + \sinh(\frac{a_0\xi^0}{c})A_{\xi^1}]$$

$$-[\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0} - \sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]A_{\xi^2}$$

$$= -(1+\frac{a_0\xi^1}{c^2})\cosh(\frac{a_0\xi^0}{c})\frac{\partial\phi_\xi}{\partial\xi^2} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\cosh(\frac{a_0\xi^0}{c})\frac{\partial A_{\xi^2}}{c\partial\xi^0}$$

$$+\sinh(\frac{a_0}{c}\xi^0)[\frac{\partial A_{\xi^2}}{\partial\xi^1} - \frac{\partial A_{\xi^1}}{\partial\xi^2}]$$

$$= \cosh(\frac{a_0}{c}\xi^0)[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^2}[\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial A_{\xi^2}}{c\partial\xi^0}]$$

$$+\sinh(\frac{a_0}{c}\xi^0)[\frac{\partial A_{\xi^2}}{\partial\xi^1} - \frac{\partial A_{\xi^1}}{\partial\xi^2}] \qquad (21)$$

The z-component E_z of electric field $\vec{E}$ is in the inertial frame,

$$E_z = -\frac{\partial \phi}{\partial z} - \frac{\partial A_z}{c\partial t} = -\frac{\partial}{\partial \xi^3}\,[\cosh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\,\xi^1)\phi_\xi + \sinh(\frac{a_0\xi^0}{c})A_{\xi^1}]$$

$$-[\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial \xi^0} - \sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial \xi^1}]A_{\xi^3}$$

$$= -(1+\frac{a_0\xi^1}{c^2})\cosh(\frac{a_0\xi^0}{c})\frac{\partial \phi_\xi}{\partial \xi^3} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\cosh(\frac{a_0\xi^0}{c})\frac{\partial A_{\xi^3}}{c\partial \xi^0}$$

$$+\sinh(\frac{a_0}{c}\,\xi^0)\,[\frac{\partial A_{\xi^3}}{\partial \xi^1} - \frac{\partial A_{\xi^1}}{\partial \xi^3}]$$

$$= \cosh(\frac{a_0}{c}\,\xi^0)\,[-\frac{1}{(1+\frac{a_0}{c^2}\,\xi^1)}\frac{\partial}{\partial \xi^3}[\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial A_{\xi^3}}{c\partial \xi^0}]$$

$$+\sinh(\frac{a_0}{c}\,\xi^0)\,[\frac{\partial A_{\xi^3}}{\partial \xi^1} - \frac{\partial A_{\xi^1}}{\partial \xi^3}]$$

(22)

The x-component B_x of magnetic field $\vec{B}$ is in the inertial frame,

$$B_x = \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} = \frac{\partial A_{\xi^3}}{\partial \xi^2} - \frac{\partial A_{\xi^2}}{\partial \xi^3} \qquad (23)$$

The y-component B_y of magnetic field $\vec{B}$ is in the inertial frame,

$$B_y = \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} = \frac{\partial A_x}{\partial \xi^3} - \frac{\partial A_{\xi^3}}{\partial x}$$

$$= \frac{\partial}{\partial \xi^3}\,[\sinh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\,\xi^1)\phi_\xi + \cosh(\frac{a_0\xi^0}{c})A_{\xi^1}]$$

$$-[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]A_{\xi^3}$$

$$=\cosh(\frac{a_0}{c}\xi^0)[\frac{\partial A_{\xi^1}}{\partial\xi^3}-\frac{\partial A_{\xi^3}}{\partial\xi^1}]$$

$$-\sinh(\frac{a_0}{c}\xi^0)[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^3}[\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2]-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial A_{\xi^3}}{c\partial\xi^0}] \quad (24)$$

The z-component B_z of magnetic field $\vec{B}$ is in the inertial frame,

$$B_z=\frac{\partial A_y}{\partial x}-\frac{\partial A_x}{\partial y}=\frac{\partial A_{\xi^2}}{\partial x}-\frac{\partial A_x}{\partial\xi^2}$$

$$=[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]A_{\xi^3}$$

$$-\frac{\partial}{\partial\xi^2}[\sinh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\phi_\xi+\cosh(\frac{a_0\xi^0}{c})A_{\xi^1}]$$

$$=\cosh(\frac{a_0}{c}\xi^0)[\frac{\partial A_{\xi^2}}{\partial\xi^1}-\frac{\partial A_{\xi^1}}{\partial\xi^2}]$$

$$+\sinh(\frac{a_0}{c}\xi^0)[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^2}[\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2]-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial A_{\xi^2}}{c\partial\xi^0}] \quad (25)$$

Hence, we can define the electro-magnetic field $(\vec{E}_\xi,\vec{B}_\xi)$ in Rindler space-time.

$$\vec{E}_\xi=-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial\vec{A}_\xi}{c\partial\xi^0}$$

$$\vec{B}_\xi=\vec{\nabla}_\xi\times\vec{A}_\xi$$

In this time, $\vec{\nabla}_{\xi} = (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3}), \vec{A}_{\xi} = (A_{\xi^1}, A_{\xi^2}, A_{\xi^3})$ (26)

We obtain the transformation of the electro-magnetic field.

$$E_x = -\frac{1}{(1+\frac{a_0 \xi^1}{c^2})}\frac{\partial}{\partial \xi^1}\{\phi_{\xi}(1+\frac{a_0\xi^1}{c^2})^2\} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial A_{\xi^1}}{c\partial \xi^0} = E_{\xi^1} \ ,$$

$$E_y = E_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) + B_{\xi^3}\sinh(\frac{a_0\xi^0}{c}),$$

$$E_z = E_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) - B_{\xi^2}\sinh(\frac{a_0\xi^0}{c})$$

$$B_x = B_{\xi^1},$$

$$B_y = B_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) - E_{\xi^3}\sinh(\frac{a_0\xi^0}{c})$$

$$B_z = B_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) + E_{\xi^2}\sinh(\frac{a_0\xi^0}{c}) \qquad (27)$$

