

# On the Cold Big Bang Cosmology and the Flatness Problem

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In my papers [3] and [4], I obtain a Cold Big Bang Cosmology, fitting the cosmological data, with an absolute zero primordial temperature, a natural cutoff for the cosmological data to a vanishingly small entropy at a singular microstate of a comoving domain of the cosmological fluid. Now, in this brief paper, we show the energy density of the  $t$ -sliced universe must actually be the critical one, following as consequence of the solution in [3] and [4]. It must be pointed out that the result obtained here in this paper on the flatness problem does not contradict the solution in [3], viz., does not contradict the open universe, with  $k = -1$ , obtained in [3], since the solution in [3] had negative pressure and negative total cosmological energy density, hence lesser than the critical positive density. The critical density we obtain here is due to the positive fluctuations I previously discussed regarding the Heisenberg mechanism in [4]. Hence, the energy density due to fluctuation turns out to be positive, the critical one, this being calculated in [4] from the fluctuations within the  $t$ -sliced spherical shell at its  $t$ -sliced hypersurface, being the total energy density that generates the fluctuations, actually, negative in [3] and [4], hence, again, lesser than the critical one and supporting  $k = -1$ . These results are complementary and support my previous results.

## 1 Explaining the observed critical energy density for the cosmological substratum

The proof is straightforward. Recalling the  $t$ -sliced spherical shell defined in [4], we proved its energy at the instant  $t$  of the cosmological time is given by:

$$N_t \delta E_\rho = E^+(t) = \frac{E_0^+}{\sqrt{1 - \dot{R}^2/c^2}}, \quad (1)$$

being the Eq. (6) in [4] and the Eqs. (56) and (59) in [3]. From the Eq. (36) in [3], the Eq. (1) is simply rewritten:

$$E^+ = \frac{c^4 R}{2G}, \quad (2)$$

from which the energy density of the  $t$ -sliced spherical shell pertaining to its  $t$ -sliced hypersurface of simultaneity, being full of its  $t$ -simultaneous cosmological points of the substratum, turns out to be given by:

$$\rho_t = \frac{c^4 R}{2G} \frac{3}{4\pi R^3} = \frac{3c^4}{8\pi G} \frac{1}{R^2}. \quad (3)$$

Now, once defined the Hubble parameter for the spherical shell of the  $t$ -sliced hypersurface, centered at its comoving origin, namely  $H_t$ :

$$H_t = \frac{\dot{R}}{R}, \quad (4)$$

one obtains the energy density of the  $t$ -sliced spherical shell pertaining to its  $t$ -sliced hypersurface of simultaneity from the Eqs. (3) and (4):

$$\rho_t = \frac{3c^2 H_t^2}{8\pi G} \frac{c^2}{\dot{R}^2}. \quad (5)$$

Now, Eq. (35) in [3] shows  $\dot{R}$  very rapidly reaches  $c$ . Hence, the energy density  $\rho_t$  of the  $t$ -sliced spherical shell pertaining to its  $t$ -sliced hypersurface of simultaneity actually is the very critical energy, viz.:

$$\rho_t \rightarrow \rho_{\text{crit}} = \frac{3c^2 H_t^2}{8\pi G}, \quad (6)$$

very rapidly, and it will continue being the critical energy density  $\rho_{\text{crit}}$  forever\*. The time domain within which  $\rho_t$  deviates from the critical is  $t \approx 0$ , at which it reads:

$$\rho_0 = \frac{3c^7}{16\pi G^2 h} \approx 10^{111} \text{ Jm}^{-3}, \quad (7)$$

provided the energy density  $\rho_t$  of the  $t$ -sliced spherical shell being given by the Eq. (3) and with the  $R_0 = \sqrt{2Gh/c^3}$  obtained in [3]. The Hubble parameter vanishes at  $t = 0$  in our model†.

## Universe is initially fine-tuned, naturally, due to its initially vanished entropy

Under the solution in [3], there is just an unique initial state for the cosmological substratum with its initial null entropy. As discussed in [3], Eq. (39), the temperature  $T$  of the cosmological substratum at each comoving domain vanishes at  $t = 0$ . Once permeated by radiation, the local entropy turns out to be proportional to  $T^3$ , with a local volume  $V$  of the  $t$ -sliced substratum proportional to  $R_0^3 \neq 0$ , with  $R_0 = \sqrt{2Gh/c^3}$  as

\*It will continue being given by the Eq. (6), but its value tends to vanish, as one easily infers from the Eq. (3). See [3] and [4].

†See [3] and [4].

discussed in [3]. Hence, the initial local entropy, at  $t = 0$ , at each comoving domain reads:

$$S|_{t=0} = \frac{32\pi^5 k_B^4}{45h^3 c^3} VT^3 \Big|_{t=0} = 0, \quad (8)$$

Hence, the solution in [3] naturally provides the required flatness, given by the Eq. (6) here, with an unique initial condition for the cosmological substratum, the minimal, initially vanished, entropy.

### Conclusion

We conclude the cosmological model obtained in [3] and [4] provides the critical energy density for the  $t$ -sliced spherical shell pertaining to its  $t$ -sliced hypersurface of simultaneity of the cosmological substratum. One may explain the initial condition *ex post* in [3], since the vanished entropy explains the unicity for the solution in [3]. It is interesting the indeterminacy principle turns out to provide such an unicity for the solution.

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