

**The Dirac Equation
is a Special Case of
the Maxwell-Cassano Equations**

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For vector Ψ_D The Klein-Gordon equation may be written [4]:

$$\left(\nabla^2 - \frac{\partial^2}{\partial t^2} - m^2\right)\Psi_D = \mathbf{0} .$$

Whenever $\Psi_D = \begin{pmatrix} \phi_D^A \\ \phi_D^B \end{pmatrix}$ is a 2^M -dimensional vector, via a matrix differential operator factorization, it may be written (in the Dirac representation):

$$\begin{pmatrix} \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) & i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ -i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) & -i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \phi_D^A \\ \phi_D^B \end{pmatrix} = \mathbf{0}$$

and, since these matrix operators are commutative:

$$\begin{pmatrix} \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) & -i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) & i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ -i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \phi_D^A \\ \phi_D^B \end{pmatrix} = \mathbf{0}$$

Where:

$$\boldsymbol{\sigma}\cdot\vec{\nabla} = \sum_{v=1}^3 \boldsymbol{\sigma}^v \frac{\partial}{\partial x^v}$$

$$\boldsymbol{\sigma}^0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \mathbf{I}_2$$

$$\boldsymbol{\sigma}^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \boldsymbol{\sigma}^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \boldsymbol{\sigma}^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\Rightarrow \boldsymbol{\sigma}\cdot\vec{\nabla} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \partial_1 + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \partial_2 + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \partial_3$$

$$= \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix}$$

$$\Rightarrow (\boldsymbol{\sigma}\cdot\vec{\nabla})^2 = \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} = (\partial_3^2 + \partial_1^2 + \partial_2^2)\mathbf{I}_2 = \nabla^2 \mathbf{I}_2$$

Let:

$$\begin{pmatrix} \Psi_D^A \\ \Psi_D^B \end{pmatrix} = \begin{pmatrix} \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) & -i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \phi_D^A \\ \phi_D^B \end{pmatrix}$$

$$\begin{pmatrix} \Psi_D^C \\ \Psi_D^D \end{pmatrix} = \begin{pmatrix} \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) & i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ -i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \phi_D^A \\ \phi_D^B \end{pmatrix}$$

then:

$$\begin{pmatrix} \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) & i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ -i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \Psi_D^A \\ \Psi_D^B \end{pmatrix} = \mathbf{0}$$

$$\begin{pmatrix} \mathbf{I}_2\left(-i\frac{\partial}{\partial t} + m\right) & -i\boldsymbol{\sigma}\cdot\vec{\nabla} \\ i\boldsymbol{\sigma}\cdot\vec{\nabla} & \mathbf{I}_2\left(i\frac{\partial}{\partial t} + m\right) \end{pmatrix} \begin{pmatrix} \Psi_D^C \\ \Psi_D^D \end{pmatrix} = \mathbf{0}$$

conformability of the matrices requires that:

$\phi_D^A \cdot \phi_D^B$; $\Psi_D^A \cdot \Psi_D^B \cdot \Psi_D^C \cdot \Psi_D^D$ are all 2×1 matrices; so setting:

Let:

$$\phi_D^A = \begin{pmatrix} \phi_D^0 \\ \phi_D^1 \end{pmatrix} \quad \phi_D^B = \begin{pmatrix} \phi_D^2 \\ \phi_D^3 \end{pmatrix}$$

$$\Psi_D^A = \begin{pmatrix} \psi_D^0 \\ \psi_D^1 \end{pmatrix} \quad \Psi_D^B = \begin{pmatrix} \psi_D^2 \\ \psi_D^3 \end{pmatrix} \quad \Psi_D^C = \begin{pmatrix} \psi_D^4 \\ \psi_D^5 \end{pmatrix} \quad \Psi_D^D = \begin{pmatrix} \psi_D^6 \\ \psi_D^7 \end{pmatrix}$$

then:

$$\begin{aligned}
& \left(\begin{array}{cc} \left(\begin{array}{cc} (i\frac{\partial}{\partial t} + m) & 0 \\ 0 & (i\frac{\partial}{\partial t} + m) \end{array} \right) & i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} \\ -i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \left(\begin{array}{cc} (-i\frac{\partial}{\partial t} + m) & 0 \\ 0 & (-i\frac{\partial}{\partial t} + m) \end{array} \right) \end{array} \right) \begin{pmatrix} \begin{pmatrix} \psi_D^0 \\ \psi_D^1 \end{pmatrix} \\ \begin{pmatrix} \psi_D^2 \\ \psi_D^3 \end{pmatrix} \end{pmatrix} = \mathbf{0} \\
& = -\left(\nabla^2 - \frac{\partial^2}{\partial t^2} - m^2\right) \Psi_D . \\
& \left(\begin{array}{cc} \left(\begin{array}{cc} (-i\frac{\partial}{\partial t} + m) & 0 \\ 0 & (-i\frac{\partial}{\partial t} + m) \end{array} \right) & -i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} \\ i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \left(\begin{array}{cc} (i\frac{\partial}{\partial t} + m) & 0 \\ 0 & (i\frac{\partial}{\partial t} + m) \end{array} \right) \end{array} \right) \begin{pmatrix} \begin{pmatrix} \psi_D^4 \\ \psi_D^5 \end{pmatrix} \\ \begin{pmatrix} \psi_D^6 \\ \psi_D^7 \end{pmatrix} \end{pmatrix} = \mathbf{0} \\
& = -\left(\nabla^2 - \frac{\partial^2}{\partial t^2} - m^2\right) \begin{pmatrix} \Psi_D^C \\ \Psi_D^D \end{pmatrix} .
\end{aligned}$$

