

Metric and algebraic tensors for 4-dimensional uncurved space.

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Abstract

The definition of the algebraic tensor for vector space by using the vector product of vectors from it's basis. And application it to the our 4 – dimensional space.

1) Metric tensor.

Let $\overset{\mathbf{r}}{e}_1, \overset{\mathbf{r}}{e}_2, \overset{\mathbf{r}}{e}_3, \overset{\mathbf{r}}{e}_4$ be the basis of uncurved space W. For these vectors we have :

$$(\overset{\mathbf{r}}{e}_1, \overset{\mathbf{r}}{e}_1) = 1 \quad (\overset{\mathbf{r}}{e}_2, \overset{\mathbf{r}}{e}_2) = -1 \quad (\overset{\mathbf{r}}{e}_3, \overset{\mathbf{r}}{e}_3) = -1 \quad (\overset{\mathbf{r}}{e}_4, \overset{\mathbf{r}}{e}_4) = -1 \quad (1.1)$$

Another scalar products of these vectors equal zero. So

$$g_{11} = 1 \quad g_{22} = -1 \quad g_{33} = -1 \quad g_{44} = -1 \quad (1.2)$$

Another components of metric tensor for W also equal zero. If we take vector

$$d\overset{\mathbf{r}}{R} = \overset{\mathbf{r}}{e}_1 \cdot i \cdot dt + \overset{\mathbf{r}}{e}_2 \cdot dx + \overset{\mathbf{r}}{e}_3 \cdot dy + \overset{\mathbf{r}}{e}_4 \cdot dz \quad (1.3)$$

then we get :

$$(d\overset{\mathbf{r}}{R}, d\overset{\mathbf{r}}{R}^*) = dt^2 - dx^2 - dy^2 - dz^2 = ds^2 \quad (1.4) \quad ds - \text{interval.}$$

If we take vector

$$\overset{\mathbf{r}}{P} = \overset{\mathbf{r}}{e}_1 \cdot i \cdot E + \overset{\mathbf{r}}{e}_2 \cdot P^x + \overset{\mathbf{r}}{e}_3 \cdot P^y + \overset{\mathbf{r}}{e}_4 \cdot P^z \quad (1.5)$$

then we get :

$$(\overset{\mathbf{r}}{P}, \overset{\mathbf{r}}{P}^*) = E^2 - (P^x)^2 - (P^y)^2 - (P^z)^2 = m^2 \quad (1.6) \quad m - \text{rest mass of a particle}$$

Here we use :

$$i \cdot i = -1 \quad i^* = -i \quad (1.7)$$

2) Algebraic tensor.

$$\text{Equation } [\overset{\mathbf{r}}{e}_\mu \times \overset{\mathbf{r}}{e}_\nu] = \overset{\mathbf{r}}{e}_\sigma \cdot F^{\sigma}_{\mu\nu} \quad (2.1)$$

defines the algebra for basis $\overset{\mathbf{r}}{e}_\mu$. In our space this algebra coincides with quaternion algebra :

	$\overset{\mathbf{r}}{e}_1$	$\overset{\mathbf{r}}{e}_2$	$\overset{\mathbf{r}}{e}_3$	$\overset{\mathbf{r}}{e}_4$	$\overset{\mathbf{r}}{e}_\nu$
$\overset{\mathbf{r}}{e}_1$	$\overset{\mathbf{r}}{e}_1$	$\overset{\mathbf{r}}{e}_2$	$\overset{\mathbf{r}}{e}_3$	$\overset{\mathbf{r}}{e}_4$	
$\overset{\mathbf{r}}{e}_2$	$\overset{\mathbf{r}}{e}_2$	$-\overset{\mathbf{r}}{e}_1$	$\overset{\mathbf{r}}{e}_4$	$-\overset{\mathbf{r}}{e}_3$	
$\overset{\mathbf{r}}{e}_3$	$\overset{\mathbf{r}}{e}_3$	$-\overset{\mathbf{r}}{e}_4$	$-\overset{\mathbf{r}}{e}_1$	$\overset{\mathbf{r}}{e}_2$	
$\overset{\mathbf{r}}{e}_4$	$\overset{\mathbf{r}}{e}_4$	$\overset{\mathbf{r}}{e}_3$	$-\overset{\mathbf{r}}{e}_2$	$-\overset{\mathbf{r}}{e}_1$	
$\overset{\mathbf{r}}{e}_\mu$					

Table 1

Let us take vector

$$\overset{\mathbf{r}}{R} = \overset{\mathbf{r}}{e}_1 \cdot i \cdot t + \overset{\mathbf{r}}{e}_2 \cdot x + \overset{\mathbf{r}}{e}_3 \cdot y + \overset{\mathbf{r}}{e}_4 \cdot z \quad (2.2)$$

Then, using also (1.5) we derive :

$$[\overset{\mathbf{r}}{P} \times \overset{\mathbf{r}}{R}^*] = \overset{\mathbf{r}}{M} = \overset{\mathbf{r}}{e}_1 \cdot M^1 + \overset{\mathbf{r}}{e}_2 \cdot M^2 + \overset{\mathbf{r}}{e}_3 \cdot M^3 + \overset{\mathbf{r}}{e}_4 \cdot M^4 \quad (2.3)$$

$$M^1 = E \cdot t - P^x \cdot x - P^y \cdot y - P^z \cdot z \quad (2.4)$$

$$M^2 = (P^y \cdot z - P^z \cdot y) + i \cdot (E \cdot x - P^x \cdot t) \quad (2.5)$$

$$M^3 = (-P^x \cdot z + P^z \cdot x) + i \cdot (E \cdot y - P^y \cdot t) \quad (2.6)$$

$$M^4 = (P^x \cdot y - P^y \cdot x) + i \cdot (E \cdot z - P^z \cdot t) \quad (2.7)$$

If we demand that M^k were real at any k , then we get :

$$E \cdot x = P^x \cdot t \quad (2.8)$$

$$E \cdot y = P^y \cdot t \quad (2.9)$$

$$E \cdot z = P^z \cdot t \quad (2.10)$$

Taking into account that $P^k = E \cdot V^k$ (2.11)

and $x = V^2 \cdot t$, $y = V^3 \cdot t$, $z = V^4 \cdot t$ (2.12)

we can say that equations (2.8), (2.9), (2.10) describe the motion of a particle. And real M^2, M^3, M^4 then describe orbital angular momentum of the particle.

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