

# Dimension of physical space

Gunn Quznetsov  
gunn@mail.ru, quznets@yahoo.com

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## Abstract

Each vector of state has its own corresponding element of the Cayley-Dickson algebra. Properties of a state vector require that this algebra was a normalized division algebra. By the Hurwitz and Frobenius theorems maximal dimension of such algebra is 8. Consequently, a dimension of corresponding complex state vectors is 4, and a dimension of the Clifford set elements is  $4 \times 4$ . Such set contains 5 matrices - among them - 3 diagonal. Hence, a dimension of the dot events space is equal to  $3+1$ .

## 1 Dimension of physical space

Further I use Cayley-Dickson algebras [1, 2]:

Let  $1, i, j, k, E, I, J, K$  be basis elements of a 8-dimensional algebra Cayley (*the octavians algebra*) [1, 2]. A product of this algebra is defined the following way [1]:

1. for every basic element e:

$$ee = -1;$$

2. If  $u_1, u_2, v_1, v_2$  are real number then

$$(u_1 + u_2i)(v_1 + v_2i) = (u_1v_1 - v_2u_2) + (v_2u_1 + u_2v_1)i.$$

3. If  $u_1, u_2, v_1, v_2$  are numbers of shape  $w = w_1 + w_2i$  ( $w_s$ , and  $s \in \{1, 2\}$  are real numbers, and  $\bar{w} = w_1 - w_2i$ ) then

$$(u_1 + u_2j)(v_1 + v_2j) = (u_1v_1 - \bar{v}_2u_2) + (v_2u_1 + u_2\bar{v}_1)j \quad (1)$$

and  $ij = k$ .

4. If  $u_1, u_2, v_1, v_2$  are number of shape  $w = w_1 + w_2i + w_3j + w_4k$  ( $w_s$ , and  $s \in \{1, 2, 3, 4\}$  are real numbers, and  $\bar{w} = w_1 - w_2i - w_3j - w_4k$ ) then

$$(u_1 + u_2E)(v_1 + v_2E) = (u_1v_1 - \bar{v}_2u_2) + (v_2u_1 + u_2\bar{v}_1)E \quad (2)$$

and

$$\begin{aligned} iE &= I, \\ jE &= J, \\ kE &= K. \end{aligned}$$

Therefore, in according with point 2.: the real numbers field ( $\mathbf{R}$ ) is extended to the complex numbers field ( $\mathbf{C}$ ), and in according with point 3.: the complex numbers field is expanded to the quaternions field ( $\mathbf{K}$ ), and point 4. expands the quaternions fields to the octavians field ( $\mathbf{O}$ ). This method of expanding of fields is called a *Dickson doubling procedure* [1].

If

$$u = a + bi + cj + dk + AE + BI + CJ + K$$

with real  $a, b, c, d, A, B, C, D$  then a real number

$$\|u\| \stackrel{def}{=} \sqrt{u\bar{u}} = (a^2 + b^2 + c^2 + d^2 + A^2 + B^2 + C^2 + D^2)^{0.5}$$

is called a *norm* of octavian  $u$  [1].

For each octavians  $u$  and  $v$ :

$$\|uv\| = \|u\| \|v\|. \quad (3)$$

Algebras with this conditions are called *normalized algebras* [1, 2].

Any 3+1-vector of a probability density can be represented by the following equations in matrix form [4], [5]

$$\begin{aligned} \rho &= \varphi^\dagger \varphi, \\ j_k &= \varphi^\dagger \beta^{[k]} \varphi \end{aligned}$$

with  $k \in \{1, 2, 3\}$ .

There  $\beta^{[k]}$  are complex 2-diagonal  $4 \times 4$ -matrices of Clifford's set of rank 4, and  $\varphi$  is matrix columns with four complex components. The light and colored pentads of Clifford's set of such rank contain in threes 2-diagonal matrices, corresponding to 3 space coordinates in according with Dirac's equation. Hence, a space of these events is 3-dimensional.