For symmetry purposes, let:

$$t \equiv ix^0$$

then, combining into a single matrix equation,

Let:

$$\begin{aligned}
\boldsymbol{\theta} &= \begin{pmatrix} \begin{pmatrix} \theta_D^0 \\ \theta_D^1 \end{pmatrix} \\ \begin{pmatrix} \theta_D^2 \\ \theta_D^3 \end{pmatrix} \\ \begin{pmatrix} \theta_D^4 \\ \theta_D^5 \end{pmatrix} \\ \begin{pmatrix} \theta_D^6 \\ \theta_D^7 \end{pmatrix} \end{pmatrix} \\
&= \left(\begin{array}{cccc} \left(\begin{array}{cc} (\partial_0 + m) & 0 \\ 0 & (\partial_0 + m) \end{array} \right) & i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \\ -i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \left(\begin{array}{cc} (-\partial_0 + m) & 0 \\ 0 & (-\partial_0 + m) \end{array} \right) & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \left(\begin{array}{cc} (-\partial_0 + m) & 0 \\ 0 & (-\partial_0 + m) \end{array} \right) & -i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \left(\begin{array}{cc} (\partial_0 + m) & 0 \\ 0 & (\partial_0 + m) \end{array} \right) \end{array} \right) \begin{pmatrix} \begin{pmatrix} \psi_D^0 \\ \psi_D^1 \end{pmatrix} \\ \begin{pmatrix} \psi_D^2 \\ \psi_D^3 \end{pmatrix} \\ \begin{pmatrix} \psi_D^4 \\ \psi_D^5 \end{pmatrix} \\ \begin{pmatrix} \psi_D^6 \\ \psi_D^7 \end{pmatrix} \end{pmatrix} \\
& \left(\begin{array}{cccc} \left(\begin{array}{cc} (-\partial_0 + m) & 0 \\ 0 & (-\partial_0 + m) \end{array} \right) & -i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \\ i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \left(\begin{array}{cc} (\partial_0 + m) & 0 \\ 0 & (\partial_0 + m) \end{array} \right) & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \left(\begin{array}{cc} (\partial_0 + m) & 0 \\ 0 & (\partial_0 + m) \end{array} \right) & i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & -i \begin{pmatrix} \partial_3 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & -\partial_3 \end{pmatrix} & \left(\begin{array}{cc} (-\partial_0 + m) & 0 \\ 0 & (-\partial_0 + m) \end{array} \right) \end{array} \right) \begin{pmatrix} \begin{pmatrix} \theta_D^0 \\ \theta_D^1 \end{pmatrix} \\ \begin{pmatrix} \theta_D^2 \\ \theta_D^3 \end{pmatrix} \\ \begin{pmatrix} \theta_D^4 \\ \theta_D^5 \end{pmatrix} \\ \begin{pmatrix} \theta_D^6 \\ \theta_D^7 \end{pmatrix} \end{pmatrix} = \\
& = -(\square - m^2) \boldsymbol{\theta} = \mathbf{0} .
\end{aligned}$$

Now,

Just as there are a number of representations of the Dirac equation, there is more than one matrix operator factorization of the Maxwell-Cassano equations [3].

One such, is:

$$\begin{pmatrix} \mathbf{H}^1 \\ \mathbf{H}^2 \\ \mathbf{H}^3 \\ \mathbf{H}^4 \\ \mathbf{H}^5 \\ \mathbf{H}^6 \\ \mathbf{H}^7 \\ \mathbf{H}^8 \end{pmatrix} = \begin{pmatrix} (\partial_1 - m_1) & (\partial_2 - m_2) & (\partial_3 - m_3) & (\partial_4 - m_4) & 0 & 0 & 0 & 0 \\ -(\partial_2 + m_2) & (\partial_1 + m_1) & 0 & 0 & 0 & 0 & (\partial_4 - m_4) & -(\partial_3 - m_3) \\ (\partial_3 + m_3) & 0 & -(\partial_1 + m_1) & 0 & 0 & (\partial_4 - m_4) & 0 & -(\partial_2 - m_2) \\ (\partial_4 + m_4) & 0 & 0 & -(\partial_1 + m_1) & 0 & -(\partial_3 - m_3) & (\partial_2 - m_2) & 0 \\ 0 & 0 & 0 & 0 & (\partial_1 + m_1) & (\partial_2 + m_2) & (\partial_3 + m_3) & (\partial_4 + m_4) \\ 0 & 0 & (\partial_4 + m_4) & -(\partial_3 + m_3) & -(\partial_2 - m_2) & (\partial_1 - m_1) & 0 & 0 \\ 0 & (\partial_4 + m_4) & 0 & -(\partial_2 + m_2) & (\partial_3 - m_3) & 0 & -(\partial_1 - m_1) & 0 \\ 0 & -(\partial_3 + m_3) & (\partial_2 + m_2) & 0 & (\partial_4 - m_4) & 0 & 0 & -(\partial_1 - m_1) \end{pmatrix} \begin{pmatrix} h^1 \\ h^2 \\ h^3 \\ h^4 \\ h^5 \\ h^6 \\ h^7 \\ h^8 \end{pmatrix}$$