Hence, we make the matrix of the transformation of the electro-magnetic field.

$$E_x = E_{\xi^1}, B_x = B_{\xi^1},$$

$$\begin{pmatrix} E_y \\ B_y \\ E_z \\ B_z \end{pmatrix} = H \begin{pmatrix} E_{\xi^2} \\ B_{\xi^2} \\ E_{\xi^3} \\ B_{\xi^3} \end{pmatrix}$$

$$H = \begin{pmatrix} \cosh(\frac{a_0\xi^0}{c}) & 0 & 0 & \sinh(\frac{a_0\xi^0}{c}) \\ 0 & \cosh(\frac{a_0\xi^0}{c}) & -\sinh(\frac{a_0\xi^0}{c}) & 0 \\ 0 & -\sinh(\frac{a_0\xi^0}{c}) & \cosh(\frac{a_0\xi^0}{c}) & 0 \\ \sinh(\frac{a_0\xi^0}{c}) & 0 & 0 & \cosh(\frac{a_0\xi^0}{c}) \end{pmatrix} \qquad (28)$$

We can obtain the matrix of the inverse-transformation of the electro-magnetic field.

$$E_{\xi^1} = E_x, B_{\xi^1} = B_x$$

$$\begin{pmatrix} E_{\xi^2} \\ B_{\xi^2} \\ E_{\xi^3} \\ B_{\xi^3} \end{pmatrix} = H^{-1} \begin{pmatrix} E_y \\ B_y \\ E_z \\ B_z \end{pmatrix}$$

$$H^{-1} = \begin{pmatrix} \cosh(\frac{a_0\xi^0}{c}) & 0 & 0 & -\sinh(\frac{a_0\xi^0}{c}) \\ 0 & \cosh(\frac{a_0\xi^0}{c}) & \sinh(\frac{a_0\xi^0}{c}) & 0 \\ 0 & \sinh(\frac{a_0\xi^0}{c}) & \cosh(\frac{a_0\xi^0}{c}) & 0 \\ -\sinh(\frac{a_0\xi^0}{c}) & 0 & 0 & \cosh(\frac{a_0\xi^0}{c}) \end{pmatrix} \quad (29)$$

(See also Ref [11]) Hence, the inverse-transformation of the electro-magnetic field is

$$E_{\xi^1} = E_x \, , B_{\xi^1} = B_x$$

$$E_{\xi^2} = E_y \cosh(\frac{a_0\xi^0}{c}) - B_z \sinh(\frac{a_0\xi^0}{c}),$$

$$B_{\xi^2} = B_y \cosh(\frac{a_0\xi^0}{c}) + E_z \sinh(\frac{a_0\xi^0}{c})$$

$$E_{\xi^3} = E_z \cosh(\frac{a_0\xi^0}{c}) + B_y \sinh(\frac{a_0\xi^0}{c})$$

$$B_{\xi^3} = B_z \cosh(\frac{a_0\xi^0}{c}) - E_y \sinh(\frac{a_0\xi^0}{c}) \quad (30)$$

If we apply Lorentz gauge transformation, Eq(10) to electro-magnetic field, Eq(26) in Rindler space-time

$$\phi_\xi \to \phi_\xi - \frac{1}{c}\frac{\partial\Lambda}{\partial\xi^0}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}, \quad \vec{A}_\xi \to \vec{A}_\xi + \vec{\nabla}_\xi\Lambda, \quad \Lambda \text{ is a scalar function.} \quad (31)$$

$$\vec{E}_\xi = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} + \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\frac{\partial\Lambda}{c\partial\xi^0}$$

$$-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial\vec{A}_\xi}{c\partial\xi^0} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}\vec{\nabla}_\xi\Lambda$$

$$= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial\vec{A}_\xi}{c\partial\xi^0}$$

$$\vec{B}_\xi = \vec{\nabla}_\xi\times\vec{A}_\xi + \vec{\nabla}_\xi\times\vec{\nabla}_\xi\Lambda = \vec{\nabla}_\xi\times\vec{A}_\xi \qquad (32)$$

If we apply Lorentz gauge transformation, Eq(10) to the transformation of the electro-magnetic 4-vector potential, Eq(13),

$$\phi - \frac{1}{c}\frac{\partial\Lambda}{\partial t} = \cosh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\{\phi_\xi - \frac{1}{c}\frac{\partial\Lambda}{\partial\xi^0}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\}$$

$$+\sinh(\frac{a_0\xi^0}{c})(A_{\xi^1} + \frac{\partial\Lambda}{\partial\xi^1})$$

$$A_x + \frac{\partial\Lambda}{\partial x} = \sinh(\frac{a_0\xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\{\phi_\xi - \frac{1}{c}\frac{\partial\Lambda}{\partial\xi^0}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\}$$

$$+\cosh(\frac{a_0\xi^0}{c})(A_{\xi^1} + \frac{\partial\Lambda}{\partial\xi^1})$$

$$A_y + \frac{\partial\Lambda}{\partial y} = A_{\xi^2} + \frac{\partial\Lambda}{\partial\xi^2}, A_z + \frac{\partial\Lambda}{\partial z} = A_{\xi^3} + \frac{\partial\Lambda}{\partial\xi^3} \qquad (33)$$

4. Electro-magnetic Field Equation (Maxwell Equation) in the Rindler space-time

Maxwell equation is in the inertial frame,

$$\vec{\nabla}\cdot\vec{E} = 4\pi\rho \qquad (34\text{-i})$$

$$\vec{\nabla}\times\vec{B} = \frac{\partial\vec{E}}{c\partial t} + \frac{4\pi}{c}\vec{j} \qquad (34\text{-ii})$$

$$\vec{\nabla}\cdot\vec{B} = 0 \qquad (34\text{-iii})$$

$$\vec{\nabla}\times\vec{E} = -\frac{\partial\vec{B}}{c\partial t} \qquad (34\text{-iv})$$

We will calculation for discovering Maxwell equations in Rindler space-time. For this purpose, we will straightforward calculate it by the electro-magnetic field transformation, Eq(27) and the transformation of differential operation, Eq(17).