Let  $\rho(t, \mathbf{x})$  be a probability density of event  $A(t, \mathbf{x})$ , and

$$\rho_c(t, \mathbf{x}|t_0, \mathbf{x}_0)$$

be a probability density of event  $A(t, \mathbf{x})$  on condition that event  $B(t_0, \mathbf{x}_0)$ .

In that case if function  $q(t, \mathbf{x}|t_0, \mathbf{x}_0)$  is fulfilled to condition:

$$\rho_c(t, \mathbf{x}|t_0, \mathbf{x}_0) = q(t, \mathbf{x}|t_0, \mathbf{x}_0)\rho(t, \mathbf{x}), \quad (4)$$

then one is called a *disturbance function*  $B$  to  $A$ .

If  $q = 1$  then  $B$  does not disturbance to  $A$ .

A conditional probability density of event  $A(t, \mathbf{x})$  on condition that event  $B(t_0, \mathbf{x}_0)$  is presented as:

$$\rho_c = \varphi_c^\dagger \varphi_c$$

like to a probability density of event  $A(t, \mathbf{x})$ .

Let

$$\varphi = \begin{bmatrix} \varphi_{1,1} + i\varphi_{1,2} \\ \varphi_{2,1} + i\varphi_{2,2} \\ \varphi_{3,1} + i\varphi_{3,2} \\ \varphi_{4,1} + i\varphi_{4,2} \end{bmatrix}$$

and

$$\varphi_c = \begin{bmatrix} \varphi_{c,1,1} + i\varphi_{c,1,2} \\ \varphi_{c,2,1} + i\varphi_{c,2,2} \\ \varphi_{c,3,1} + i\varphi_{c,3,2} \\ \varphi_{c,4,1} + i\varphi_{c,4,2} \end{bmatrix}$$

(all  $\varphi_{r,s}$  and  $\varphi_{c,r,s}$  are real numbers).

In that case octavian

$$u = \varphi_{1,1} + \varphi_{1,2}i + \varphi_{2,1}j + \varphi_{2,2}k + \varphi_{3,1}E + \varphi_{3,2}I + \varphi_{4,1}J + \varphi_{4,2}K$$

is called a *Caylean* of  $\varphi$ . Therefore, octavian

$$u_c = \varphi_{c,1,1} + \varphi_{c,1,2}i + \varphi_{c,2,1}j + \varphi_{c,2,2}k + \varphi_{c,3,1}E + \varphi_{c,3,2}I + \varphi_{c,4,1}J + \varphi_{c,4,2}K$$

is Caylean of  $\varphi_c$ .

In accordance with the octavian norm definition:

$$\begin{aligned} \|u_c\|^2 &= \rho_c \\ \|u\|^2 &= \rho \end{aligned} \tag{5}$$

Because the octavian algebra is a division algebra [1, 2] then for each octavians  $u$  and  $u_c$  there exists an octavian  $w$  such that

$$u_c = wu,$$

Because the octavians algebra is normalized then

$$\|u_c\|^2 = \|w\|^2 \|u\|^2.$$

Hence, from (4) and (5):

$$q = \|w\|^2.$$

Therefore, in a 3+1-dimensional space-time there exists an octavian-Caylean for a disturbance function of any event to any event.

In order to increase a space dimensionality the octavian algebra can be expanded by a Dickson doubling procedure:

Another 8 elements should be added to basic octavians:

$$z_1, z_2, z_3, z_4, z_5, z_6, z_7, z_8,$$

such that:

$$\begin{aligned} z_2 &= iz_1, \\ z_3 &= jz_1, \\ z_4 &= kz_1, \\ z_5 &= Ez_1, \\ z_6 &= Iz_1, \\ z_7 &= Jz_1, \\ z_8 &= Kz_1, \end{aligned}$$

and for every octavians  $u_1, u_2, v_1, v_2$ :

$$(u_1 + u_2z_1)(v_1 + v_2z_1) = (u_1v_1 - \bar{v}_2u_2) + (v_2u_1 + u_2\bar{v}_1)z_1$$

(here: if  $w = w_1 + w_2i + w_3j + w_4k + w_5E + w_6I + w_7J + w_8K$  with real  $w_s$  then  $\bar{w} = w_1 - w_2i - w_3j - w_4k - w_5E - w_6I - w_7J - w_8K$ ).