Then:

$$\begin{pmatrix} (\partial_1 + m_1) & -(\partial_2 - m_2) & (\partial_3 - m_3) & (\partial_4 - m_4) & 0 & 0 & 0 & 0 \\ (\partial_2 + m_2) & (\partial_1 - m_1) & 0 & 0 & 0 & 0 & (\partial_4 - m_4) & -(\partial_3 - m_3) \\ (\partial_3 + m_3) & 0 & -(\partial_1 - m_1) & 0 & 0 & (\partial_4 - m_4) & 0 & (\partial_2 - m_2) \\ (\partial_4 + m_4) & 0 & 0 & -(\partial_1 - m_1) & 0 & -(\partial_3 - m_3) & -(\partial_2 - m_2) & 0 \\ 0 & 0 & 0 & 0 & (\partial_1 - m_1) & -(\partial_2 + m_2) & (\partial_3 + m_3) & (\partial_4 + m_4) \\ 0 & 0 & (\partial_4 + m_4) & -(\partial_3 + m_3) & (\partial_2 - m_2) & (\partial_1 + m_1) & 0 & 0 \\ 0 & (\partial_4 + m_4) & 0 & (\partial_2 + m_2) & (\partial_3 - m_3) & 0 & -(\partial_1 + m_1) & 0 \\ 0 & -(\partial_3 + m_3) & -(\partial_2 + m_2) & 0 & (\partial_4 - m_4) & 0 & 0 & -(\partial_1 + m_1) \end{pmatrix} \begin{pmatrix} \mathbf{H}^1 \\ \mathbf{H}^2 \\ \mathbf{H}^3 \\ \mathbf{H}^4 \\ \mathbf{H}^5 \\ \mathbf{H}^6 \\ \mathbf{H}^7 \\ \mathbf{H}^8 \end{pmatrix} = \begin{pmatrix} (\square - |m|^2) & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & (\square - |m|^2) & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & (\square - |m|^2) & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & (\square - |m|^2) & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & (\square - |m|^2) & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & (\square - |m|^2) & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & (\square - |m|^2) & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & (\square - |m|^2) \end{pmatrix} \begin{pmatrix} h^1 \\ h^2 \\ h^3 \\ h^4 \\ h^5 \\ h^6 \\ h^7 \\ h^8 \end{pmatrix}$$

Another matrix operator factorization of the Maxwell-Cassano equations may be compactly written, as follows [1][3]:

From the definitions:

$$\mathbf{f} \equiv \mathbf{w}^{4,i} f^i, \text{ where } : f^i \equiv \begin{pmatrix} f_+^i \\ f_-^i \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} f_+^i \\ f_-^i \end{pmatrix}, f_+^i, f_-^i \in \mathbf{R}$$

$$D_i^+ \equiv (\partial_i + m_i), \quad D_i^- \equiv (\partial_i - m_i)$$

$$D_i \equiv \begin{pmatrix} D_i^+ & 0 \\ 0 & D_i^- \end{pmatrix}, \quad D_i^\updownarrow \equiv \begin{pmatrix} D_i^- & 0 \\ 0 & D_i^+ \end{pmatrix},$$

$$D_i^{\leftrightarrow} \equiv \begin{pmatrix} 0 & D_i^- \\ D_i^+ & 0 \end{pmatrix}, \quad D_i^{\leftrightarrow\updownarrow} \equiv \begin{pmatrix} 0 & D_i^+ \\ D_i^- & 0 \end{pmatrix}$$

$$\hat{\mathbf{f}} \equiv \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}$$

such a factorization may be

$$\begin{pmatrix} D_0 & D_3^{\leftrightarrow} & -D_2^{\leftrightarrow} & D_1 \\ -D_3^{\leftrightarrow} & D_0 & D_1^{\leftrightarrow} & D_2 \\ D_2^{\leftrightarrow} & -D_1^{\leftrightarrow} & D_0 & D_3 \\ D_1^\updownarrow & D_2^\updownarrow & D_3^\updownarrow & -D_0^\updownarrow \end{pmatrix} \begin{pmatrix} D_0^\updownarrow & -D_3^{\leftrightarrow} & D_2^{\leftrightarrow} & D_1 \\ D_3^{\leftrightarrow} & D_0^\updownarrow & -D_1^{\leftrightarrow} & D_2 \\ -D_2^{\leftrightarrow} & D_1^{\leftrightarrow} & D_0^\updownarrow & D_3 \\ D_1^\updownarrow & D_2^\updownarrow & D_3^\updownarrow & -D_0^\updownarrow \end{pmatrix} \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}$$

$$= \begin{pmatrix} -D_0 & D_3^{\leftrightarrow} & -D_2^{\leftrightarrow} & -D_1 \\ -D_3^{\leftrightarrow} & -D_0 & D_1^{\leftrightarrow} & -D_2 \\ D_2^{\leftrightarrow} & -D_1^{\leftrightarrow} & -D_0 & -D_3 \\ -D_1^\updownarrow & -D_2^\updownarrow & -D_3^\updownarrow & D_0^\updownarrow \end{pmatrix} \begin{pmatrix} -D_0^\updownarrow & -D_3^{\leftrightarrow} & D_2^{\leftrightarrow} & -D_1 \\ D_3^{\leftrightarrow} & -D_0^\updownarrow & -D_1^{\leftrightarrow} & -D_2 \\ -D_2^{\leftrightarrow} & D_1^{\leftrightarrow} & -D_0^\updownarrow & -D_3 \\ -D_1^\updownarrow & -D_2^\updownarrow & -D_3^\updownarrow & D_0^\updownarrow \end{pmatrix} \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}$$

$$= \begin{pmatrix} D_0 & -D_3^{\leftrightarrow} & D_2^{\leftrightarrow} & D_1 \\ D_3^{\leftrightarrow} & D_0 & -D_1^{\leftrightarrow} & D_2 \\ -D_2^{\leftrightarrow} & D_1^{\leftrightarrow} & D_0 & D_3 \\ D_1^\updownarrow & D_2^\updownarrow & D_3^\updownarrow & -D_0^\updownarrow \end{pmatrix} \begin{pmatrix} D_0^\updownarrow & D_3^{\leftrightarrow} & -D_2^{\leftrightarrow} & D_1 \\ -D_3^{\leftrightarrow} & D_0^\updownarrow & D_1^{\leftrightarrow} & D_2 \\ D_2^{\leftrightarrow} & -D_1^{\leftrightarrow} & D_0^\updownarrow & D_3 \\ D_1^\updownarrow & D_2^\updownarrow & D_3^\updownarrow & -D_0^\updownarrow \end{pmatrix} \begin{pmatrix} f^1 \\ f^2 \\ f^3 \\ f^0 \end{pmatrix}$$

$$= \begin{pmatrix} (\square - |m|^2)f^1 \\ (\square - |m|^2)f^2 \\ (\square - |m|^2)f^3 \\ (\square - |m|^2)f^0 \end{pmatrix} = (\square - |m|^2)\hat{\mathbf{f}}$$

For the stationary state the source/sink density term vanishes in the Maxwell-Cassano equations, which allows an equating of the Maxwell-Cassano equation & Dirac equation factorizations.