At first, the transformation of 4-vector $(c\rho,\vec{j}) = \rho_0\frac{dx^\alpha}{d\tau}$ is

$$\rho = \rho_\xi (1+\frac{a_0\xi^1}{c^2})\cosh(\frac{a_0\xi^0}{c}) + \frac{j_{\xi^1}}{c}\sinh(\frac{a_0\xi^0}{c})$$

$$j_x = j_{\xi^1}\cosh(\frac{a_0\xi^0}{c}) + c\rho_\xi(1+\frac{a_0}{c^2}\xi^1)\sinh(\frac{a_0\xi^0}{c}), \quad j_y = j_{\xi^2}, j_z = j_{\xi^3} \qquad (35)$$

In this time, 4-vector $(c\rho_\xi, \vec{j}_\xi) = \rho_0 \frac{d\xi^\alpha}{d\tau}$ is defined in Rindler space-time.

Maxwell equation's first law is in the inertial frame,

1. $\vec{\nabla}\cdot\vec{E} = 4\pi\rho$

We write transformation of the electric field for easy calculating.

$E_x = E_{\xi^1}$,

$$E_y = E_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) + B_{\xi^3}\sinh(\frac{a_0\xi^0}{c}),$$

$$E_z = E_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) - B_{\xi^2}\sinh(\frac{a_0\xi^0}{c})$$

$$4\pi\rho = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z}$$

$$= [-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0} + \cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]E_{\xi^1}$$

$$+\frac{\partial}{\partial\xi^2}[E_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) + B_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

$$+\frac{\partial}{\partial\xi^3}[E_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) - B_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

$$= \cosh(\frac{a_0}{c}\xi^0)(\vec{\nabla}_\xi\cdot\vec{E}_\xi) + \sinh(\frac{a_0}{c}\xi^0)[\frac{\partial B_{\xi^3}}{\partial\xi^2} - \frac{\partial B_{\xi^2}}{\partial\xi^3} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial E_{\xi^1}}{c\partial\xi^0}] \qquad (36)$$

Maxwell equation's second law is in the inertial frame,

2. $\vec{\nabla}\times\vec{B} = \frac{\partial\vec{E}}{c\partial t} + \frac{4\pi}{c}\vec{j}$

We write the transformation of the magnetic field for easy calculating.

$$B_x = B_{\xi^1}$$

$$B_y = B_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) - E_{\xi^3}\sinh(\frac{a_0\xi^0}{c})$$

$$B_z = B_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) + E_{\xi^2}\sinh(\frac{a_0\xi^0}{c})$$

The X-component of Maxwell equation's second law is in the inertial frame,

X-component) $\frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z}$

$$= \frac{\partial}{\partial \xi^2}[B_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) + E_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

$$- \frac{\partial}{\partial \xi^3}[B_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) - E_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

$$= \cosh(\frac{a_0}{c}\xi^0)[\frac{\partial B_{\xi^3}}{\partial \xi^2} - \frac{\partial B_{\xi^2}}{\partial \xi^3}] + \sinh(\frac{a_0\xi^0}{c})[\frac{\partial E_{\xi^2}}{\partial \xi^2} + \frac{\partial E_{\xi^3}}{\partial \xi^3}]$$

$$= \frac{\partial E_x}{c\partial t} + \frac{4\pi}{c} j_x$$

$$= [\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial \xi^0} - \sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial \xi^1}]E_{\xi^1} + \frac{4\pi}{c} j_x$$

Hence,

$$\frac{4\pi}{c} j_x$$

$$= \sinh(\frac{a_0\xi^0}{c})(\vec{\nabla}_\xi \cdot \vec{E}_\xi) + \cosh(\frac{a_0\xi^0}{c})[\frac{\partial B_{\xi^3}}{\partial \xi^2} - \frac{\partial B_{\xi^2}}{\partial \xi^3} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial E_{\xi^1}}{c\partial \xi^0}] \quad (37)$$

The y-component of Maxwell equation's second law is in the inertial frame,

Y-component) $\frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x}$

$$= \frac{\partial B_{\xi^1}}{\partial \xi^3}$$

$$-[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[B_{\xi^3}\cosh(\frac{a_0\xi^0}{c})+E_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

$$=\frac{\partial E_y}{c\partial t}+\frac{4\pi}{c}j_y$$

$$=[\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}-\sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[E_{\xi^2}\cosh(\frac{a_0\xi^0}{c})+B_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

$$+\frac{4\pi}{c}j_y$$

Hence,

$$\frac{4\pi}{c}j_y=\frac{\partial B_{\xi^1}}{\partial\xi^3}-\frac{\partial B_{\xi^3}}{\partial\xi^1}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}B_{\xi^3}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial E_{\xi^2}}{c\partial\xi^0}$$

$$=\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^3}\{B_{\xi^1}(1+\frac{a_0}{c^2}\xi^1)\}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^1}\{B_{\xi^3}(1+\frac{a_0\xi^1}{c^2})\}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial E_{\xi^2}}{c\partial\xi^0}$$

(38)

