

Compactly reproducing the fine structure constant inverse and the muon-, neutron-, and proton-electron mass ratios

J. S. Markovitch
P.O. Box 752
Batavia, IL 60510
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Compact equations are introduced that reproduce the fine structure constant inverse and the muon-, neutron-, and proton-electron mass ratios near their experimental limits.

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The fine structure constant (FSC) and the muon-, neutron-, and proton-electron mass ratios can be economically reproduced as follows. Firstly, define

$$\begin{aligned} l_0 &= \frac{1}{M^2} & q_0 &= \frac{1}{M^3} \\ l_1 &= \frac{[M - l_0/3M^2]^2}{N^{-2}} & q_1 &= \frac{M^2 - q_0}{N^{-2}} \end{aligned} .$$

Similarly, define

$$\begin{aligned} l_2 &= \frac{M^3 - l_0}{N} & q_2 &= \frac{M^3 - q_0}{N} \\ l_3 &= \frac{[M - l_0/3M^2]^3}{N} & q_3 &= \frac{[M - q_0/3M^2]^3}{N} \end{aligned}$$

which are symmetric under $l \leftrightarrow q$, so that for

$$M = 10 \quad \text{and} \quad N = 3$$

the FSC inverse can be approximated four ways

$$\begin{aligned} \frac{l_1 + l_2}{N^2} &= 137.036\ 000\ 001\ 111 & \frac{q_1 + q_2}{N^2} &= 137.036 \\ \frac{l_1 + l_3}{N^2} &= 137.036\ 000\ 002\ 346 & \frac{q_1 + q_3}{N^2} &= 137.036\ 000\ 000\ 012 \end{aligned}$$

which are also symmetric under $l \leftrightarrow q$. Up to this point, all definitions and approximations exactly follow [1]. Now define three functions

$$\begin{aligned} R(n, p) &= \frac{(n + M^2)^n}{(mn - M^n)^{n-2}} (|p| - M^{2p})^p & e(p) &= R(-1, p) \\ & & s(p) &= R(+1, p) \end{aligned}$$

where $n = \pm 1$, and p equals -1 , 0 , or $+1$. Then, for $m = 4$ these functions allow the muon- and neutron-electron mass ratios to be economically approximated by

$$\begin{aligned} \frac{e(1)/q_0 - 1}{l_2 - l_0} &= 206.768\ 270\ 731 & \frac{e(1)/l_0 - s(0)/q_0}{q_2 - q_0} &= 1838.683\ 654\ 735 \\ \frac{e(1)/q_0 - 1}{l_3 - l_0} &= 206.768\ 270\ 724 & \frac{e(1)/l_0 - s(0)/q_0}{q_3 - q_0} &= 1838.683\ 654\ 734 \end{aligned}$$

each with three terms (of four) also symmetric under $l \leftrightarrow q$. Moreover, the proton-electron mass ratio can be approximated by

$$\frac{e(-1)/l_0 - s(-1)/q_0}{\frac{M^3}{N}} = 1836.152\ 675\ 237 \quad .$$

Also note that

$$\frac{M^3}{N} = \frac{10^3}{3} \approx l_2 - l_0 \approx l_3 - l_0 \approx q_2 - q_0 \approx q_3 - q_0 \quad .$$

With the exception of the less precisely measured muon-electron mass ratio, which above is reproduced at its experimental limit, all of these values are within just a few parts per billion (ppb) of their 2006 CODATA values [2]:

137.036 000 002	206.768 270 7	1838.683 654 7	1836.152 675 24
137.035 999 679 (94)	206.768 282 3 (52)	1838.683 660 5 (11)	1836.152 672 47 (80)

All values are identical to those produced earlier in [1, 3].

At this point the reader might speculate that n in the expression $n + M^2$ allows the above approximations to be *fine-tuned* to fit their corresponding experimental mass ratios. But n simply does not permit the precise fine-tuning needed to fit these ratios. Two out of three of the mass ratios are known and fit to within a few ppb, whereas the value n (equaling ± 1) represents 10 *million* ppb of M^2 . It follows that although n is responsible for *some* fine-tuning, it cannot begin to explain the precise fit achieved for the mass data. Identical reasoning applies to p .

In the same way, the neutron- and proton-electron mass ratios may be viewed as having values close enough to each other to assure that a good approximation of one can readily be fine-tuned into a good approximation of the other (which is to say that their earlier approximations are not truly independent). Actually, the neutron- and proton-electron mass ratios differ by 1.4 *million* ppb, an enormous gap to overcome by fine-tuning.

Likewise, the use of -1 in the muon equation's numerator $e(1)/q_0 - 1$ might also be seen as a case of fine-tuning. But letting $m = 0$ causes $e(1)/q_0 - 1$ to equal 0, implying that -1 actually reflects underlying order.

Similarly, M and N may be regarded as "adjusted to fit experiment." But, as demonstrated elsewhere (see [1]), the values used for M and N (10 and 3, respectively) are the smallest positive integers that cause all four of the above FSC inverse approximations to produce nearly equal values. Hence, M and N are not adjusted to fit experiment, but acquire their values automatically.

Moreover, in [4] two simple symmetric mathematical identities are employed as a starting point to automatically generate the values 10, 3, and 137.036, while in [5] a brute-force computer search for efficient approximations of the FSC inverse also automatically finds 10, 3, and 137.036 (more specifically, it finds $(q_1 + q_2)/N^2 = 137.036$).

And, perhaps most significantly, in [6] a mixing model from 2007 is described, one which employs the constants 10 and 3 in a way that derives from $(l_1 + l_3)/N^2$. This model is shown to have correctly predicted the (arguably improbable) changes subsequently observed in the experimental quark mixing angles.

More generally, in [7] information theory and number theory are employed to demonstrate that even the comparatively primitive 2004 versions of the above muon- and neutron-electron mass ratio approximations manage to *compress* the mass data they reproduce, something one would not expect to happen by chance (see Eqs. (14a) and (14b) in [7]).

Finally, note that:

$$\begin{array}{llll}
 l_0 = 0.01 & l_1 = 899.99400001 & l_2 = 333.33 & l_3 = 333.330000011111 \dots \\
 q_0 = 0.001 & q_1 = 899.991 & q_2 = 333.333 & q_3 = 333.333000000111 \dots
 \end{array}$$

And:

$$\begin{array}{ll}
 e(+1) = +4.1^3 & -s(0) = 600 \times 101 \\
 e(-1) = -4.1^3/9.9^2 & -s(-1) = 600 \times 101/0.99
 \end{array}$$

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