These may imply correlations between the Dirac equation and the Maxwell-Cassano equations as the correspondences/mappings:

$$m \Leftrightarrow |m| \ \& \ -\theta_D^i \Leftrightarrow f^h.$$

The Dirac equation may be expanded with the above notation into:

$$\begin{aligned} (-\partial_0 + m)\theta_D^0 - i\partial_3\theta_D^2 - i(\partial_1 - i\partial_2)\theta_D^3 &= 0 \\ (-\partial_0 + m)\theta_D^1 - i(\partial_1 + i\partial_2)\theta_D^2 + i\partial_3\theta_D^3 &= 0 \\ i\partial_3\theta_D^0 + i(\partial_1 - i\partial_2)\theta_D^1 + (\partial_0 + m)\theta_D^2 &= 0 \\ i(\partial_1 + i\partial_2)\theta_D^0 - i\partial_3\theta_D^1 + (\partial_0 + m)\theta_D^3 &= 0 \\ (\partial_0 + m)\theta_D^4 + i\partial_3\theta_D^6 + i(\partial_1 - i\partial_2)\theta_D^7 &= 0 \\ (\partial_0 + m)\theta_D^5 + i(\partial_1 + i\partial_2)\theta_D^6 - i\partial_3\theta_D^7 &= 0 \\ -i\partial_3\theta_D^4 - i(\partial_1 - i\partial_2)\theta_D^5 + (-\partial_0 + m)\theta_D^6 &= 0 \\ -i(\partial_1 + i\partial_2)\theta_D^4 + i\partial_3\theta_D^5 + (-\partial_0 + m)\theta_D^7 &= 0 \end{aligned}$$

or:

$$\begin{aligned} (-\partial_0 + m)\theta_D^0 & & -i\partial_3\theta_D^2 - i(\partial_1 - i\partial_2)\theta_D^3 &= 0 \\ (\partial_0 + m)\theta_D^4 & & +i\partial_3\theta_D^6 + i(\partial_1 - i\partial_2)\theta_D^7 &= 0 \\ (-\partial_0 + m)\theta_D^1 - i(\partial_1 + i\partial_2)\theta_D^2 & & +i\partial_3\theta_D^3 &= 0 \\ (\partial_0 + m)\theta_D^5 + i(\partial_1 + i\partial_2)\theta_D^6 & & -i\partial_3\theta_D^7 &= 0 \\ i\partial_3\theta_D^0 + i(\partial_1 - i\partial_2)\theta_D^1 & & +(\partial_0 + m)\theta_D^2 &= 0 \\ -i\partial_3\theta_D^4 - i(\partial_1 - i\partial_2)\theta_D^5 & & +(-\partial_0 + m)\theta_D^6 &= 0 \\ i(\partial_1 + i\partial_2)\theta_D^0 & & -i\partial_3\theta_D^1 & +(\partial_0 + m)\theta_D^3 = 0 \\ -i(\partial_1 + i\partial_2)\theta_D^4 & & +i\partial_3\theta_D^5 & +(-\partial_0 + m)\theta_D^7 = 0 \end{aligned}$$

As [1] shows, the component pairs may be organized such that this organization exhibits the mass-generalization of Maxwell's equations, but organizing them while comparing them analogously to the Dirac equations yields:

$$\begin{aligned} (-\partial_0 + m)\theta_D^0 & & -i\partial_3\theta_D^2 & & -i(\partial_1 - i\partial_2)\theta_D^3 &= 0 \\ (\partial_0 + m)\theta_D^4 & & +i\partial_3\theta_D^6 & & +i(\partial_1 - i\partial_2)\theta_D^7 &= 0 \\ (\partial_0 + m_0)Z^0 & +(\partial_3 - m_3)Z^5 & -(\partial_2 - m_2)Z^6 & +(\partial_1 + m_1)Z^3 &= J^0 \\ (\partial_0 - m_0)Z^4 & +(\partial_3 + m_3)Z^1 & -(\partial_2 + m_2)Z^2 & +(\partial_1 - m_1)Z^7 &= J^4 \\ (-\partial_0 + m)\theta_D^1 & -i(\partial_1 + i\partial_2)\theta_D^2 & & & +i\partial_3\theta_D^3 &= 0 \\ (\partial_0 + m)\theta_D^5 & +i(\partial_1 + i\partial_2)\theta_D^6 & & & -i\partial_3\theta_D^7 &= 0 \\ (\partial_0 + m_0)Z^1 & +(\partial_1 - m_1)Z^6 & +(\partial_2 + m_2)Z^3 & -(\partial_3 - m_3)Z^4 &= J^1 \\ (\partial_0 - m_0)Z^5 & +(\partial_1 + m_1)Z^2 & +(\partial_2 - m_2)Z^7 & -(\partial_3 + m_3)Z^0 &= J^5 \\ i\partial_3\theta_D^0 & +i(\partial_1 - i\partial_2)\theta_D^1 & +(\partial_0 + m)\theta_D^2 & & &= 0 \\ -i\partial_3\theta_D^4 & -i(\partial_1 - i\partial_2)\theta_D^5 & +(-\partial_0 + m)\theta_D^6 & & &= 0 \\ (\partial_3 + m_3)Z^3 & -(\partial_1 - m_1)Z^5 & +(\partial_0 + m_0)Z^2 & +(\partial_2 - m_2)Z^4 &= J^2 \\ (\partial_3 - m_3)Z^7 & -(\partial_1 + m_1)Z^1 & +(\partial_0 - m_0)Z^6 & +(\partial_2 + m_2)Z^0 &= J^6 \\ i(\partial_1 + i\partial_2)\theta_D^0 & & -i\partial_3\theta_D^1 & & +(\partial_0 + m)\theta_D^3 &= 0 \\ -i(\partial_1 + i\partial_2)\theta_D^4 & & +i\partial_3\theta_D^5 & & +(-\partial_0 + m)\theta_D^7 &= 0 \\ (\partial_1 - m_1)Z^0 & +(\partial_3 - m_3)Z^2 & +(\partial_2 - m_2)Z^1 & -(\partial_0 - m_0)Z^3 &= J^3 \\ (\partial_1 + m_1)Z^4 & +(\partial_3 + m_3)Z^6 & +(\partial_2 + m_2)Z^5 & -(\partial_0 + m_0)Z^7 &= J^7 \end{aligned}$$