The Z-component of Maxwell equation's second law is in the inertial frame,

Z-component) $\frac{\partial B_y}{\partial x}-\frac{\partial B_x}{\partial y}$

$$=[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[B_{\xi^2}\cosh(\frac{a_0\xi^0}{c})-E_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

$$-\frac{\partial B_{\xi^1}}{\partial\xi^2}$$

$$=\frac{\partial E_z}{c\partial t}+\frac{4\pi}{c}j_z$$

$$= [\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0} - \sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[E_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) - B_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

$$+\frac{4\pi}{c}j_z$$

Therefore,

$$\frac{4\pi}{c}j_z = \frac{\partial B_{\xi^2}}{\partial\xi^1} - \frac{\partial B_{\xi^1}}{\partial\xi^2} + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}B_{\xi^2} - \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial E_{\xi^3}}{c\partial\xi^0}$$

$$= \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^1}\{B_{\xi^2}(1+\frac{a_0}{c^2}\xi^1)\} - \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^2}\{B_{\xi^1}(1+\frac{a_0\xi^1}{c^2})\} - \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial E_{\xi^3}}{c\partial\xi^0}$$

(39)

Maxwell equation's third law is in the inertial frame,

3. $\vec{\nabla}\cdot\vec{B} = 0$

$$\vec{\nabla}\cdot\vec{B} = \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z}$$

$$= [-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0} + \cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]B_{\xi^1}$$

$$+\frac{\partial}{\partial\xi^2}[B_{\xi^2}\cosh(\frac{a_0\xi^0}{c}) - E_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

$$+\frac{\partial}{\partial\xi^3}[B_{\xi^3}\cosh(\frac{a_0\xi^0}{c}) + E_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

$$= \cosh(\frac{a_0\xi^0}{c})(\vec{\nabla}_\xi\cdot\vec{B}_\xi) + \sinh(\frac{a_0\xi^0}{c})[-(-\frac{\partial E_{\xi^2}}{\partial\xi^3} + \frac{\partial E_{\xi^3}}{\partial\xi^2}) - \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial B_{\xi^1}}{c\partial\xi^0}] = 0$$

(40)

Maxwell equation's fourth law is in the inertial frame,

4. $\vec{\nabla}\times\vec{E} = -\frac{\partial\vec{B}}{c\partial t}$

We write the transformation of the electric field for easy calculating.

$$E_x = E_{\xi^1} \quad ,$$

$$E_y = E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c}),$$

$$E_z = E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})$$

The x-component of Maxwell equation's fourth law is in the inertial frame,

X-component) $\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z}$

$$= \frac{\partial}{\partial \xi^2} [E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})]$$

$$- \frac{\partial}{\partial \xi^3} [E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})]$$

$$= \cosh(\frac{a_0}{c} \xi^0) [\frac{\partial E_{\xi^3}}{\partial \xi^2} - \frac{\partial E_{\xi^2}}{\partial \xi^3}] - \sinh(\frac{a_0 \xi^0}{c}) [\frac{\partial B_{\xi^2}}{\partial \xi^2} + \frac{\partial B_{\xi^3}}{\partial \xi^3}]$$

$$= -\frac{\partial B_x}{c \partial t}$$

$$= -[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}] B_{\xi^1}$$

Hence,

$$-\sinh(\frac{a_0 \xi^0}{c})(\vec{\nabla}_\xi \cdot \vec{B}_\xi) + \cosh(\frac{a_0 \xi^0}{c}) [(\frac{\partial E_{\xi^3}}{\partial \xi^2} - \frac{\partial E_{\xi^2}}{\partial \xi^3}) + \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial B_{\xi^1}}{c \partial \xi^0}] = 0 \qquad (41)$$

The y-component of Maxwell equation's fourth law is in the inertial frame,

Y-component) $\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x}$

$$= \frac{\partial E_{\xi^1}}{\partial \xi^3}$$

$$-[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[E_{\xi^3}\cosh(\frac{a_0\xi^0}{c})-B_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

$$=-\frac{\partial B_y}{c\partial t}$$

$$=-[\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}-\sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[B_{\xi^2}\cosh(\frac{a_0\xi^0}{c})-E_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

Therefore,

$$\frac{\partial E_{\xi^1}}{\partial\xi^3}-\frac{\partial E_{\xi^3}}{\partial\xi^1}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}E_{\xi^3}+\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial B_{\xi^2}}{c\partial\xi^0}$$

$$=\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^3}\{E_{\xi^1}(1+\frac{a_0}{c^2}\xi^1)\}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial\xi^1}\{E_{\xi^3}(1+\frac{a_0\xi^1}{c^2})\}+\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial B_{\xi^2}}{c\partial\xi^0}$$

$$=0$$

(42)

The Z-component of Maxwell equation's fourth law is in the inertial frame,

$$\text{Z-component)}\ \frac{\partial E_y}{\partial x}-\frac{\partial E_x}{\partial y}$$

$$=[-\frac{\sinh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}+\cosh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[E_{\xi^2}\cosh(\frac{a_0\xi^0}{c})+B_{\xi^3}\sinh(\frac{a_0\xi^0}{c})]$$

$$-\frac{\partial E_{\xi^1}}{\partial\xi^2}$$

$$=-\frac{\partial B_z}{c\partial t}$$

$$=-[\frac{\cosh(\frac{a_0\xi^0}{c})}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}-\sinh(\frac{a_0\xi^0}{c})\frac{\partial}{\partial\xi^1}]\cdot[B_{\xi^3}\cosh(\frac{a_0\xi^0}{c})+E_{\xi^2}\sinh(\frac{a_0\xi^0}{c})]$$

Hence,

$$\frac{\partial E_{\xi^2}}{\partial \xi^1} - \frac{\partial E_{\xi^1}}{\partial \xi^2} + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}E_{\xi^2} + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial B_{\xi^3}}{c\partial \xi^0}$$

$$= \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial \xi^1}\{E_{\xi^2}(1+\frac{a_0}{c^2}\xi^1)\} - \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial}{\partial \xi^2}\{E_{\xi^1}(1+\frac{a_0\xi^1}{c^2})\} + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{\partial B_{\xi^3}}{c\partial \xi^0}$$

$$= 0$$

(43)