It is a 16-dimensional Cayley-Dickson algebra.

In according with [3]: for any natural number  $z$  there exists a Clifford set of rank  $2^z$ . In considering case for  $z = 3$  there is Clifford's seven:

$$\begin{aligned} \underline{\beta}^{[1]} &= \begin{bmatrix} \beta^{[1]} & 0_4 \\ 0_4 & -\beta^{[1]} \end{bmatrix}, \underline{\beta}^{[2]} = \begin{bmatrix} \beta^{[2]} & 0_4 \\ 0_4 & -\beta^{[2]} \end{bmatrix}, \\ \underline{\beta}^{[3]} &= \begin{bmatrix} \beta^{[3]} & 0_4 \\ 0_4 & -\beta^{[3]} \end{bmatrix}, \underline{\beta}^{[4]} = \begin{bmatrix} \beta^{[4]} & 0_4 \\ 0_4 & -\beta^{[4]} \end{bmatrix}, \\ \underline{\beta}^{[5]} &= \begin{bmatrix} \gamma^{[0]} & 0_4 \\ 0_4 & -\gamma^{[0]} \end{bmatrix}, \end{aligned} \quad (6)$$

$$\underline{\beta}^{[6]} = \begin{bmatrix} 0_4 & 1_4 \\ 1_4 & 0_4 \end{bmatrix}, \underline{\beta}^{[7]} = i \begin{bmatrix} 0_4 & -1_4 \\ 1_4 & 0_4 \end{bmatrix}, \quad (7)$$

Therefore, in this seven five 4-diagonal matrices (6) define a 5-dimensional space of events, and two 4-antidiagonal matrices (7) defined a 2-dimensional space for the electroweak transformations.

It is evident that such procedure of dimensions building up can be continued endlessly. But in accordance with the Hurwitz theorem<sup>1</sup> and with the generalized Frobenius theorem<sup>2</sup> a more than 8-dimensional Cayley-Dickson algebra does not a division algebra. Hence, there in a more than 3-dimensional space exist events such that a disturbance function between these events does not hold a Caylean. I call such disturbance *supernatural*.

Therefore, supernatural disturbance do not exist in a 3-dimensional space, but in a more than 3- dimensional space such supernatural disturbance act.

<sup>1</sup>Every normalized algebra with unit is isomorphous to one of the following: the real numbers algebra  $\mathbf{R}$ , the complex numbers algebra  $\mathbf{C}$ , the quaternions algebra  $\mathbf{K}$ , the octavians algebra  $\mathbf{O}$  [1]

<sup>2</sup>A division algebra can be only either 1 or 2 or 4 or 8-dimensional [2]

## References

- [1] {M. L. Cantor, A. S. Solodovnikov, *Hipercomplex numbers*, Moscow, (1973), p.99}, Kantor, I. L.; Solodovnikov, A. S. (1978), *Hyperkomplexe Zahlen*, Leipzig: B.G. Teubner
- [2] O. V. Mel'nikov, V. N. Remeslennikov, et al., *General Algebra*, Moscow, (1990), p.396
- [3] V. A. Zhelnorovich, *Theory of spinors. Application to mathematics and physics*, Moscow, (1982), p.21
- [4] E. Abers, *Quantum Mechanics*. Addison Wesley, (2004), p.423. I
- [5] Gunn Quznetsov, *Final Book on Fundamental Theoretical Physics*, viXra:1111.0051 (2014), pp.60–62