So:

$$\begin{aligned} (\partial_0 - m)\theta_D^0 & +i(\partial_1 - i\partial_2)\theta_D^3 & & +i\partial_3\theta_D^2 & &= 0 \\ (\partial_0 + m)\theta_D^4 & +i(\partial_1 - i\partial_2)\theta_D^7 & & +i\partial_3\theta_D^6 & &= 0 \\ (\partial_0 + m_0)Z^0 & +(\partial_1 + m_1)Z^3 & -(\partial_2 - m_2)Z^6 & +(\partial_3 - m_3)Z^5 &= J^0 \\ (\partial_0 - m_0)Z^4 & +(\partial_1 - m_1)Z^7 & -(\partial_2 + m_2)Z^2 & +(\partial_3 + m_3)Z^1 &= J^4 \\ (\partial_0 - m)\theta_D^1 & +i(\partial_1 + i\partial_2)\theta_D^2 & & & -i\partial_3\theta_D^3 &= 0 \\ (\partial_0 + m)\theta_D^5 & +i(\partial_1 + i\partial_2)\theta_D^6 & & & -i\partial_3\theta_D^7 &= 0 \\ (\partial_0 - m_0)Z^5 & +(\partial_1 + m_1)Z^2 & +(\partial_2 - m_2)Z^7 & -(\partial_3 + m_3)Z^0 &= J^5 \\ (\partial_0 + m_0)Z^1 & +(\partial_1 - m_1)Z^6 & +(\partial_2 + m_2)Z^3 & -(\partial_3 - m_3)Z^4 &= J^1 \\ (\partial_0 + m)\theta_D^2 & +i(\partial_1 - i\partial_2)\theta_D^1 & & +i\partial_3\theta_D^0 & &= 0 \\ (\partial_0 - m)\theta_D^6 & +i(\partial_1 - i\partial_2)\theta_D^5 & & +i\partial_3\theta_D^4 & &= 0 \\ (\partial_0 + m_0)Z^2 & -(\partial_1 - m_1)Z^5 & +(\partial_2 - m_2)Z^4 & +(\partial_3 + m_3)Z^3 &= J^2 \\ (\partial_0 - m_0)Z^6 & -(\partial_1 + m_1)Z^1 & +(\partial_2 + m_2)Z^0 & +(\partial_3 - m_3)Z^7 &= J^6 \\ (\partial_0 + m)\theta_D^3 & +i(\partial_1 + i\partial_2)\theta_D^0 & & & -i\partial_3\theta_D^1 &= 0 \\ (\partial_0 - m)\theta_D^7 & +i(\partial_1 + i\partial_2)\theta_D^4 & & & -i\partial_3\theta_D^5 &= 0 \\ (\partial_0 + m_0)Z^7 & -(\partial_1 + m_1)Z^4 & -(\partial_2 + m_2)Z^5 & & -(\partial_3 + m_3)Z^6 &= J^7 \\ (\partial_0 - m_0)Z^3 & -(\partial_1 - m_1)Z^0 & -(\partial_2 - m_2)Z^1 & & -(\partial_3 - m_3)Z^2 &= J^3 \end{aligned}$$

Continuing the comparison with the Maxwell-Cassano equations in the special case:

$$m_0 \rightarrow -m, m_1 = m_2 = m_3 = 0 :$$

$$(\partial_0 - m)\theta_D^0 + i(\partial_1 - i\partial_2)\theta_D^3 + i\partial_3\theta_D^2 = 0$$

$$\begin{aligned}
(\partial_0 + m)\theta_D^4 + i(\partial_1 - i\partial_2)\theta_D^7 + i\partial_3\theta_D^6 &= 0 \\
(\partial_0 - m)Z^0 + \partial_1Z^3 - \partial_2Z^6 + \partial_3Z^5 &= J^0 \\
(\partial_0 + m)Z^4 + \partial_1Z^7 - \partial_2Z^2 + \partial_3Z^1 &= J^4 \\
(\partial_0 - m)\theta_D^1 + i(\partial_1 + i\partial_2)\theta_D^2 - i\partial_3\theta_D^3 &= 0 \\
(\partial_0 + m)\theta_D^5 + i(\partial_1 + i\partial_2)\theta_D^6 - i\partial_3\theta_D^7 &= 0 \\
(\partial_0 - m)Z^1 + \partial_1Z^6 + \partial_2Z^3 - \partial_3Z^4 &= J^1 \\
(\partial_0 + m)Z^5 + \partial_1Z^2 + \partial_2Z^7 - \partial_3Z^0 &= J^5 \\
(\partial_0 + m)\theta_D^2 + i(\partial_1 - i\partial_2)\theta_D^1 + i\partial_3\theta_D^0 &= 0 \\
(\partial_0 - m)\theta_D^6 + i(\partial_1 - i\partial_2)\theta_D^5 + i\partial_3\theta_D^4 &= 0 \\
(\partial_0 - m)Z^2 - \partial_1Z^5 + \partial_2Z^4 + \partial_3Z^3 &= J^2 \\
(\partial_0 + m)Z^6 - \partial_1Z^1 + \partial_2Z^0 + \partial_3Z^7 &= J^6 \\
(\partial_0 + m)\theta_D^3 + i(\partial_1 + i\partial_2)\theta_D^0 - i\partial_3\theta_D^1 &= 0 \\
(\partial_0 - m)\theta_D^7 + i(\partial_1 + i\partial_2)\theta_D^4 - i\partial_3\theta_D^5 &= 0 \\
(\partial_0 - m)Z^7 - \partial_1Z^4 - \partial_2Z^5 - \partial_3Z^6 &= J^7 \\
(\partial_0 + m)Z^3 - \partial_1Z^0 - \partial_2Z^1 - \partial_3Z^2 &= J^3
\end{aligned}$$