Therefore, we obtain the electro-magnetic field equation by Eq (35)-Eq(43) in Rindler space-time. (See also Ref [13])

$$\vec{\nabla}_{\xi} \cdot \vec{E}_{\xi} = 4\pi\rho_{\xi}(1+\frac{a_0\xi^1}{c^2}) \quad \text{(44-i)}$$

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_{\xi} \times \{\vec{B}_{\xi}(1+\frac{a_0\xi^1}{c^2})\} = \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial \vec{E}_{\xi}}{c\partial \xi^0} + \frac{4\pi}{c}\vec{j}_{\xi} \quad \text{(44-ii)}$$

$$\vec{\nabla}_{\xi} \cdot \vec{B}_{\xi} = 0 \quad \text{(44-iii)}$$

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_{\xi} \times \{\vec{E}_{\xi}(1+\frac{a_0\xi^1}{c^2})\} = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial \vec{B}_{\xi}}{c\partial \xi^0} \quad \text{(44-iv)}$$

$$\vec{E}_{\xi} = (E_{\xi^1}, E_{\xi^2}, E_{\xi^3}), \vec{B}_{\xi} = (B_{\xi^1}, B_{\xi^2}, B_{\xi^3}),$$

$$\vec{\nabla}_{\xi} = (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3})$$

For instant, we think the spherical charge density ρ_{ξ} of a stationary accelerated frame is in a charged huge sphere.

$$t = \xi^0 = 0, 0 \le x = \xi^1, y = \xi^2, z = \xi^3 \le R,$$

$$\vec{E} = \vec{E}_{\xi} = \frac{Q}{R^3}\vec{r}, \;\; \rho = \rho_0 = \frac{Q}{V} = \frac{3Q}{4\pi R^3}, \;\; \vec{j} = \vec{j}_{\xi} = \vec{0}$$

$$\frac{3Q}{4\pi R^3}\frac{1}{(1+\frac{a_0 R}{c^2})} \le \rho_\xi = \frac{\rho}{(1+\frac{a_0}{c^2}\xi^1)} = \frac{3Q}{4\pi R^3}\frac{1}{(1+\frac{a_0 x}{c^2})} \le \frac{3Q}{4\pi R^3}$$

(44-v)

Generally, the continuity equation is in Rindler space-time,

$$0 = j^\mu{}_{;\mu} = \frac{\partial j^\mu}{\partial \xi^\mu} + \Gamma^\mu{}_{\mu\rho} j^\rho,$$

$$\Gamma^\mu{}_{\mu\rho} = \Gamma^0{}_{01} = \frac{1}{2} g^{00}(\frac{\partial g_{00}}{\partial \xi^1}) = \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}$$

$$g^{00} = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}, g^{11} = g^{22} = g^{33} = 1$$

$$0 = \frac{\partial \rho}{\partial t} + \vec{\nabla}\cdot\vec{j} = \frac{\partial \rho_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{j}_\xi + \frac{j_{\xi^1} a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \tag{45}$$

We treat Lorentz gauge transformation, Eq(10) about the electro-magnetic field equation, Eq(44-i)-Eq(44-iv) in Rindler space-time. We will straightforward calculate it by Lorentz gauge transformation, Eq(10) and Lorentz gauge, Eq(11), Lorentz gauge fixing condition, Eq(12) in Rindler space-time.

Eq(44-i) is

$$\vec{\nabla}_\xi \cdot \vec{E}_\xi = \vec{\nabla}_\xi \cdot \{-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial \vec{A}_\xi}{c\partial \xi^0}\}$$

$$= -\vec{\nabla}_\xi\{\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\}\cdot[\vec{\nabla}_\xi\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} + \frac{\partial \vec{A}_\xi}{c\partial \xi^0}]$$

$$-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} + \frac{\partial}{c\partial \xi^0}(\vec{\nabla}_\xi \cdot \vec{A}_\xi)]$$

$$= \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}[\frac{\partial}{\partial \xi^1}\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} + \frac{\partial A_{\xi^1}}{c\partial \xi^0}]$$

$$-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}$$

$$-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}[-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{A_{\xi^1}a_0}{c^2}],$$

Lorentz gauge is Rindler space-time

$$\frac{1}{c}\frac{\partial\phi_\xi}{\partial\xi^0}+\vec{\nabla}_\xi\cdot\vec{A}_\xi=-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}A_{\xi^1}$$

Hence,

$$=-\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}E_{\xi^1}-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{\partial A_{\xi^1}}{c\partial\xi^0}\frac{a_0}{c^2}$$

$$=4\pi\rho_\xi(1+\frac{a_0\xi^1}{c^2}) \qquad (46)$$

If we apply Lorentz gauge transformation to Eq (46),

$\phi_\xi\to\phi_\xi-\frac{1}{c}\frac{\partial\Lambda}{\partial\xi^0}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}$, $\vec{A}_\xi\to\vec{A}_\xi+\vec{\nabla}_\xi\Lambda$, Λ is a scalar function.

$$=-\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}E_{\xi^1}-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\frac{1}{c}\frac{\partial\Lambda}{\partial\xi^0}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{\partial A_{\xi^1}}{c\partial\xi^0}\frac{a_0}{c^2}+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{a_0}{c^2}\frac{\partial}{c\partial\xi^0}\frac{\partial\Lambda}{\partial\xi^1}$$

$$=-\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}E_{\xi^1}-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial}{c\partial\xi^0}\{[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\Lambda+\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial\Lambda}{\partial\xi^1}\}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{\partial A_{\xi^1}}{c\partial\xi^0}\frac{a_0}{c^2} \tag{47}$$