or:

$$\begin{aligned}
(\partial_0 - m)\theta_D^0 + i(\partial_1 - i\partial_2)\theta_D^3 + i\partial_3\theta_D^2 &= 0 \\
(\partial_0 + m)\theta_D^4 + i(\partial_1 - i\partial_2)\theta_D^7 + i\partial_3\theta_D^6 &= 0 \\
(\partial_0 - m)\theta_D^1 + i(\partial_1 + i\partial_2)\theta_D^2 - i\partial_3\theta_D^3 &= 0 \\
(\partial_0 + m)\theta_D^5 + i(\partial_1 + i\partial_2)\theta_D^6 - i\partial_3\theta_D^7 &= 0 \\
(\partial_0 + m)\theta_D^2 + i(\partial_1 - i\partial_2)\theta_D^1 + i\partial_3\theta_D^0 &= 0 \\
(\partial_0 - m)\theta_D^6 + i(\partial_1 - i\partial_2)\theta_D^5 + i\partial_3\theta_D^4 &= 0 \\
(\partial_0 + m)\theta_D^3 + i(\partial_1 + i\partial_2)\theta_D^0 - i\partial_3\theta_D^1 &= 0 \\
(\partial_0 - m)\theta_D^7 + i(\partial_1 + i\partial_2)\theta_D^4 - i\partial_3\theta_D^5 &= 0 \\
(\partial_0 - m)Z^0 + \partial_1Z^3 - \partial_2Z^6 + \partial_3Z^5 &= J^0 \\
(\partial_0 + m)Z^4 + \partial_1Z^7 - \partial_2Z^2 + \partial_3Z^1 &= J^4 \\
(\partial_0 - m)Z^1 + \partial_1Z^6 + \partial_2Z^3 - \partial_3Z^4 &= J^1 \\
(\partial_0 + m)Z^5 + \partial_1Z^2 + \partial_2Z^7 - \partial_3Z^0 &= J^5 \\
(\partial_0 + m)Z^6 - \partial_1Z^1 + \partial_2Z^0 + \partial_3Z^7 &= J^6 \\
(\partial_0 - m)Z^2 - \partial_1Z^5 + \partial_2Z^4 + \partial_3Z^3 &= J^2 \\
(\partial_0 + m)Z^3 - \partial_1Z^0 - \partial_2Z^1 - \partial_3Z^2 &= J^3 \\
(\partial_0 - m)Z^7 - \partial_1Z^4 - \partial_2Z^5 - \partial_3Z^6 &= J^7
\end{aligned}$$

So, extending the Dirac equation beyond the source/sink free case (so looking beyond just eigenvalues and eigenvectors); and writing in matrix form, and comparing:

$$\begin{pmatrix}
(\partial_0 - m) & 0 & 0 & 0 & i\partial_3 & 0 & i(\partial_1 - i\partial_2) & 0 \\
0 & (\partial_0 + m) & 0 & 0 & 0 & i\partial_3 & 0 & i(\partial_1 - i\partial_2) \\
0 & 0 & (\partial_0 - m) & 0 & i(\partial_1 + i\partial_2) & 0 & -i\partial_3 & 0 \\
0 & 0 & 0 & (\partial_0 + m) & 0 & i(\partial_1 + i\partial_2) & 0 & -i\partial_3 \\
i\partial_3 & 0 & i(\partial_1 - i\partial_2) & 0 & (\partial_0 + m) & 0 & 0 & 0 \\
0 & i\partial_3 & 0 & i(\partial_1 - i\partial_2) & 0 & (\partial_0 - m) & 0 & 0 \\
i(\partial_1 + i\partial_2) & 0 & -i\partial_3 & 0 & 0 & 0 & (\partial_0 + m) & 0 \\
0 & i(\partial_1 + i\partial_2) & 0 & -i\partial_3 & 0 & 0 & 0 & (\partial_0 - m)
\end{pmatrix}
\begin{pmatrix}
\theta_D^0 \\
\theta_D^4 \\
\theta_D^1 \\
\theta_D^5 \\
\theta_D^2 \\
\theta_D^6 \\
\theta_D^3 \\
\theta_D^7
\end{pmatrix}
=
\begin{pmatrix}
\Phi^0 \\
\Phi^4 \\
\Phi^1 \\
\Phi^5 \\
\Phi^2 \\
\Phi^6 \\
\Phi^3 \\
\Phi^7
\end{pmatrix}$$

$$\begin{aligned}
(\partial_0 - m)Z^2 - \partial_1Z^5 + \partial_2Z^4 + \partial_3Z^3 &= J^2 \\
(\partial_0 + m)Z^6 - \partial_1Z^1 + \partial_2Z^0 + \partial_3Z^7 &= J^6 \\
(\partial_0 + m)Z^3 - \partial_1Z^0 - \partial_2Z^1 - \partial_3Z^2 &= J^3 \\
(\partial_0 - m)Z^7 - \partial_1Z^4 - \partial_2Z^5 - \partial_3Z^6 &= J^7
\end{aligned}$$