In this time, Lorentz gauge fixing condition, Eq(12) is in Rindler space-time,

$$[\frac{1}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}(\frac{\partial}{\partial\xi^0})^2-\nabla_\xi{}^2]\Lambda-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial\Lambda}{\partial\xi^1}=0 \tag{48}$$

Hence, Eq(44-i) is

$$\vec{\nabla}_\xi\cdot\vec{E}_\xi$$

$$=-\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}E_{\xi^1}-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}[\nabla_\xi{}^2-\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{\partial A_{\xi^1}}{c\partial\xi^0}\frac{a_0}{c^2}$$

$$=4\pi\rho_\xi(1+\frac{a_0\xi^1}{c^2}) \tag{49}$$

Eq(44-i) is invariant about Lorentz gauge transformation in Rindler space-time.

Eq (44-ii) is

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\times\{\vec{B}_\xi(1+\frac{a_0\xi^1}{c^2})\}$$

$$=\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi\times\{\vec{\nabla}_\xi\times\vec{A}_\xi(1+\frac{a_0\xi^1}{c^2})\}$$

$$=\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi(1+\frac{a_0}{c^2}\xi^1)\times\{\vec{\nabla}_\xi\times\vec{A}_\xi\}+\vec{\nabla}_\xi\times\vec{\nabla}_\xi\times\vec{A}_\xi$$

$$=\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}(1,0,0)\times\vec{B}_\xi+\{-\nabla_\xi{}^2\vec{A}_\xi+\vec{\nabla}_\xi(\vec{\nabla}_\xi\cdot\vec{A}_\xi)\}$$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}(0,-B_{\xi^3},B_{\xi^2}) + \{-\nabla_\xi{}^2\vec{A}_\xi + \vec{\nabla}_\xi(\vec{\nabla}_\xi\cdot\vec{A}_\xi)\}$$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial\vec{E}_\xi}{c\partial\xi^0} + \frac{4\pi}{c}\vec{j}_\xi$$

$$= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{\partial}{c\partial\xi^0}[\vec{\nabla}_\xi\{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\}] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\vec{A}_\xi + \frac{4\pi\vec{j}_\xi}{c}$$

$$= -\frac{\partial}{c\partial\xi^0}\vec{\nabla}_\xi\phi_\xi - \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{2a_0}{c^2}\frac{\partial\phi_\xi}{c\partial\xi^0}(1,0,0) - \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\vec{A}_\xi + \frac{4\pi\vec{j}_\xi}{c}$$

(50)

Therefore,

$$\frac{4\pi}{c}\vec{j}_\xi$$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}(0,-B_{\xi^3},B_{\xi^2}) + \{-\nabla_\xi{}^2\vec{A}_\xi + \vec{\nabla}_\xi(\vec{\nabla}_\xi\cdot\vec{A}_\xi)\}$$

$$+ \frac{\partial}{c\partial\xi^0}\vec{\nabla}_\xi\phi_\xi + \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{2a_0}{c^2}\frac{\partial\phi_\xi}{c\partial\xi^0}(1,0,0) + \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\vec{A}_\xi$$

$$= \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(0,-B_{\xi^3},B_{\xi^2}) + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{2a_0}{c^2}\frac{\partial\phi_\xi}{c\partial\xi^0}(1,0,0)$$

$$+ [-\nabla_\xi{}^2 + \frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{A}_\xi + \vec{\nabla}_\xi[-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}A_{\xi^1}] \quad (51)$$

Lorentz gauge is Rindler space-time

$$\frac{1}{c}\frac{\partial\phi_\xi}{\partial\xi^0} + \vec{\nabla}_\xi\cdot\vec{A}_\xi = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}A_{\xi^1}$$

Hence,

$$\frac{4\pi}{c}\vec{j}_{\xi}$$

$$= \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(0,-B_{\xi^3},B_{\xi^2}) + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{2a_0}{c^2}\frac{\partial\phi_{\xi}}{c\partial\xi^0}(1,0,0)$$

$$+[-\nabla_{\xi}{}^2 + \frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{A}_{\xi}$$

$$+\vec{\nabla}_{\xi}[-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}A_{\xi^1}] \qquad (52)$$

If we apply Lorentz gauge transformation to Eq (52),

$\phi_{\xi} \to \phi_{\xi} - \frac{1}{c}\frac{\partial\Lambda}{\partial\xi^0}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}$, $\vec{A}_{\xi} \to \vec{A}_{\xi} + \vec{\nabla}_{\xi}\Lambda$, Λ is a scalar function.

$$\frac{4\pi}{c}\vec{j}_{\xi}$$

$$= \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(0,-B_{\xi^3},B_{\xi^2}) + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{2a_0}{c^2}\frac{\partial\phi_{\xi}}{c\partial\xi^0}(1,0,0)$$

$$-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^3}\frac{2a_0}{c^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\Lambda(1,0,0)$$

$$+[-\nabla_{\xi}{}^2 + \frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{A}_{\xi} + [-\nabla_{\xi}{}^2 + \frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{\nabla}_{\xi}\Lambda$$

$$+\vec{\nabla}_{\xi}[-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}A_{\xi^1}] + \vec{\nabla}_{\xi}[-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial\Lambda}{\partial\xi^1}]$$

$$= \frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(0,-B_{\xi^3},B_{\xi^2}) + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{2a_0}{c^2}\frac{\partial\phi_{\xi}}{c\partial\xi^0}(1,0,0)$$

$$-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^3}\frac{2a_0}{c^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\Lambda(1,0,0)$$

$$+[-\nabla_{\xi}{}^{2}+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{A}_{\xi}+[-\nabla_{\xi}{}^{2}+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{\nabla}_{\xi}\Lambda$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{a_0{}^2}{c^4}A_{\xi^1}(1,0,0)-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\vec{\nabla}_{\xi}A_{\xi^1}$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{a_0{}^2}{c^4}\frac{\partial\Lambda}{\partial\xi^1}(1,0,0)-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\vec{\nabla}_{\xi}\Lambda$$

(53)

In this time, Lorentz gauge fixing condition, Eq(12) is in Rindler space-time,

$$[\frac{1}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}(\frac{\partial}{\partial\xi^0})^2-\nabla_{\xi}{}^{2}]\Lambda-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial\Lambda}{\partial\xi^1}=0$$

$$0=\vec{\nabla}_{\xi}[\{-\nabla_{\xi}{}^{2}+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\}\Lambda]$$

$$=\vec{\nabla}_{\xi}\{\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}\}(\frac{\partial}{\partial\xi^0})^2\Lambda+[-\nabla_{\xi}{}^{2}+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{\nabla}_{\xi}\Lambda$$

$$-\vec{\nabla}_{\xi}\{\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\}\frac{\partial\Lambda}{\partial\xi^1}-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\vec{\nabla}_{\xi}\Lambda$$

$$=-\frac{2}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^3}\frac{a_0}{c^2}(\frac{\partial}{\partial\xi^0})^2\Lambda(1,0,0)+[-\nabla_{\xi}{}^{2}+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{\nabla}_{\xi}\Lambda$$

$$+\frac{a_0{}^2}{c^4}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}\frac{\partial\Lambda}{\partial\xi^1}(1,0,0)-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\vec{\nabla}_{\xi}\Lambda \qquad (54)$$

Therefore, Eq(53) is

$$\frac{4\pi}{c}\vec{j}_{\xi}$$

$$=\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(0,-B_{\xi^3},B_{\xi^2})+\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{2a_0}{c^2}\frac{\partial\phi_\xi}{c\partial\xi^0}(1,0,0)$$

$$-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^3}\frac{2a_0}{c^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\Lambda(1,0,0)$$

$$+[-\nabla_\xi{}^2+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{A}_\xi+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{a_0{}^2}{c^4}A_{\xi^1}(1,0,0)$$

$$-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\vec{\nabla}_\xi A_{\xi^1}+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{a_0{}^2}{c^4}\frac{\partial\Lambda}{\partial\xi^1}(1,0,0)-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\vec{\nabla}_\xi\Lambda$$

$$+\vec{\nabla}_\xi[\{-\nabla_\xi{}^2+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\}\Lambda]$$

$$+\frac{2}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^3}\frac{a_0}{c^2}(\frac{\partial}{\partial\xi^0})^2\Lambda(1,0,0)-\frac{a_0{}^2}{c^4}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}\frac{\partial\Lambda}{\partial\xi^1}(1,0,0)$$

$$+\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{a_0}{c^2}\frac{\partial}{\partial\xi^1}\vec{\nabla}_\xi\Lambda$$

$$=\frac{a_0}{c^2}\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(0,-B_{\xi^3},B_{\xi^2})+\frac{1}{(1+\frac{a_0}{c^2}\xi^1)}\frac{2a_0}{c^2}\frac{\partial\phi_\xi}{c\partial\xi^0}(1,0,0)$$

$$+[-\nabla_\xi{}^2+\frac{1}{c^2}\frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}(\frac{\partial}{\partial\xi^0})^2]\vec{A}_\xi$$

$$+\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{a_0{}^2}{c^4}A_{\xi^1}(1,0,0)-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{a_0}{c^2}\vec{\nabla}_\xi A_{\xi^1} \tag{55}$$

Hence, Eq(44-ii) is invariant about Lorentz gauge transformation in Rindler space-time.

Eq (44-iii) is

$$\vec{\nabla}_{\xi}\cdot\vec{B}_{\xi}=\vec{\nabla}_{\xi}\cdot(\vec{\nabla}_{\xi}\times\vec{A}_{\xi}+\vec{\nabla}_{\xi}\times\vec{\nabla}_{\xi}\Lambda)=\vec{\nabla}_{\xi}\times\vec{\nabla}_{\xi}\cdot\vec{A}_{\xi}=0 \tag{56}$$

Eq (44-iv) is

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_{\xi}\times\{\vec{E}_{\xi}(1+\frac{a_0\xi^1}{c^2})\}$$

$$=-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_{\xi}\times[\vec{\nabla}_{\xi}\{\phi_{\xi}(1+\frac{a_0\xi^1}{c^2})^2\}-\vec{\nabla}_{\xi}(\frac{\partial\Lambda}{c\partial\xi^0})+\frac{\partial\vec{A}_{\xi}}{c\partial\xi^0}+\frac{\partial}{c\partial\xi^0}(\vec{\nabla}_{\xi}\Lambda)]$$

$$=-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_{\xi}\times\frac{\partial\vec{A}_{\xi}}{c\partial\xi^0}=-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial(\vec{\nabla}_{\xi}\times\vec{A}_{\xi})}{c\partial\xi^0}=-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial\vec{B}_{\xi}}{c\partial\xi^0} \tag{57}$$

Hence, Eq (44-iii), Eq (44-iv) are invariant about Lorentz gauge transformation in Rindler space-time. Hence, the electro-magnetic field equations (Maxwell Equations) in Rindler space-time, Eq(44-i)-Eq(44-iv) are invariant about Lorentz gauge transformation.

5. Conclusion

We find the electro-magnetic field transformation and the electro-magnetic equation in uniformly accelerated frame in one theory.