$$\begin{pmatrix}
(\partial_0 - m) & 0 & 0 & \partial_3 & 0 & -\partial_2 & \partial_1 & 0 \\
0 & (\partial_0 + m) & \partial_3 & 0 & -\partial_2 & 0 & 0 & \partial_1 \\
0 & -\partial_3 & (\partial_0 - m) & 0 & 0 & \partial_1 & \partial_2 & 0 \\
-\partial_3 & 0 & 0 & (\partial_0 + m) & \partial_1 & 0 & 0 & \partial_2 \\
0 & \partial_2 & 0 & -\partial_1 & (\partial_0 - m) & 0 & \partial_3 & 0 \\
\partial_2 & 0 & -\partial_1 & 0 & 0 & (\partial_0 + m) & 0 & \partial_3 \\
-\partial_1 & 0 & -\partial_2 & 0 & -\partial_3 & 0 & (\partial_0 + m) & 0 \\
0 & -\partial_1 & 0 & -\partial_2 & 0 & -\partial_3 & 0 & (\partial_0 - m)
\end{pmatrix}
\begin{pmatrix}
Z^0 \\
Z^4 \\
Z^1 \\
Z^5 \\
Z^2 \\
Z^6 \\
Z^3 \\
Z^7
\end{pmatrix}
=
\begin{pmatrix}
J^0 \\
J^4 \\
J^1 \\
J^5 \\
J^2 \\
J^6 \\
J^3 \\
J^7
\end{pmatrix}$$

The matrix is equivalent (have the same solution set) under the elementary row operation of interchanging rows, so interchanging rows 4 & 5:

$$\begin{pmatrix} (\partial_0 - m) & 0 & 0 & \partial_3 & 0 & -\partial_2 & \partial_1 & 0 \\ 0 & (\partial_0 + m) & \partial_3 & 0 & -\partial_2 & 0 & 0 & \partial_1 \\ 0 & -\partial_3 & (\partial_0 - m) & 0 & 0 & \partial_1 & \partial_2 & 0 \\ -\partial_3 & 0 & 0 & (\partial_0 + m) & \partial_1 & 0 & 0 & \partial_2 \\ \partial_2 & 0 & -\partial_1 & 0 & 0 & (\partial_0 + m) & 0 & \partial_3 \\ 0 & \partial_2 & 0 & -\partial_1 & (\partial_0 - m) & 0 & \partial_3 & 0 \\ -\partial_1 & 0 & -\partial_2 & 0 & -\partial_3 & 0 & (\partial_0 + m) & 0 \\ 0 & -\partial_1 & 0 & -\partial_2 & 0 & -\partial_3 & 0 & (\partial_0 - m) \end{pmatrix} \begin{pmatrix} Z^0 \\ Z^4 \\ Z^1 \\ Z^5 \\ Z^2 \\ Z^6 \\ Z^3 \\ Z^7 \end{pmatrix} = \begin{pmatrix} J^0 \\ J^4 \\ J^1 \\ J^5 \\ J^2 \\ J^6 \\ J^3 \\ J^7 \end{pmatrix}$$

To retain equality under the elementary row operation of interchanging columns 4 & 5, the rows of vectors \mathbf{Z} & \mathbf{J} are interchanged:

$$\begin{pmatrix} (\partial_0 - m) & 0 & 0 & \partial_3 & -\partial_2 & 0 & \partial_1 & 0 \\ 0 & (\partial_0 + m) & \partial_3 & 0 & 0 & -\partial_2 & 0 & \partial_1 \\ 0 & -\partial_3 & (\partial_0 - m) & 0 & \partial_1 & 0 & \partial_2 & 0 \\ -\partial_3 & 0 & 0 & (\partial_0 + m) & 0 & \partial_1 & 0 & \partial_2 \\ \partial_2 & 0 & -\partial_1 & 0 & (\partial_0 + m) & 0 & 0 & \partial_3 \\ 0 & \partial_2 & 0 & -\partial_1 & 0 & (\partial_0 - m) & \partial_3 & 0 \\ -\partial_1 & 0 & -\partial_2 & 0 & 0 & -\partial_3 & (\partial_0 + m) & 0 \\ 0 & -\partial_1 & 0 & -\partial_2 & -\partial_3 & 0 & 0 & (\partial_0 - m) \end{pmatrix} \begin{pmatrix} Z^0 \\ Z^4 \\ Z^1 \\ Z^5 \\ Z^6 \\ Z^2 \\ Z^3 \\ Z^7 \end{pmatrix} = \begin{pmatrix} J^0 \\ J^4 \\ J^1 \\ J^5 \\ J^6 \\ J^2 \\ J^3 \\ J^7 \end{pmatrix}$$