Generally, the coordinate transformation of accelerated frame is (see Ref [9])

(I) $$ct=(\frac{c^2}{a_0}+\xi^1)\sinh(\frac{a_0\xi^0}{c})$$

$$x=(\frac{c^2}{a_0}+\xi^1)\cosh(\frac{a_0\xi^0}{c})-\frac{c^2}{a_0}\quad ,y=\xi^2,z=\xi^3 \tag{58}$$

(II) $$ct=\frac{c^2}{a_0}\exp(\frac{a_0}{c^2}\xi^1)\sinh(\frac{a_0\xi^0}{c})$$

$$x=\frac{c^2}{a_0}\exp(\frac{a_0}{c^2}\xi^1)\cosh(\frac{a_0\xi^0}{c})-\frac{c^2}{a_0},y=\xi^2,z=\xi^3 \tag{59}$$

If you try to use Eq(59) for make Maxwell equation in Rindler space-time, You have to fail it.

In A.Einstein's article (see Ref [10]), Einstein obtain Lorenz transformation by Maxwell equation in inertial frame, Einstein give up Galilei transformation in inertial frame. In accelerated frame, we think our article's choice is Rindler coordinate (I) can treat electro-magnetic field equation likely Einstein's election.

APPENDIX A

In 2-Dimension Rindler space-time, if we mistake calculation, we can think the electro-magnetic wave function.

$$E_{\xi^1} = E_x = E_{X0} \sin\omega(t - \frac{x}{c})$$
$$= E_{x0} \sin\omega[-(\{\frac{c}{a_0} + \frac{\xi^1}{c}\}\exp(-\frac{a_0}{c}\xi^0) + \frac{c}{a_0})] = E_{x0}\sin\Phi \qquad \text{(A-1)}$$

$$ct = (\frac{c^2}{a_0} + \xi^1)\sinh(\frac{a_0\xi^0}{c}) \ , \quad x = (\frac{c^2}{a_0} + \xi^1)\cosh(\frac{a_0\xi^0}{c}) - \frac{c^2}{a_0}$$

In this time,

$$-\exp(-\frac{a_0}{c}\xi^0) = \sinh(\frac{a_0\xi^0}{c}) - \cosh(\frac{a_0\xi^0}{c}) \qquad \text{(A-2)}$$

(A-1) have to satisfy the following equation.

$$\frac{1}{(1+\frac{a_0\xi^1}{c})^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\sin\Phi = (\frac{\partial}{\partial\xi^1})^2\sin\Phi \qquad \text{(A-3)}$$

If we calculate Eq(A-3),

$$(\frac{\partial}{\partial\xi^1})^2\sin\Phi = \frac{\partial}{\partial\xi^1}\{\frac{\partial}{\partial\xi^1}\sin\omega[-\{\frac{c}{a_0} + \frac{\xi^1}{c}\}\exp(-\frac{a_0}{c}\xi^0) + \frac{c}{a_0}]\}$$
$$= \frac{\partial}{\partial\xi^1}\{(\cos\Phi)\cdot -\omega\frac{1}{c}\exp(-\frac{a_0}{c}\xi^0)\} = (-\sin\Phi)\cdot\frac{\omega^2}{c^2}\exp(-2\frac{a_0}{c}\xi^0)$$
$$\text{(A-4)}$$

But next calculation's situation is different.

$$\frac{1}{(1+\frac{a_0\xi^1}{c})^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\sin\Phi$$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}\frac{\partial}{\partial\xi^0}\{\frac{\partial}{\partial\xi^0}\sin\omega[-\{\frac{c}{a_0} + \frac{\xi^1}{c}\}\exp(-\frac{a_0}{c}\xi^0) + \frac{c}{a_0}]\}$$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}\frac{\partial}{\partial\xi^0}\{(\cos\Phi)\cdot\omega\{\frac{c}{a_0} + \frac{\xi^1}{c}\}\frac{a_0}{c}\exp(-\frac{a_0}{c}\xi^0)\}$$

In this time, $D(\alpha\beta) = \beta D\alpha + \alpha D\beta$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}\{\frac{\partial}{\partial\xi^0}(\cos\Phi)\cdot\omega\{\frac{c}{a_0}+\frac{\xi^1}{c}\}\frac{a_0}{c}\exp(-\frac{a_0}{c}\xi^0)$$

$$+(\cos\Phi)\cdot\omega\{\frac{c}{a_0}+\frac{\xi^1}{c}\}\frac{a_0}{c}\frac{\partial}{\partial\xi^0}(\exp(-\frac{a_0}{c}\xi^0))\}$$

$$= \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2}\frac{1}{c^2}\{-\sin\Phi\cdot\omega^2\{\frac{c}{a_0}+\frac{\xi^1}{c}\}^2\frac{{a_0}^2}{c^2}\exp(-2\frac{a_0}{c}\xi^0)$$

$$-(\cos\Phi)\cdot\omega\{\frac{c}{a_0}+\frac{\xi^1}{c}\}\frac{{a_0}^2}{c^2}\exp(-\frac{a_0}{c}\xi^0)\} \qquad \text{(A-5)}$$

Hence, if we compare Eq(A-4) and Eq(A-5),

$$(\frac{\partial}{\partial\xi^1})^2\sin\Phi = (-\sin\Phi)\cdot\frac{\omega^2}{c^2}\exp(-2\frac{a_0}{c}\xi^0)$$

$$\neq -\sin\Phi\cdot\omega^2\frac{1}{c^2}\exp(-2\frac{a_0}{c}\xi^0)$$

$$-\frac{1}{(1+\frac{a_0\xi^1}{c^2})}(\cos\Phi)\cdot\omega\frac{a_0}{c}\frac{1}{c^2}\exp(-\frac{a_0}{c}\xi^0) = \frac{1}{(1+\frac{a_0\xi^1}{c})^2}\frac{1}{c^2}(\frac{\partial}{\partial\xi^0})^2\sin\Phi \qquad \text{(A-6)}$$

Therefore, we think it cannot exist electro-magnetic wave function in Rindler space-time.

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