Now, consider each matrix as a sum:

$$\begin{pmatrix} (\partial_0 - m) & 0 & 0 & 0 & i\partial_3 & 0 & i(\partial_1 - i\partial_2) & 0 \\ 0 & (\partial_0 + m) & 0 & 0 & 0 & i\partial_3 & 0 & i(\partial_1 - i\partial_2) \\ 0 & 0 & (\partial_0 - m) & 0 & i(\partial_1 + i\partial_2) & 0 & -i\partial_3 & 0 \\ 0 & 0 & 0 & (\partial_0 + m) & 0 & i(\partial_1 + i\partial_2) & 0 & -i\partial_3 \\ i\partial_3 & 0 & i(\partial_1 - i\partial_2) & 0 & (\partial_0 + m) & 0 & 0 & 0 \\ 0 & i\partial_3 & 0 & i(\partial_1 - i\partial_2) & 0 & (\partial_0 - m) & 0 & 0 \\ i(\partial_1 + i\partial_2) & 0 & -i\partial_3 & 0 & 0 & 0 & (\partial_0 + m) & 0 \\ 0 & i(\partial_1 + i\partial_2) & 0 & -i\partial_3 & 0 & 0 & 0 & (\partial_0 - m) \end{pmatrix} = \begin{pmatrix} (\partial_0 - m) & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & (\partial_0 + m) & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & (\partial_0 - m) & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & (\partial_0 + m) & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & (\partial_0 + m) & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & (\partial_0 - m) & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & (\partial_0 + m) & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & (\partial_0 - m) \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 0 & i\partial_3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & i\partial_3 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -i\partial_3 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -i\partial_3 \\ i\partial_3 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & i\partial_3 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -i\partial_3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i\partial_3 & 0 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & i\partial_1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & i\partial_1 \\ 0 & 0 & 0 & 0 & i\partial_1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & i\partial_1 & 0 & 0 \\ 0 & 0 & i\partial_1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & i\partial_1 & 0 & 0 & 0 & 0 \\ i\partial_1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & i\partial_1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & \partial_2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \partial_2 \\ 0 & 0 & 0 & 0 & -\partial_2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -\partial_2 & 0 & 0 \\ 0 & 0 & \partial_2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \partial_2 & 0 & 0 & 0 & 0 \\ -\partial_2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -\partial_2 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

and:

$$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \stackrel{\bar{\partial}_2}{\Leftrightarrow} \begin{pmatrix} 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \end{pmatrix} \stackrel{\partial_3}{}$$

By the invertible matrix theorem each matrix is invertible. Thus, these transformations are onto, bijective.

From the first matrices on each side of the sum, the rest of the transformations are even more easily seen. The full set of transformations follow.

$(\partial_0 \pm m) \Leftrightarrow (\bar{\partial}_0 \pm m)$	
$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \stackrel{\partial_1}{\Leftrightarrow}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \stackrel{i\bar{\partial}_1}{}$
$\begin{pmatrix} 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \end{pmatrix} \stackrel{\partial_2}{\Leftrightarrow}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \end{pmatrix} \stackrel{i\bar{\partial}_3}{}$
$\begin{pmatrix} 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \end{pmatrix} \stackrel{\partial_3}{\Leftrightarrow}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \stackrel{\bar{\partial}_2}{}$

[Dirac (barred)]

Let:

$$\sigma_2^0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \sigma_2^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \mathbf{0}_2 = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

then, more compactly:

$$\begin{array}{c}
(\partial_0 \pm m) \Leftrightarrow (\bar{\partial}_0 \pm m) \quad \Rightarrow x^0 = \bar{x}^0 \\
\left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 & 0 \\ \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \\ -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) \partial_1 \Leftrightarrow \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 & 0 \\ \mathbf{0}_2 & \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \\ \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) i\bar{\partial}_1 \Rightarrow x^1 = -i \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 & 0 \\ \mathbf{0}_2 & \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \\ \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right)^{-1} \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 & 0 \\ \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \\ -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) \bar{x}^1 \\
\left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) \partial_2 \Leftrightarrow \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^0 \\ \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) i\bar{\partial}_3 \Rightarrow x^2 = -i \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^0 \\ \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right)^{-1} \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) x^3 \\
\left(\begin{array}{cccc} \mathbf{0}_2 & \sigma_2^1 & \mathbf{0}_2 & \mathbf{0}_2 \\ -\sigma_2^1 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^1 \\ \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^1 & \mathbf{0}_2 \end{array} \right) \partial_3 \Leftrightarrow \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \\ -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right) \bar{\partial}_2 \Rightarrow x^2 = \left(\begin{array}{cccc} \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^0 \\ \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 \\ -\sigma_2^0 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \end{array} \right)^{-1} \left(\begin{array}{cccc} \mathbf{0}_2 & \sigma_2^1 & \mathbf{0}_2 & \mathbf{0}_2 \\ -\sigma_2^1 & \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{0}_2 & \mathbf{0}_2 & \sigma_2^1 \\ \mathbf{0}_2 & \mathbf{0}_2 & -\sigma_2^1 & \mathbf{0}_2 \end{array} \right) \bar{x}^3 \\
\text{[Dirac (barred)]}
\end{array}$$

$$\left(\begin{array}{c} \theta_D^0 \\ \theta_D^4 \\ \theta_D^1 \\ \theta_D^5 \\ \theta_D^2 \\ \theta_D^6 \\ \theta_D^3 \\ \theta_D^7 \end{array} \right) = \left(\begin{array}{c} Z^0 \\ Z^4 \\ Z^1 \\ Z^5 \\ Z^6 \\ Z^2 \\ Z^3 \\ Z^7 \end{array} \right) \quad \left(\begin{array}{c} \Phi^0 \\ \Phi^4 \\ \Phi^1 \\ \Phi^5 \\ \Phi^2 \\ \Phi^6 \\ \Phi^3 \\ \Phi^7 \end{array} \right) = \left(\begin{array}{c} J^0 \\ J^4 \\ J^1 \\ J^5 \\ J^6 \\ J^2 \\ J^3 \\ J^7 \end{array} \right)$$

This proves that the mass-generalized Maxwell's equations (Maxwell-Cassano equations) is a more general analysis of fundamental-elementary particle phenomena.

It further proves that the Lagrangian is far simpler than that consisting of the Glashow-Salam-Weinberg + fermion + Higgs + Yukawa kludge.

Also, it explains the group structure and architecture of the fermions [2].

It also proves that those with wealth to seek the truth choose not to do so, but with all deceivableness and unrighteousness in them they have not the love of the truth, but rather embrace strong delusion, that they profess a lie.

References and further readings

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